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The sound of high winds:
the effect of atmospheric stability
on wind turbine sound and microphone noise

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* This is a copy of GP (Frits) van den Berg's doctoral thesis, from the University of Groningen, Netherlands, completed May 2006. Dr. van den Berg kindly sent this to me earlier this month.

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The sound of high winds:

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and microphone noise

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I WIND POWER, SOCIETY, THIS BOOK: an introduction

Bobby asks: 'Do you ever hear the windmills?'
'What sound do they make?'
'It's a clanking metal noise, but when the wind is really strong the blades blur and the air starts screaming in pain.' He shudders.
'What are the windmills for?'
'They keep everything running.
If you put your ear to the ground you can hear them.'
'What do you mean by everything?'
'The lights, the factories, the railways. Without the windmills it all stops.'

This is the story of the discovery of a new phenomenon: why wind turbines sound different at night time. This discovery was related to a problem in society, namely that of perceived noise by residents living close to such turbines.

This introduction sketches the context in which my work proceeded: how the questions came up, why noise is an inseparable part of wind power development, and that being critical does not need to imply a negative attitude towards wind power. Let's start at the beginning.

1.1 A ‘new’ phenomenon

The discovery was modest: I have not found a new law of nature or a new way to make money. It was rather the idea to apply existing knowledge in a new context: the application of atmospheric physics to solve the mystery why people complained about noise from wind turbines that according to wind developers and acoustic consultants they should not even be able to hear. In principle it was not very difficult to find out why. When Walter Flight (a very Dutch citizen despite his name) told me he could see the wind turbines near his house rotating at high speed while at the same time his garden was completely calm, I thought: oh yes, I know that, that’s

1 'The suspect', by Michael Robotham, Time Warner Paperbacks, 2003 (p. 151)
because at night, especially on nice summer evenings, the atmosphere becomes stable. I teach this in a course, Environmental Techniques. The phenomenon is treated extensively in this book, but for now it is sufficient to know that, due to strong winds at greater heights coupled with very light winds at ground level, wind turbines can be a lot noisier in a night time atmosphere than they are in daytime. This was why Walter and his neighbours complained. Also the nature of the sound changes: a thumping character can become very pronounced at night.

In this book I will often use the terms 'day' and 'night', though the distinction is more accurately stated as the atmosphere being unstable (which is usually in daytime, that is: sun up) or stable (night time, sun down). The heat coming in from the sun or radiated out at night is the real cause of the difference in stability. In between is another state, namely neutral, where heating or cooling are unimportant because of heavy clouding and/or strong wind and which can occur in day as well as night time, though not very often in a temperate climate and over land. Atmospheric stability means that vertical movements in the air are damped and as a consequence horizontal layers of air can have a greater difference in velocity: close to the ground the wind can be weak while higher up there is a strong wind.

Though in principle the explanation is simple and easily understood, it of course had to be shown from solid theory and with sufficient data that the explanation was correct. The first steps were extensive measurements in Bellingwolde, where severe complaints had arisen about noise from the nearby Rhede wind farm. This I did together with Richard de Graaf, then a physics student.

After this simple discovery, a new mystery (to me) was why this did not play a role in the assessment of wind turbine noise? Every meteorologist knows about atmospheric stability, so why had none of the experts dealing with wind turbine sound ever come across it? Wind turbines have been built for several decades and since the 1980’s in ever larger numbers, so there should be a lot of accumulated experience. Had no one (except some
residents) noticed the discrepancy between predicted and real noise exposure?
There are probably several reasons. One of them is that for a long time wind turbines were not big enough for the effects of atmospheric stability to be clearly noticeable. Since wind turbines have grown taller the effect manifests itself more clearly. Secondly, as the more distant locations have become scarce, more and more turbines are being built closer to where people live, so more people now experience the sound of wind turbines. Thirdly, atmospheric stability over flat land is easier to understand and quantify than in a mountainous or coastal area where the atmosphere is more complex so the effect on wind turbines may be less easily recognizable.
Wind turbines as such have not become that much noisier, despite their increase in height and blade span (the sound power depends more on speed than on physical dimensions of the towers). Earlier machines could be quite noisy due to whining or severe thumping, and modern designs are certainly better. The point is they now reach into less familiar parts of the atmosphere.

Finally, an important reason to not recognize the unexpected high sound levels certainly is the fact that it impedes commercial interests and national policy. The positive ring of the term 'sustainability' helps investors in wind energy and local authorities (applying national policy) to counterbalance objections concerning possible disadvantages of new projects. As these objections are sometimes strong enough to torpedo projects, investors and authorities don't welcome more negative news. Though the population widely supports sustainable energy, reactions are less positive when a new project adversely affects their lives. This 'contradictory behaviour' is in fact quite understandable: when a new project is planned in an area, residents for the first time have to balance the positive social consequences to the negative local impact: visual impact, flickering shadows, noise and possibly ice throw from turbine blades.

The first reaction of wind energy proponents, represented by the Windkoepel (‘Wind dome’), to our research results was to pay a consultant
to comment on our report [Van den Berg et al 2002]. This consultant boasted of having advised a large number of wind farm projects, so he clearly understood the position of the wind power industry. In the resulting ‘second opinion’ [Kerkers 2003] no material critique was presented, only procedural arguments were used to declare our results inaccurate and thus irrelevant. The Windkoepel issued a press statement concluding that we had made a lot of fuss, but had not contributed any new insights.\(^1\) They could get back to business.

1.2 Digging deeper

I too went back to my business, which can be summarized as helping citizen groups to defend their position by objective arguments using known principles of physics. In 2004 an article about my research was published in a scientific journal [Van den Berg 2004a] lending my results the respectability of peer review and triggering an international e-mail influx from interested consultants as well as worried residents, as our first report had done earlier on a national scale.

What still puzzled me at that time was how a single turbine could start thumping at night. I thought I understood how the modest blade swish of a single turbine could evolve into louder thumping: the small sound variations due to blade swish from several turbines could add up to louder pulses. But with a single turbine there is nothing to add! Apart from this, in news media in the UK there were complaints that low frequency wind turbine noise had been underestimated and had been making people sick.\(^2\)

Some thoughts about this were presented at a conference in Maastricht [Van den Berg 2004b]. I agreed with delegate Jørgen Jakobsen, who presented a paper on low frequency wind turbine noise [Jakobsen 2004],

\(^1\) Press statement February 2, 2003 “Onlangs is opschudding ontstaan …..,” (“Recently an upheaval was caused…”), De Windkoepel, Arnhem

that even though wind turbines did produce an appreciable amount of infrasound, the level was so far below the average human hearing threshold that it could not be a large scale problem. But it was possible that complaints had been expressed in a way not understood by experts. Perhaps people bothered by the endless thumping of a relatively low pitched sound (such as I had heard myself on several occasions), thought that 'low frequency sound' was a term to use, as official sounding jargon. They might not be aware that the term 'low frequency sound' makes acousticians think of frequencies below 100 to 200 hertz, and in that range the sound level was not considered to be problematic. A classical misunderstanding perhaps, that could be clarified. After the Maastricht conference I wanted to quantify my ideas on the origin of the night time thumping of wind turbines and the relevance of low frequencies. This resulted in a second scientific article [Van den Berg 2005a] in which I tried to put these ideas together.

What had surprised me from early on was that people in the wind power business seemed to know so little about their raw material, the wind. In the Windkoepel press statement (see footnote previous page) a wind turbine manufacturer’s spokesman argued that if the hub height wind velocity indeed was structurally higher at night, this must be visible in production statistics. This indeed seems plausible, so why not investigate that? If the wind industry had done so, they might have come up with results I found from measured wind profiles at Cabauw over an entire year [Van den Berg 2005b]. Indeed for an 80 m high turbine the night time yield is significantly higher than expected, whereas the daytime yield is lower. The net result was that in the real atmosphere at Cabauw annual production was 14% to 20% (depending on wind turbine power settings) higher than in an atmosphere extrapolated from 10-m wind velocities with a perpetual neutral wind profile. For wind power production forecasting there is a method that incorporates a correction for atmospheric stability [Troen et al 1989], but such knowledge has never been used for sound exposure forecasting.
I.3 **Commercial and policy implications**

So from an energy point of view a stable atmosphere is very attractive. The challenge is to use that potential, but not put the burden on those living nearby. One solution is to build wind farms offshore where no people are affected if enough distance is kept (and calculation models are used that accurately model long range sound propagation over water). Over large bodies of water seasonal, not diurnal atmospheric stability will boost production in part of the year but lower it when the water has warmed. Another solution is to improve turbine design from two perspectives: decreasing sound power without substantially decreasing electric power, and reducing annoyance by minimizing fluctuations in the sound. Part of any solution is to respect complainants and try to achieve a better balance between national benefits and local costs.

Oblivious of any research, residents had already noticed a discrepancy between predicted and real noise exposure. Opponents of wind farms have organized themselves in recent years in the Netherlands and elsewhere, and word had spread that noise exposure in some cases was worse than predicted. Though atmospheric stability and sometimes a malfunctioning turbine could explain this, most wind farm developers and their consultants relied on the old prediction methods. An energy firm’s spokesman complained that each and every new project attracted complaints (from local groups) and called this “a new Dutch disease”.¹ This is a very narrow view on the problem, denying the detrimental effects for residents. If their real concerns are denied it is not unreasonable for residents to oppose a new project, because practical experience shows that once the wind farm is there (or any other noise producer) and problems do arise, complaints will very probably not alter the situation for at least several years. Social scientists are familiar with such situations and suggest better strategies such as being honest and respectful, treating residents as equal partners, and not being arrogant: already in 1990 Wolsink mentioned this in a study on acceptance of wind energy and warned that it was wrong to label opposition as NIMBY (Not In My Back Yard) and refuse to recognize

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¹ NRC Handelsblad, August 26 2005: "Verzet tegen windmolens succesvol" ("Opposition to wind mills succesful")
It is sad that most of the proponents still emanate a WARYDU attitude (We Are Right but You Don’t Understand).

When real complaints are not addressed seriously, the “new Dutch disease” may well become an Australian, British, Chinese or any nation’s disease. In the Netherlands assessment of wind turbine noise still is according to the old standard procedure (with one exception, see chapter VII), assuming a neutral atmosphere at all times, even though this has been admitted to be wrong for more than a year now. Consultants apparently are afraid to be critical, perhaps because they don’t want to jeopardize new assignments or because a change in assessment implies they were not correct before (they were not correct, but we were wrong collectively). Though most consultants claim to be impartial, the problem of ‘not biting the hand that feeds’ is more subtle, as I concluded in an earlier desk study on the quality of acoustic reports [Van den Berg 2000]. E.g., it involves authorities who do not question the position of paid experts, and a society hiding political decisions behind the demand for more research.

I hope other countries do not to follow the Dutch way: first denying the consistency and legitimacy of the complaints, then being late in addressing them and in the end finding this has created more opposition. It is evident that also in the UK there are (a few?) serious complaints from honest people that are not dealt with adequately. In at least some cases atmospheric stability again seems to offer an explanation for observations of unpleasant wind turbine noise by residents (see example in box on next page), but the matter has not been investigated correctly.

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1 In March 2004 I showed in an article in ‘Geluid’, a Dutch professional journal, how to deal with non-neutral atmospheric conditions within the existing legal procedures [Van den Berg 2004c]; in July 2004 the Ministry of Housing, Environment and Spatial Planning advised to investigate the ‘wind climate’ at new wind farm locations (letter on “Beoordeling geluidmetingen Natuurkundewinkel RUG bij De Lethe, gem. Bellingwedde” to Parliament by State Secretary van Geel, June 21, 2004); in the 2005 Annual report of BLOW, a union of local, provincial and national authorities to promote wind energy development, it is recognized that the effect of wind shear still should be addressed, but no action is announced (Annual report BLOW 2005, January 2006).
NOISE FROM WINDFARM MAKING LIFE A MISERY

A recent settler in Caithness claimed yesterday his life is being blighted by ghostly noises from his new neighbours, the county's first large-scale windfarm. (.....) Mr Bellamy said: "The problem is particularly bad at night when I try to get to sleep and there's a strong wind coming from the direction of the turbines. "They just keep on droning on. It's a wooh wooh type of sound, a ghostly sort of noise. It's like torture and would drive anyone mad."

Mr Bellamy believes the noise is being transmitted through the ground since it seems to intensify when he lies down. He said he has got nowhere with complaints to the wind company and environmental health officers. "I feel I'm just getting fobbed off and can't get anyone to treat me seriously," he said. Mr Bellamy has been asked to take noise readings every 10 minutes during problem times, something he claims is unrealistic to expect him to do. He said the company's project manager Stuart Quinton-Tulloch said they could not act until it had proof of unacceptable noise levels. Mr Bellamy said: "I'm not the moaning type and I have no problem with the look of the windmills. I'm not anti-windfarm. It's just the noise which is obviously not going to go away." (.....)

Highland Council's principal environment officer Tom Foy who has been dealing with Mr Bellamy's complaint was unavailable for comment. His colleague David Proudfoot said he was aware of noise complaints about the Causewaymire turbines being lodged by two other residents, but said he had gone out several times and found no evidence to support the concerns.

Part of an article in Press and Journal of Aberdeen, 25 May 2005
Thinking that this could perhaps be solved by the Sustainable Development Commission (SDC), the UK government’s ‘independent advisory body on sustainable development’. I wrote to the SDC about remarks on wind turbine noise in their report “Wind power in the UK” [SDC 2005], which was in my opinion too positive and somewhat overly optimistic regarding wind turbine noise. The SDC replied, on authority of its (unknown) consultants, that they had no detailed knowledge of atmospheric conditions in the UK but still thought an impulsive character of the noise ‘likely to be very rare’. After I presented some examples the SDC preferred to close the discussion.

The situation in the Netherlands is not very different. In the latest annual report of the body of national, provincial and local authorities responsible for wind energy development it is acknowledged that the problem of underrated noise has justly been brought to the policy agenda. Nevertheless, no activity is undertaken to remedy this.

1.4 Large scale benefits and small scale impact

Though wind turbine noise is the main topic of this book, it is not the main problem in wind power development. Visual impact is usually considered the most important and most discussed local or regional effect. It is often presented as a matter of individual taste, though there are some common factors in ‘public taste’. One such factor is the perceived contrast of a wind turbine (farm) and its environment: a higher contrast will have more impact, either in a positive or negative way. A peculiarity of turbines is that the rotational movement makes them more conspicuous and thus enhances visual impact. This common notion suggests that wind turbines in a built up area will have less impact relative to a remote natural area (though this may be overruled by the number of people perceiving the impact).

A second factor is attitude: e.g. farmers usually have a different attitude to the countryside than ‘city folk” have, and hence they differ in judgments on the appropriateness of a building, construction or activity in the

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countryside. It is predictable that when residents have a positive association with a neighbouring wind farm they will experience less annoyance from the visual impact. For a wind turbine owner the sound of each blade passing means another half kWh is generated\(^1\) and is perhaps associated with the sound of coins falling into his lap, a lullaby. The very same rhythm, like the proverbial leaking faucet tap, might prevent his neighbour from falling asleep.

Other issues have gained attention in the public discussion, such as the modest contribution of wind energy to total energy consumption and the problematic variability of wind power. This is not the place to discuss these issues, except that they partially depend on a person’s world view and expectations of the future. But I would like to show my personal position here. I find it astounding to realize that *all* wind turbine energy generated in the Netherlands in one year (2004) is equal to two months’ growth of the total Dutch energy consumption. And even though wind turbine energy now provides about 2% of the total Dutch electricity consumption, this is only 0.2% of our total energy consumption.\(^2\) This is also true on a global scale as is clear from figure 1.1: wind power is now negligible and expected to supply 0.5% in 2030.

Despite the disappointingly low percentages I still think that wind energy need not be insignificant. In my view the problem is rather that we use such vast amounts of energy and keep on using ever more, which is a problem that no source, including wind power, can solve. Society will need to find a stand in the variety of opinions that have been brought forward since the 1970’s. In a recent newspaper discussion about the liberalization of the energy market an opinion maker stated: “It is now generally appreciated that the end of the rich era of energy approaches rapidly, and the competition has begun for the last stocks”, whilst his opponent the Minister or Economic Affairs wrote: “The lights must be kept burning, the

\(^1\) when the turbine generates 2 MW at 20 rpm

\(^2\) the percentages are based on data from Statistics Netherlands (Centraal Bureau voor Statistiek) for the Netherlands for the year 2004: wind energy production: 1.9 TWh; total electricity consumption: 108.5 TWh; total energy consumption: 919 TWh. Growth in total energy consumption in period 1995 – 2004: + 100 TWh or 1.7 TWh per two months. Growth in total electricity consumption 1995 - 2004: +23 TWh or 2.3 TWh per year.
gas must keep flowing”. I do not agree with the Minister: I think that a limited resource should require limited consumption, even at the cost of some discomfort to our spoiled society. If we can curb our Joule addiction, wind power may help us to produce part of the sustainable energy we need to satisfy basic needs.

Wind turbine noise is a problem that may grow due to neglect by wind energy proponents and thus it may be another reason for part of the public, with politicians following, to turn away from wind power. This problem can be solved when it is also addressed at the level of local impact: sustainability must also apply at the local level. Some technical possibilities for noise reduction are given in this book and more competent, hardware oriented people may come up with better solutions. In addition to this, the social side of the problems must not be neglected. In a recent study [Van As et al 2005] it was concluded that “growing public resistance

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1 NRC Handelsblad 8-11-2005, articles “Bezinning nodig over energiebeleid” (“Energy policy needs reflection” by W. van Dieren) and “Nieuw debat schept slechts onzekerheid” (“New debate only creates uncertainty” by Laurens Jan Brinkhorst); my translations
to onshore wind turbines” obstructs wind energy development in the Netherlands. According to the report this opposition is now the main bottle-neck: local communities and residents are faced with the disadvantages whilst others (proponents, society at large) reap the benefits. The report recommends that the former share in the benefits too.

1.5 Microphone wind noise

In contrast to the impact my wind turbine research has had in society, the same knowledge of atmospheric physics helped me solve a non-controversial problem of interest to only a few: what is the nature of the noise that wind creates in a microphone? It occurred to me that if atmospheric turbulence was the cause, then one must be able to calculate the level of this noise. I was delighted when I found out how well theoretical considerations fitted hitherto only vaguely understood measurement results. Eureka!, such is the joy of work in science.

Somewhat unexpectedly this second discovery turns out to be related to wind turbine sound, which is why it is in this book. Originally it was considered difficult to measure wind turbine sound, because the strong winds that were supposed to cause high wind turbine sound levels, also were believed to be responsible for a lot of microphone wind noise. Solutions to this problem were either to put the microphone out of the wind on the ground or use several microphones and decrease microphone noise by averaging over all microphone signals. A new solution offered in this book is to take measurements in a stable atmosphere where near-ground wind velocity is so low that microphone noise is far less of a problem. One can measure sound at distances from a wind farm most researchers would not now believe to be possible.

The relationship is even stronger. In some countries the level of ambient background sound determines (part of) the limit imposed on sound exposure. To measure the level of this background sound the microphone must be put up in a place where residents stay outdoors, also in stronger winds. In this case it is important to discriminate between real ambient
sound and the noise that wind produces in the microphone. With the calculation methods in this book it is now possible to do so.

1.6 Research aims
The issues raised above concerning wind turbine noise and its relationship to altitude dependent wind velocity led to the following issues to be investigated:

- what is the influence of atmospheric stability on the speed and sound power of a wind turbine?
- what is the influence of atmospheric stability on the character of wind turbine sound?
- how widespread is the impact of atmospheric stability on wind turbine performance: is it relevant for new wind turbine projects?; how can noise prediction take this stability into account?
- what can be done to deal with the resultant higher impact of wind turbine sound?

Apart from these directly wind turbine related issues, a final aim was to address a measurement problem:

- how does wind on a microphone affect the measurement of the ambient sound level?

1.7 Text outline and original work
This book gives an overview of results of the wind turbine noise research that has been presented in the international arena in the last few years, as well as some opinions on this topic in the Introduction and Epilogue. Most of the text in this book has been published in scientific journals or presented at conferences. However, the texts have been adapted somewhat so as to form a continuous story without too much overlap. Other changes have been listed below.

- Chapter II is a reflection on some problems I encountered in doing research and presenting the results, most of it concerning wind turbine noise, but set against a more general background. It corresponds to a
paper presented at Euronoise 2003 [Van den Berg 2003], but some overlap with later chapters is taken out and some new information concerning the variation of wind turbine sound has been added (last paragraph in II.2). The remaining text has been edited slightly.

* Chapter III gives some numbers on wind energy development in the European Union, as well as an introduction on atmospheric wind gradients and the origins of aerodynamic wind turbine sound. It corresponds to sections of two published papers [Van den Berg 2004a and 2005a] to which remarks on the local wind speed at the turbine blade (section III.3) and on the spectrum of thickness sound (footnote in III.4) has been added. Also a description of sound and effects as given by a residential group with practical experience is added (box at end of chapter) and a remark on constant speed and variable speed wind turbines (in III.4).

* Chapter IV corresponds to my first paper on this topic [Van den Berg 2004a] on measurements at the Rhede wind farm. The section on Impulsive Sound has been taken out here and transferred to the next chapter. A new section (IV.10) has been added describing previously unpublished measurements at the Rhede wind farm as well as a comparison with calculated sound levels. Chapter IV demonstrates the fact that sound levels due to wind turbines have been systematically underestimated because hub height wind velocities were not correctly predicted. This effect is becoming more important for modern, tall wind turbines particularly when the atmosphere is ‘non standard’ (i.e. diverging from neutrality).

* In chapter V a second effect of atmospheric stability is investigated. Not only has the sound level been underestimated, but also the effect on the sound character: when the atmosphere turns stable, a more pronounced beating sound evolves. Most of the data are from the Rhede wind farm, complemented by data from a smaller single turbine elsewhere and theoretical calculations. In a section on the perception of fluctuating sound, it is explained how an apparently weak sound level variation can indeed turn into audibly pronounced beating. This chapter corresponds to a published paper [Van den Berg 2005a], but the section on interaction of several turbines (V.2.4) has been
combined with the corresponding section of the first paper [Van den Berg 2004a]. In this chapter the fact that wind velocity in the rotor is not equal to the free wind velocity, which was neglected in the paper, has been taken into account.

♦ In chapter VI data on atmospheric stability and wind statistics are presented. The raw data are from a location in the mid west of the Netherlands and have been provided by the KNMI. The analysis and application to a reference wind turbine help us to understand the behaviour of wind turbines and, together with research results from other countries, show that the atmospheric conditions found at the Rhede wind farm certainly were no exception. This chapter is the text of a paper presented at the WindTurbineNoise2005 conference [Van den Berg 2005b], with some results from other presentations at that conference added (in section VI.6).

♦ In chapter VII some possibilities are discussed to cope with the effects of atmospheric stability on wind turbine noise, either by controlling wind turbine performance or by new designs. In part this is derived from a project in the town of Houten where the town council wants to permit a wind farm, taking into account the effect on residents, especially at night. This chapter is a somewhat expanded version (a concluding section has been added) of a second paper presented at the WindTurbineNoise2005 conference [Van den Berg 2005c].

♦ In chapter VIII a new topic is introduced: how does wind affect sound from a microphone? It shows that atmospheric turbulence, closely related to -again- atmospheric stability, is the main cause of wind induced microphone noise. The chapter corresponds to a published article [Van den Berg 2006].

♦ In Chapter IX all results are summarized. Based on these general conclusions recommendations are given for a fresh look at wind turbine noise.

♦ Finally, in chapter X, some thoughts are given to conclude the text. After that the appendices give additional information.
II ACOUSTICAL PRACTICE
AND SOUND RESEARCH

II.1. Different points of view

In 2001 the German wind farm Rhede was put into operation close to the Dutch border. Local authorities as well as residents at the Dutch side had opposed the construction of the 17 wind turbines because of the effects on landscape and environment: with 98 m hub height the 1.8 MW turbines would dominate the skyline of the early 20th century village of Bellingwolde and introduce noise in the quiet area.

With the turbines in operation, residents at 500 m and more from the wind farm found the noise (and intermittent or flicker shadow, which will not be dealt with here) worse than they had expected. The wind farm operator declined to take measures as acoustic reports showed that German as well as Dutch noise limits were not exceeded. When the residents brought the case to a German court, they failed on procedural grounds. For a Dutch court they had to produce arguments that could only be provided by experts.

Science Shops are specifically intended to help non-profit groups by doing research on their behalf. For the Science Shop for Physics in Groningen noise problems constitute the majority of problems that citizens, as a group or individually, come up with. Although the aim of our research is the same as for acoustic consultants – to quantify sound levels relevant for annoyance – the customers are different: consultants mostly work for the party responsible for the sound production, whereas the Science Shop mostly works for the party that is affected by the sound. This may lead to different research questions. In the case of wind farm Rhede a consultancy will check the sound production of the turbines and check compliance of the calculated sound immission level with relevant limits. However, the Science Shop, taking the strong reaction from the residents as a starting point, wanted to check whether the real sound immission agrees with the
calculated one and whether sound character could explain extra annoyance.

In the Dutch professional journal ‘Geluid’ it was shown, on the basis of 30 acoustic reports, that acoustic consultants tend to rely too much on information from their customers, even when they had reason to be critical about it [Van den Berg 2000]. As consultants’ customers are usually noise producers and authorities, the point of view of those that are affected by noise is not usually very prominent. This book shows that for wind turbines a similar case can be made.

II.2 Results from our wind turbine research

The results of the investigation of the sound from the wind farm Rhede are given in the next chapters. Here the results will be dealt with briefly. The main cause for the high sound level perceived by residents is the fact that wind velocities at night can, at 100 m height, be substantially higher than expected. As a consequence a wind turbine produces more sound. As measured immission levels near the wind farm Rhede show, the discrepancy may be very large: sound levels are up to 15 dB higher than expected at 400 m from the wind farm. The important point is not so much that the maximum measured sound level is higher than the maximum expected sound level (it was, around +2 dB, but this was not an effect of the wind velocity profile). The point is that this maximum does not only occur at high wind velocities as expected, accompanied by high wind induced ambient sound levels, but already at relatively low wind velocities (4 m/s at 10 m height) when there is little wind at the surface and therefore little wind induced background sound. Thus, the discrepancy of 15 dB occurs at quiet nights, but yet with wind turbines at almost maximum power. This situation occurs quite frequently.

A second effect that adds to the sound annoyance is that the sound has an impulsive character. The primary factor for this appeared to be the well known swishing sound one hears close to a turbine. For a single turbine these 1 – 2 dB broad band sound pressure fluctuations would not classify as impulsive, but at night this swish seems to evolve into a less gentle thumping. Also, when several turbines operate nearly synchronously the
pulses may occur in phase increasing pulse strength further. At some distance from the wind farm this sound characteristic, described as thumping or beating, can be very pronounced though in the wind farm, close to a turbine, we never heard this impulsiveness. Indeed, close to a turbine it seems that most sound is coming from the downgoing blade, not when it passes the tower. One has to be careful in estimating blade position, as an observer at, say, 100 m from the foot of the tower is 140 m from a 100 m hub and therefore hears the sound from a blade approximately half a second after it was produced, in which time a blade may have rotated over some 30°. At the Berlin WindTurbineNoise conference Oerlemans [2005] explained this phenomenon: when the blade comes down and heads towards the observer, the observer is at an angle to the blade where most sound is radiated (see remark on directivity just below equation B.5 in Appendix B). On top of that the high tip velocity (70 m/s) causes a Doppler amplification. Both effects increase the sound level for our observer. However, this observation cannot be used for a distant turbine as in that case the observer sees the rotor sideways. Then the change due to the directivity of the sound is small, and also the Doppler effect is nil as the change in the velocity component towards the observer is negligible.

II.3 Early warnings of noisy wind turbines?

One may wonder why the strong effect of the nightly wind profile or the thumping was not noticed before. In the 1998 publication IEC 16400 only the neutral logarithmic wind profile is used [IEC, 1998]. As recent as 2002 it was stated that wind turbine sound is not impulsive [Kerkers et al 2002], which was concluded from assumed, not from measured sound level variations.

There have been some warnings, though. In 1998 Rudolphi concluded from measurements that wind velocity at 10 m height is not a good measure for the sound level: at night the (58 m hub height) turbine sound level was 5 dB higher than expected [Rudolphi 1998]. This conclusion was not followed by more thorough investigation. Since several years residential groups in the Netherlands and abroad complained about
annoying turbine sound at distances where they are not even expected to be able to hear the sound. Recently Pedersen et al [2003, 2004] found that annoyance was relatively high at calculated maximum sound immission levels below 40 dB(A) where one would not expect strong annoyance. As wind turbines become taller, the discrepancy between real and expected levels grows and as more tall wind turbines are constructed complaints may become more widespread. In the Netherlands residents near the German border were the first Dutch to be acquainted with turbines of 100 m hub heights.

It may be that earlier discrepancies between real and projected sound immission were not sufficient to evoke strong community reactions and that only recently turbines have become so tall that the discrepancy now is intolerable. There are other reasons that early warnings perhaps did not make much impression. One is that sound emission measurements are usually done in daytime. It is hard to imagine the sound would be very different at night time, so (almost) no one did. Until some years ago, I myself could not imagine how people could hear wind turbines 2 km away when at 300 to 400 m distance the (calculated) immission level was, for a given wind velocity, already equal to the ambient background sound level (L_{eq}). But it proved I had not listened in a relevant period: an atmospherically stable night.

What is probably also a reason is the rather common attitude that ‘there are always people complaining’. Complaints are a normal feature, not as such a reason to re-investigate. Indeed Dutch noise policy is not to prevent any noise annoyance, but to limit it to acceptable proportions. Added to this is a rather general conviction of Dutch authorities and consultants that routine noise assessment in compliance with legal standards must yield correct results. If measurements are performed it is to check actual emission levels - usually in normal working hours, so in daytime. It is quite unusual to compare the calculated sound immission from a wind turbine (farm) with measured immission levels (so unusual that it is likely that we were the first to do so).
A third reason may be partiality to the outcome of the results. Wind turbine operators are not keen on spending money that may show that sound levels do not comply with legal standards. And if, as expected, they do comply, the money is effectively wasted. Apart from this, we have the experience that at least some organisations that advocate wind energy are not interested in finding out why residents oppose wind farms.

II.4 The use of standard procedures

Although our objective was to measure immission sound levels, we also wanted to understand what was going on: if levels were higher than expected, was that because emission was higher or attenuation less? Could there be focussing or interference? We therefore also measured sound emission as a function of rotational speed of the variable speed turbines. An interesting point that came up with the emission measurement was that compliance with the recommended standard [Ljunggren 1997 or IEC 1998] was impossible. As the farm operator withdrew the co-operation that was previously agreed upon, we had to measure emission levels with the full wind farm in operation, as we obviously did not have the means to stop all turbines except the one to be measured, as the standard prescribes. To measure ambient background sound level, even the last turbine should be stopped.

According to the recommended standard the sound emission should be measured within 20% of the distance to the turbine equal to hub height + blade length. However, to prevent interference from the sound from other turbines the measurement location had to be chosen closer to the turbine. The primary check on the correctness of the distance (i.e. not too close to other turbines) was by listening: the closest turbine should be the dominant source. If not, no measurement was done, and usually a measurement near another turbine was possible. Afterwards we were able to perform a second check by comparing the measured sound immission of the wind farm at a distance of 400 m with the level calculated with a sound propagation model with the measured emission level of all (identical) turbines as input. The calculated difference between a single turbine sound power level and the immission level was 58.0 dB (assuming a constant spectrum this is independent from the power level itself). The measured average difference
was 57.9 dB, with a maximum deviation of individual measurement points of 1.0 dB. So our measurements proved to be quite accurate, deviating only 0.1 ± 1.0 dB from the expected value! In fact, from our measurements one may conclude that, to determine turbine sound power level, it is easier and cheaper to determine total sound emission by measurements at some distance from a wind farm than measuring separate turbines. The wind induced ambient sound, that easily spoils daytime measurements, is not an important disturbance in many nights!

Using a 1 m diameter round hard board, again to comply with the standard, was quite impractical and sometimes impossible. E.g. at one place potato plants would have to be cleared away, at another place one would have to create a flat area in clumps of grass in a nature reserve, both unnecessarily. Instead of the large board we used the side (30-44 cm²) of a plastic sound meter case. We convinced ourselves that (in this case) this was still a good procedure by comparing at one location sound levels measured on the case on soft ground with sound levels measured on a smooth tarmac road surface a few meters away, both at the same distance to the turbine as in the other measurements: there was no difference.

Whether a turbine produces impulsive sound is usually determined by listening to and measuring the sound near a single turbine (along with measurements to determine sound power and spectral distribution). In the Netherlands impulsivity is judged subjectively (by ear), not by a technical procedure as in Germany, though judgement can be supported with a sound registration showing the pulses. Interestingly, in Dutch practice only an acoustician’s ear seems reliable, though even their opinions may disagree. From our measurements the impulsive character can be explained by the wind profile and the interaction of the sound of several turbines. Even at a time the impulsive character can be heard near residents’ dwellings, it cannot clearly be heard close to the turbines in the wind farm (as explained in section II.2). So here also there was need to do measurements where people are actually annoyed, and not to rely on source measurements only, certainly not from a single turbine.
When noise disputes are brought to court, it is clearly advantageous to have objective procedures and standards to assure that the technical quality, which can hardly be judged by non-experts, is sufficient and therefore the results are reliable. In the case made here however, a standard may be non-applicable for valid reasons. Nonetheless, the emission measurements have been contested on procedural grounds (viz. we have not complied to the standard [Kerkers 2003]), even though the immission sound levels were the primary research targets and we did not really need the sound emission measurement results (which, however, proved very accurate).

The tendency to put all noise assessment into technical standard procedures has the disadvantage that when there is a flaw in a legally enforced standard, still the standard is followed, not reality. It is hardly possible for non-experts, such as residents, to bring other arguments to court. They, the annoyed, will have to hire an expert to objectify their annoyance. This is not something every citizen can afford.

**II.5 Modelling versus measurements**

Being able to calculate sound levels from physical models is a huge advantage over having to do measurements (if that, indeed, is possible) especially as in practical situations conditions keep changing and other sounds disturb the measurements. Because of its obvious advantages models have become far more important for noise assessment than measurements. In the Netherlands usually sound emission measurements are carried out close to a source to determine sound power levels. Then, with the sound power level, the immission level is calculated, usually on façades of residences close to the sound source. It is not common to measure immission levels in the Netherlands; in some cases (e.g. railway, aircraft noise) there is not even a measurement method (legally) available to check calculated levels.

However, a physical model is never the same as reality. As will be shown in this book, the widely used standard to quantify sound emission from wind turbines is implicitly based on a specific wind profile. This profile is
not correct at night, although the night is the critical period for wind turbine noise assessment.

Even a perfect physical model will not reproduce reality if input values are not according to reality. An example is to apply sound power levels from new sources (cars, road surfaces, aeroplanes, mopeds, vacuum cleaners, etc.), maybe acquired in a specific test environment, to real life situations and conditions. Another example is a wind farm south of the Rhede wind farm where a turbine produced a clearly audible and measurable tonal sound, probably caused by damage on a blade. It is very hard for residents to convince the operator and authorities of this annoying fact, partly because most experts say that modern wind turbines do not produce tonal sound.

Incorrect models and incorrect input may well occur together and be difficult to separate. It is important that calculation models are checked for correctness when they are used in new applications. Situations where (strong) complaints arise may indicate just those cases where models do not cover reality.

II.6 Conclusion

In modelling wind turbine sound very relevant atmospheric behaviour has been 'overlooked'. As a consequence, at low surface wind velocities such as often occur at night, wind turbine noise immission levels may be much higher than expected. The discrepancy between real and modelled noise levels is greater for tall wind turbines. International models used to assess wind turbine noise on dwellings should be revised for this atmospheric effect, at least by giving less attention to the 'standard' neutral atmosphere.

A discrepancy between noise forecasts and real noise perception, as a result of limited or even defective models, cannot always be avoided, even not in principle. However, its consequences can be minimised if immission levels are measured at relevant times and places. This relevancy is also determined by observations of those affected. It should always be possible to check noise forecasts by measurement.
For wind turbine noise (and other noise sources) standard measurement procedures require co-operation of the operator to be able to check emission sound levels. This introduces an element of partiality to the advantage of the noise producer. This is also generally a weak point in noise assessment: the source of information is usually the noise producer. Hence there should always be a procedure to determine noise exposure independently of the noise producer.

Standard technical procedures have the benefit of providing quality assurance: when research has been conducted in compliance with a standard procedure lay persons should be able to rely on the results. It may however also have a distinct disadvantage for lay people opposing a noise source: when an assessment does not comply with a standard procedure it is not accepted in court, regardless of the content of the claim. A consequence is they have to depend on legal as well as acoustical expertise. If citizens are forced to use expert knowledge, one may argue that they should be given access to that knowledge. An important obstacle is the cost of that access.
III BASIC FACTS: wind power and the origins of modern wind turbine sound

III.1 Wind energy in the EU

Modern onshore wind turbines have peak electric power outputs up to 3 MW and tower heights of 80 to 100 meters. In 2003, 75% of the global wind power peak electric output of 40 GW was installed in the European Union. The original European target for 2010 was 40 GW, but the European Wind Energy Association have already set a new target for 2010 of 75 GW, of which 10 GW is projected off-shore, while others have forecasted a peak output of 120 GW for that year [EWEA 2004]. Whether this growth will actually occur is uncertain; with the proportional increase of wind energy in total electric power the difficulties and costs of integrating large scale windpower with respect to grid capacity and stability, reserve capacity and CO₂ emission reductions are becoming more prominent [see, e.g., E.On 2004, ESB 2004]). However, further expansion of wind energy is to be expected, and as a result of this (predominantly on-shore) growth an increasing number of people may face the prospect of living near wind farms, and have reason to inquire and perhaps be worried about their environmental impact. Visual intrusion, intermittent reflections on the turbine blades, as well as intermittent shadows (caused when the rotating blades pass between the viewer and the sun), and sound, are usually considered potentially negative impacts.

III.2 Wind profiles and atmospheric stability

Atmospheric stability has a profound effect on the vertical wind profile and on atmospheric turbulence strength. Stability is determined by the net heat flux to the ground, which is a sum of incoming solar and outgoing thermal radiation, and of latent and sensible heat exchanged with the air and the subsoil. When incoming radiation dominates (clear summer days) air is heated from below and rises: the atmosphere is unstable. Thus, thermal turbulence implies vertical air movements, preventing large
variations in the vertical wind velocity gradient (i.e. the change in time averaged wind velocity with height). When outgoing radiation dominates (clear nights) air is cooled from below; air density will increase closer to the ground, leading to a stable configuration where vertical movements are damped. The ‘decoupling’ of horizontal layers of air allows a higher vertical wind velocity gradient. A neutral state occurs when thermal effects are less significant, which is under heavy clouding and/or in strong winds.

Wind velocity at altitude $h_2$ can be deduced from wind velocity at altitude $h_1$ with a simple power law function:

$$\frac{V_{h_2}}{V_{h_1}} = (\frac{h_2}{h_1})^m$$

Equation III.1 is an engineering formula used to express the degree of stability in a single number (the shear exponent $m$), but has no physical basis. The relation is suitable where $h$ is at least several times the roughness height (a height related to the height of vegetation or obstacles on the ground). Also, at high altitudes the wind profile will not follow (III.1), as eventually a more or less constant wind velocity (the geostrophic wind) will be attained. At higher altitudes in a stable atmosphere there may be a decrease in wind velocity when a nocturnal ‘jet’ develops. The maximum in this jet is caused by a transfer of kinetic energy from the near-ground air that decouples from higher air masses as large, thermally induced eddies vanish because of ground cooling. In fact, reversal of the usual near-ground diurnal pattern of low wind velocities at night and higher wind velocities in daytime is a common phenomenon at higher altitudes over land in clear nights as will be shown further below (Chapter VI). Over large bodies of water the phenomenon may be seasonal as atmospheric stability occurs more often when the water is relatively cold (winter, spring). This may also be accompanied by a maximum in wind velocity at a higher altitude [Smedman et al 1996].

In flat terrain the shear exponent $m$ has a value of 0.1 and more. For a neutral atmosphere $m$ has a value of approximately 1/7. In an unstable atmosphere - occurring in daytime - thermal effects caused by ground heating are dominant. Then $m$ has a lower value, down to approximately
0.1. In a stable atmosphere vertical movements are damped because of
ground cooling and \( m \) has a higher value. One would eventually expect a
parabolic wind profile, as is found in laminar flow, corresponding to a
value of \( m \) of \( 0.7 = \sqrt{\frac{1}{2}} \). Our measurements near the Rhede wind farm
yielded values of \( m \) up to 0.6. A sample (averages over 0:00–0:30 GMT of
each first night of the month in 1973) from data from a 200 m high tower
in flat, agricultural land [Van Ulden et al 1976] shows that the theoretical
value is indeed reached: in ten out of the twelve samples there was a
temperature inversion in the lower 120 m, indicating atmospheric stability.
In six samples the temperature increased with more than 1 °C from 10 to
120 m height and the exponent \( m \) (calculated from (III.1): \( m = \log(V_{80}/V_{10})/\log(8) \)) was 0.43, 0.44, 0.55, 0.58, 0.67 and 0.72. More data
from this site (Cabauw) and other areas will be presented in chapter VI.

A physical model to calculate wind velocity \( V_h \) at height \( h \) is ([Garrat
1992], p. 53):

\[
V_h = (u*/\kappa)[\ln(h/z_o) - \Psi]
\]  

(III.2)

where \( \kappa = 0.4 \) is von Karman’s constant, \( z_o \) is roughness height and \( u* \) is
friction velocity, defined by \( u^2 = \sqrt{\langle uw^2 + vw^2 \rangle} = \tau/\rho \), where \( \tau \) equals
the momentum flux due to turbulent friction across a horizontal plane, \( \rho \) is
air density and \( u, v \) and \( w \) are the time-varying components of in-wind,
cross-wind and vertical wind velocity, with \( \langle x \rangle \) the time average of \( x \). The
stability function \( \Psi = \Psi(\zeta) \) (with \( \zeta = h/L \)) corrects for atmospheric stability.
Here Monin-Obukhov length \( L \) is an important length scale for stability
and can be thought of as the height above which thermal turbulence
dominate over friction turbulence; the atmosphere at heights \( 0 < h < L \) (if
\( L \) is positive and not very large) is the stable boundary layer. The following
approximations for \( \Psi \), mentioned in many text books on atmospheric
physics (e.g. [Garrat 1992]), are used:

- in a stable atmosphere (\( L > 0 \)) \( \Psi(\zeta) = -5\zeta < 0 \).
- in a neutral atmosphere (\( |L| \) large \( \rightarrow 1/L \approx 0 \)) \( \Psi(0) = 0 \).
- in an unstable atmosphere (\( L < 0 \)) \( \Psi(\zeta) = 2\cdot\ln[(1+x)/2] + \ln[(1+x^2)/2] \)
  \( - 2/\tan(x) + \pi/2 > 0 \), where \( x = (1-16\cdot\zeta)^{1/4} \).
For $\Psi = 0$ equation (III.2) reduces to $V_{h,\text{log}} = (u*/\kappa) \cdot \ln(h/z_o)$, the widely used logarithmic wind profile. With this profile the ratio of wind velocities at two heights can be written as:

$$V_{h_2,\text{log}}/V_{h_1} = \log(h_2/z_o)/\log(h_1/z_o)$$

(III.3)

For a roughness length of $z_o = 2$ cm (pasture) and $m = 0.14$, the wind profiles according to equations III.1 and III.3 coincide within 2% for $h < 100$ m. In figure III.1 wind profiles are given as measured by Holtslag [1984], as well as wind profiles according to formulae (III.1) and (III.3).

Formula III.3 is an approximation of the wind profile in the turbulent boundary layer of a neutral atmosphere, when the air is mixed by turbulence resulting from friction with the surface of the earth. In daytime thermal turbulence is added, especially when there is strong insolation. At night time a neutral atmosphere, characterized by the adiabatic temperature gradient of $-1$ °C per 100 m, occurs under heavy clouding and/or at relatively high wind velocities. When there is some clear sky and in the absence of strong winds the atmosphere becomes stable because of radiative cooling of the surface: the wind profile changes and can no longer be adequately described by (III.3). The effect of the change to a stable atmosphere is that, relative to a given wind velocity at 10 m height in daytime, at night there is a higher wind velocity at hub height and thus a higher turbine sound power level; also there is a lower wind velocity below 10 m and thus less wind-induced sound in vegetation.
With regard to wind power some attention is being paid to stability effects and thus to other wind profile models such as the diabatic wind velocity model (III.2) [see, e.g., Archer et al 2003, Baidya Roy et al 2004, Pérez et al 2004, Smedman et al 1996, Smith et al 2002]. In relation to wind turbine sound, much less attention has been given to atmospheric stability (see section II.3).

Stability can also be categorized in Pasquill classes that depend on observations of wind velocity and cloud cover (see, e.g., [LLNL 2004]). They are usually referred to as classes A (very unstable) through F (very stable). In a German guideline [TA-Luft 1986] a closely related classification is given (again closely related to the international Turner classification [Kühner 1998]). An overview of stability classes with the appropriate value of m is given in table III.1.

<table>
<thead>
<tr>
<th>Pasquill class</th>
<th>name</th>
<th>comparable stability class [TA-Luft 1986]</th>
<th>m</th>
</tr>
</thead>
<tbody>
<tr>
<td>A</td>
<td>very unstable</td>
<td>V</td>
<td>0.09</td>
</tr>
<tr>
<td>B</td>
<td>moderately unstable</td>
<td>IV</td>
<td>0.20</td>
</tr>
<tr>
<td>C</td>
<td>neutral</td>
<td>IV2</td>
<td>0.22</td>
</tr>
<tr>
<td>D</td>
<td>slightly stable</td>
<td>IV1</td>
<td>0.28</td>
</tr>
<tr>
<td>E</td>
<td>moderately stable</td>
<td>II</td>
<td>0.37</td>
</tr>
<tr>
<td>F</td>
<td>(very) stable</td>
<td>I</td>
<td>0.41</td>
</tr>
</tbody>
</table>

According to long-term data from Eelde and Leeuwarden [KNMI 1972], two meteorological measurement sites of the KNMI (Royal Netherlands Meteorological Institute) in the northern part of the Netherlands, a stable atmosphere (Pasquill classes E and F) at night occurs for a considerable proportion of night time: 34% and 32% respectively.

From formula (III.3) the ratio of wind velocities at hub height (98 m) and reference height, over land with low vegetation ($z_o = 3$ cm), is $f_{log} = V_{98}/V_{10} = 1.4$. According to formula (III.1) and table III.1 this ratio would
be $f_{\text{unstable}} = 1.2 = 0.85 f_{\log}$ in a very unstable atmosphere and $f_{\text{stable}} = 2.5 = 1.8 f_{\log}$ in a (very) stable atmosphere.

The shear exponent $m$ can be determined from the measured ratio of wind velocities at two heights ($V_{h2}/V_{h1}$) using equation III.1:

$$m_{h1,h2} = \ln(V_{h2}/V_{h1})/\ln(h_2/h_1)$$  \hspace{1cm} (III.4)

### III.3 Air flow on the blade

As is the case for aircraft wings, the air flow around a wind turbine blade generates lift. An air foil performs best when lift is maximised and drag (flow resistance) is minimised. Both are determined by the angle of attack: the angle ($\alpha$) between the incoming flow and the blade chord (line between front and rear edge; see figure III.2). The optimum angle of attack for turbine blades is usually between 0 and 4°, depending on the blade profile.

![Figure III.2: flow impinging on a turbine blade with flow angle $\varphi$, blade pitch angle $\theta$ and angle of attack on blade $\alpha = \varphi - \theta$](image)

The local wind at the blade is not the unobstructed wind velocity. The rotor extracts energy from the air at the cost of the kinetic energy of the wind. The velocity of the air passing through the rotor is thus reduced to $V_b = (1 - a)V_h$, where $a$ is the induction factor. The highest efficiency of a wind turbine is reached at the Betz limit: at this theoretical limit the induction factor is 1/3 and the efficiency is 16/27 (≈ 60%) [Hansen 2000]. The wind velocity at the blade is thus:

$$V_b = V_h \cdot 2/3$$  \hspace{1cm} (III.5)
**III.4 Main sources of wind turbine sound**

There are many publications on the nature and power of turbine sound: original studies [e.g. Lowson 1985, Grosveld 1985] and reviews [e.g. Hubbard et al 2004, Wagner et al 1996]. A short introduction on wind aeroacoustics will be given to elucidate the most important sound producing mechanisms. If an air flow is smooth around a (streamlined) body, it will generate very little sound. For high velocities and/or over longer lengths the flow in the boundary layer between the body and the main flow becomes turbulent. The rapid turbulent velocity changes at the surface cause sound with frequencies related to the rate of the velocity changes. The turbulent boundary layer at the downstream end of an airfoil produces *trailing edge sound*, which is the dominant audible sound from modern turbines. When the angle of attack increases from its optimal value the turbulent boundary layer on the suction (low pressure) side grows in thickness, thereby decreasing power performance and increasing sound level. For high angles of attack this eventually leads to stall, that is: a dramatic increase of drag on the blades. Apart from this turbulence inherent to an airfoil, the atmosphere itself is turbulent over a wide range of frequencies and sizes.

![Figure III.3: 15 m blades for Altamont Pass, Ca (photo: Alex Haag)](image)

Turbulence can be defined as changes over time and space in wind velocity and direction, resulting in velocity components normal to the airfoil varying with the turbulence frequency causing *in-flow turbulent sound*. Atmospheric turbulence energy has a maximum at a frequency that depends on altitude and on atmospheric stability. For wind turbine altitudes...
this peak frequency is of an order of magnitude of once per minute (0.017 Hz). The associated eddy (whirl) scale is of the order of magnitude of several hundreds of meters [Petersen et al 1998] in an unstable atmosphere, less in a stable atmosphere. Eddy size and turbulence strength decrease at higher frequency, and vanish due to viscous friction when the eddies have reached a size of approximately one millimetre.¹

A third sound producing mechanism is the response of the blade to the change in lift when it passes the tower. The wind is slowed down by the tower which changes the angle of attack on the blade; as a result the lift and drag forces on the blade suddenly change. The resulting sideways movement of the blade causes thickness sound at the blade passing frequency and its harmonics.² Thickness sound is also mentioned as sound originating from the (free) rotating blade pushing the air sideways. However, the associated air movement is relatively smooth and is not a relevant source of sound.

A more thorough review of these three sound production mechanisms is given in appendix B, where frequency ranges and sound levels are quantified in so far as relevant for this book.

Sound originating from the generator or the transmission gear has decreased in level in the past decades and has become all but irrelevant if considering annoyance for residents.

To summarize, a modern wind turbine sound spectrum can be divided in (overlapping) regions corresponding to the three mechanisms mentioned:

¹ for more information on atmospheric turbulence: see chapter VIII
² a thickness sound pulse has a length \( t_{\text{pulse}} \) with an order of magnitude of (tower diameter/tip speed \( \approx \)) 0.1 s, so its spectrum has a maximum at \( 1/t_{\text{pulse}} \approx 10 \text{ Hz} \). The spectrum of a periodic series of Dirac pulses (unit energy 'spikes' with, here, a period of \( T_{\text{blade}} \)) is a series of spikes at frequencies \( n/T_{\text{blade}} \) (\( n = 1, 2, 3, 4, \ldots \)). When periodic thickness sound is considered as a convolution of the single sound pulse with a series of Dirac pulses, the Fourier transform is the product of the transforms of both, that is: the product of the sound pulse spectrum centered at \( 1/t_{\text{pulse}} \) and spikes at \( n/T_{\text{blade}} \). The result is a series of spikes with the single sound pulse spectrum as an envelope, determining each spike level. In practice \( 1/T_{\text{pulse}} \) usually has a value of 4 to 8 Hz (see e.g. [Wagner 1996]) and the harmonic closest to this frequency carries most energy.
- High frequency: trailing edge (TE) sound is noise with a maximum level at 500–1000 Hz for the central octave band, decreasing with 11 dB for neighbouring octave bands and more for further octave bands.
- Low frequency: in-flow turbulent sound is broad band noise with a maximum level of approximately 10 Hz and a slope of 3–6 dB per octave.
- Infrasound frequency ($f < 30$ Hz): the thickness sound is tonal, the spectrum containing peaks at the blade passing frequency $f_B$ and its harmonics.

As thickness sound is not relevant for direct perception, turbulent flow is the dominant cause of (audible) sound for modern wind turbines. It is broad band noise with no tonal components and only a little variation, known as blade swish. Trailing edge sound level is proportional to $50\cdot\log M$ (see equation B.4 in appendix B), where $M$ is the Mach number of the air impinging on the blade. TE sound level, the dominant audible sound source in a modern turbine, therefore increases steeply with blade speed and is highest at the high velocity blade tips. Writing Mach number at the blade tip as $M = V_{tip}/c$, wind turbine sound level strongly depends on blade tip speed $V_{tip}$:

$$L_{TE} \sim 50\cdot\log(V_{tip}/c)$$  \hspace{1cm} (III.6)
Swish, which is the variation in TE sound, thus also originates predominantly at the tips.

This book deals with modern variable speed turbines where the angle of attack is constant over a wide range of wind speeds. Keeping blade pitch (the angle between the blade chord and the rotor plane) constant, the rotational speed increases with wind speed usually up to a rated wind speed of some 14 m/s. At higher wind speeds the pitch angle is decreased at constant rotational speed to keep a constant angle of attack until for safety reasons the rotor is stopped. The effect on sound production is that first the sound power level increases up to the rated wind speed, then remains almost constant at higher wind speeds.

In a constant speed turbine the rotational speed has a fixed value, though usually a turbine then has two speeds to accommodate for low and high wind speeds. Here the blade pitch is set to optimize the angle of attack up to the rated power. Above rated power, a situation that will not occur very often, the pitch angle is kept constant, so the angle of attack increases with wind speed and the turbine becomes less efficient. The result is that the sound power at low speed is almost constant, then increases sharply at the change to the higher speed. After that it is again almost constant, increasing again above the rated power when the angle of attack drifts away from the optimum value.

Sound from downwind rotors, i.e. with the rotor downwind from the tower, was considered problematic as it was perceived as a pulsating sound (see appendix B). For modern upwind rotors this variation in sound level is weaker. It is not thought to be relevant for annoyance and considered to become less pronounced with increasing distance due to loss of the effect of directivity, due to relatively high absorption at swish frequencies, and because of the increased masking effect of background noise [ETSU 1996]. However, an increase in the level of the swishing sound related to increasing atmospheric stability has not been taken into account as yet. In this context the periodic change in angle of attack near the tower proves to be important, not in relation to thickness sound but as a modulation period.
So, what's the sound like...?

(.....) Our experience is that mechanical noise is insignificant compared to the aerodynamic noise, or 'blade thump' as we call it. At "our" windfarm the mechanical noise is usually only audible when within about 100 metres of the turbine, but the blade thump can be heard at distances of up to 1.5 Km away.

(.....)

Some residents describe this noise as an old boot in a tumble dryer, others as a Whumph! Whumph! Whumph! Either way its not particularly loud at 1.5 km distance but closer than that and it can be extremely irritating when exposed to it for any period of time. Some residents have even resorted to stuffing chimney stacks with newspaper as the sound reverberates down the stack.

Because it is generally rhythmic, it's not the kind of noise that you can shut out of your mind, like, say, distant road noise - this is why we think the noise level stipulation on the planning conditions of such a windfarm development is woefully inadequate for protecting local residents from the noise effects of a windfarm.

All of us agree that the most disturbing aspect of the noise is the beat that we think is caused by the blades passing the tower of the turbine. As the rotational speed of the 3 bladed turbines is about 28 rpm "on full song" this results in a sound of about 84 beats per minute from each turbine.

The sound rises and falls in volume due to slight changes in wind direction but the end result for those in the affected area is a feeling of anxiety, and sometimes nausea, as the rate continually speeds and slows - we think that is maybe because this frequency of the pulses is close to the human heart rate and some residents feel that their own pulse rate is trying to match that of the turbines. (.....)

When does it strike?

The windfarm makes a noise all the time it is operating, however there are times when it becomes less of a nuisance.

When the wind is very strong, the background noise created by the wind whistling around trees etc. drowns out the noise of the turbines and the problem is reduced. (.....)

In this area we all agree that the worst conditions are when the wind is blowing lightly and the background noise is minimal. Under these conditions residents up to 1 kilometre have complained to the Environmental Health department about the drone from the turbines. Unfortunately these are just the sort of weather conditions that you would wish to be outside enjoying your garden. (.....)

During the summer nights it is not possible for some residents, even as far away as 1000 metres, to sleep with the window open due to the blade thump. (.....)

Excerpts describing wind turbine sound and its effects, from a page of the website of MAIWAG (consulted December 3, 2005), a group of residents in three villages in the south of Cumbria (UK)
IV LOUD SOUNDS IN WEAK WINDS: effect of the wind profile on turbine sound level

IV.1 The Rhede wind farm

In Germany several wind turbine farms have been and are being established in sparsely populated areas near the Dutch border. One of these is the Rhede wind farm in northwestern Germany (53° 6.2' latitude, 7° 12.6' longitude) with seventeen Enercon E-66 1.8 MW turbines of 98 m hub height and with 3-blade propellers of 35 m blade length. The turbines have a variable speed increasing with wind velocity, starting with 10 rpm (revolutions per minute) at a wind velocity of 2.5 m/s at hub height up to 22 rpm at wind velocities of 12 m/s and over.

At the Dutch side of the border is a residential area along the Oude Laan and Veendijk in De Lethe (see figure IV.2): countryside dwellings surrounded by trees and agricultural fields. The dwelling nearest to the wind farm is some 500 m west of the nearest wind turbine (nr. 16). According to a German noise assessment study a maximum immission level of 43 dB(A) was expected, 2 dB below the relevant German noise limit. According to a Dutch consultancy immission levels would comply with Dutch (wind velocity dependent) noise limits.

After the farm was put into operation residents made complaints about the noise, especially at (late) evening and night. The residents, united in a neighbourhood group, could not persuade the German operator into mitigation measures or an investigation of the noise problem and brought the case to court. The Science Shop for Physics had just released a report explaining a possible discrepancy between calculated and real sound immission levels of wind turbines because of changes in wind profile, and was asked to investigate the consequences of this discrepancy by sound measurements. Although at first the operator agreed to supply measurement data from the wind turbines (such as power output, rotation speed, axle direction), this was withdrawn after the measurements had started. All relevant data therefore had to be supplied or deduced from our own measurements.
Figure IV.2: turbines (dots W1,...,W17) in and measurement locations (crosses A,...,X) near the Rhede wind farm; Duch – German border indicated by line of +++ (through A); grid lines are 1 km apart, north is at top
IV.2 Noise impact assessment

In the Netherlands and Germany noise impact on dwellings near a wind turbine or wind farm is calculated with a sound propagation model. Wind turbine sound power levels $L_w$ are used as input for the model, based on measured or estimated data. In Germany a single ‘maximum’ sound power level (at 95% of maximum electric power) is used to assess sound impact. In the Netherlands sound power levels related to wind velocities at 10 m height are used; the resulting sound immission levels are compared to wind velocity dependent noise limits (see figure VII.1). Implicitly this assessment is based on measurements in daytime and does not take into account atmospheric conditions affecting the wind profile, especially at night.

In the Netherlands a national calculation model is used [VROM 1999] to assess noise impact, as is the case in Germany [TA-Lärm 1998]. According to Kerkers [Kerkers 1999] there are, at least in the case of these wind turbines, no significant differences between both models. In both sound propagation models the sound immission level $L_{imm}$ at a specific observation point is a summation over $j$ sound power octave band levels $L_{wj}$ of $k$ sources (turbines), reduced with attenuation factors $D_{j,k}$:

$$L_{imm} = 10 \cdot \log \left[ \Sigma_j \Sigma_k 10^{0.1 \cdot (L_{wj} - D_{j,k})} \right]$$

(IV.1)
L_{wj}, assumed identical for all k turbines, is a function of rotational speed. D_j is the attenuation due to geometrical spreading (D_{geo}), air absorption (D_{j-air}) and ground absorption (D_{j-ground}): D_{j,k} = D_{geo,k} + D_{j-air,k} + D_{j-ground,k}. Formula (IV.1) is valid for a downwind situation. For long term assessment purposes a meteorological correction factor is applied to (IV.1) to account for 'average atmospheric conditions'. When comparing calculated and measured sound immission levels in this study no such meteo-correction is applied because measurements were always downwind of a turbine or the wind farm.

**IV.3 Wind turbine noise perception**

There is a distinct audible difference between the night and daytime wind turbine sound at some distance from the turbines. On a summer's day in a moderate or even strong wind the turbines may only be heard within a few hundred meters and one might wonder why residents should complain of the sound produced by the wind farm. However, in quiet nights the wind farm can be heard at distances of up to several kilometers when the turbines rotate at high speed. In these nights, certainly at distances from 500 to 1000 m from the wind farm, one can hear a low pitched thumping sound with a repetition rate of about once a second (coinciding with the frequency of blades passing a turbine mast), not unlike distant pile driving, superimposed on a constant broad band 'noisy' sound. A resident living at 1 km from the nearest turbine says it is the rhythmic character of the sound that attracts attention: beats are clearly audible for some time, then fade away to come back again a little later. A resident living at 2.3 km from the wind farm describes the sound as 'an endless train'. In daytime these pulses are usually not audible and the sound from the wind farm is less intrusive or even inaudible (especially in strong winds because of the then high ambient sound level).

In the wind farm the turbines are audible for most of the (day and night) time, but the thumping is not evident, although a ‘swishing’ sound – a regular variation in sound level- is readily discernible. Sometimes a rumbling sound can be heard, but it is difficult to assign it, by ear, to a specific turbine or to assess it’s direction.
**IV.5 Measurement instruments and method**

Sound immission measurements were made over 1435 hours, of which 417 hours at night, within four months on two consecutive locations with an unmanned Sound and Weather Measurement System (SWMS) consisting of a sound level meter (type 1 accuracy) with a microphone at 4.5 m height fitted with a 9 cm diameter foam wind shield, and a wind meter at 10 m as well as at 2 m height. Every second wind velocity and wind direction (at 10 m and at 2 m height) and the A-weighted sound level were measured; the measured data were stored as statistical distributions over 5 minute intervals. From these distributions all necessary wind data and sound levels can be calculated, such as average wind velocity, median wind direction or equivalent sound level and any percentile (steps of 5%) wind velocity, wind direction or sound level, in intervals of 5 minutes or multiples thereof.

Also complementary measurements were done with logging sound level meters (type 1 and 2 accuracy) and a spectrum analyser (type 1) to measure immission sound levels in the residential area over limited periods, and emission levels near wind turbines. Emission levels were measured according to international standards [IEC 1998, Ljunggren 1997], but for practical purposes they could not be adhered to in detail: with respect to the recommended values a smaller reflecting board was used for the microphone (30-44 cm² instead of a 1 m diameter circular board) and a smaller distance to the turbine (equal to tower height instead of tower height + blade length); reasons for this were given in Chapter II. Also it was not possible to do emission measurements with only one turbine in operation.

**IV.6 Results: sound emission**

Emission levels $L_{eq}$ measured very close to the centre of a horizontal, flat board at a distance $R$ from a turbine hub can be converted to a turbine sound power level $L_w$ [IEC 1998, Ljunggren 1997]:

$$L_w = L_{eq} - 6 + 10 \cdot \log(4\pi R^2/A_o)$$  \hspace{1cm} (IV.2)

where $A_o$ is a unit surface (1 m²). From earlier measurements [Kerkers 1999] a wind velocity dependence of $L_w$ was established as given in table
IV.1. As explained above, the wind velocity at 10 m height was not considered a reliable single measure for the turbine sound power, but rotational speed was a better measure. Emission levels have been measured, typically for 5 minutes per measurement, at nine turbines on seven different days with different wind conditions. The results are plotted in figure IV.3; the sound power level is plotted as a function of rotational speed $N$. $N$ is proportional to wind velocity at hub height and could be determined by counting, typically during one minute, blades passing the turbine mast. This counting procedure is not very accurate (accuracy per measurement is $\leq 2$ counts, corresponding to $2/3$ rpm) and is probably the dominant reason for the spread in figure IV.3. The best logarithmic least squares fit to the data points in figure IV.3 is:

$$L_W = 67.1 \cdot \log(N) + 15.4 \text{ dB(A)}$$

(IV.3)

with a correlation coefficient of 0.98. The standard deviation of measurement values with respect to this fit is 1.0 dB.

<table>
<thead>
<tr>
<th>Table IV.1: sound power level of wind turbines [Kerkers 1999]</th>
</tr>
</thead>
<tbody>
<tr>
<td>wind velocity $V_{10}$</td>
</tr>
<tr>
<td>sound power level $L_W$</td>
</tr>
</tbody>
</table>

<table>
<thead>
<tr>
<th>Table IV.2: octave band spectra of wind turbines at $L_W = 103 \text{ dB(A)}$</th>
</tr>
</thead>
<tbody>
<tr>
<td>frequency</td>
</tr>
<tr>
<td>this report</td>
</tr>
<tr>
<td>[Kerkers 1999]</td>
</tr>
</tbody>
</table>
At the specification extremes of 10 rpm and 22 rpm the (individual) wind turbine sound power level $L_W$ is 82.8 dB(A) and 105.7 dB(A), respectively. In table IV.2 earlier measurement results [Kerkers 1999] are given for the octave band sound power spectrum. Also in table IV.2 the results of this study are given: the logarithmic average of four different spectra at different rotational speeds. In all cases spectra are scaled, with formula IV.3, to the same sound power level of 103 dB(A).

To calculate sound immission levels at a specific rotational speed (or vice versa) the sound power level given in formula (IV.3), and the spectral form in table IV.2 (‘this report’) have been used.

### IV.7 Results: sound immission

The sound immission level has been measured with the unmanned SWMS on two locations. From May 13 until June 22, 2002 it was placed amidst open fields with barren earth and later low vegetation at 400 meters west of the westernmost row of wind turbines (location A, see figure IV.2). This site was a few meters west of the Dutch-German border, visible as a ditch and a 1.5 to 2 m high dike. From June 22 until September 13, 2002 the SWMS was placed on a lawn near a dwelling at 1500 m west of the westernmost row (location B), with low as well as tall trees in the vicinity. On both locations there were no reflections of turbine sound towards the microphone, except via the ground, and no objects (such as trees) in the line of sight between the turbines and the microphone. Apart from possible wind induced sound in vegetation relevant sound sources are traffic on rather quiet roads, agricultural activities, and birds. As, because of the trees, the correct (potential) wind velocity and direction could not be measured on location B, wind measurement data provided by the KNMI were used from their Nieuw Beerta site 10 km to the north. These data fitted well with the measurements on location A.

At times when the wind turbine sound is dominant, the sound level is relatively constant within 5 minute intervals. In figure IV.4 this is demonstrated for two nights. Thus measurement intervals with dominant turbine sound could be selected with a criterion based on a low variation in sound level: $L_5 - L_{95} \leq 4$ dB, where $L_5$ and $L_{95}$ are the 5 and 95 percentile
sound level in the measurement interval. In a normal (Gaussian) distribution this would equal $\sigma \leq 1.2$ dB, with $\sigma$ the standard deviation.

On location A, 400 m from the nearest turbine, the total measurement time was 371 hours. In 25% of this time the wind turbine sound was dominant, predominantly at night (23:00 – 6:00 hours: 72% of all 105 nightly hours) and hardly in daytime (6:00 – 19:00 hours: 4% of 191 hours). See table IV.3.

On location B, 1500 m from the nearest turbine, these percentages are almost halved, but still the turbine sound is dominant for over one third of the time at night (38% of 312 hours). The trend in percentages agree with complaints concerning mostly noise in the (late) evening and at night and their being more strongly expressed by residents closer to the wind farm.
Table IV.3: total measurement time in hours and selected time
with dominant wind turbine sound

<table>
<thead>
<tr>
<th>Location</th>
<th>total time (hours and % of total measurement time at location)</th>
<th>Night 23:00-6:00</th>
<th>Evening 19:00-23:00</th>
<th>Day 6:00-19:00</th>
</tr>
</thead>
<tbody>
<tr>
<td>A: total</td>
<td>371 h</td>
<td>105</td>
<td>75</td>
<td>191</td>
</tr>
<tr>
<td></td>
<td>92 h 25%</td>
<td>76</td>
<td>9</td>
<td>7</td>
</tr>
<tr>
<td>A: selected</td>
<td>136 h 13%</td>
<td>119</td>
<td>13</td>
<td>4</td>
</tr>
<tr>
<td></td>
<td>92 h 25%</td>
<td>76</td>
<td>9</td>
<td>7</td>
</tr>
<tr>
<td></td>
<td>1064 h</td>
<td>312</td>
<td>183</td>
<td>569</td>
</tr>
<tr>
<td>B: total</td>
<td>136 h 13%</td>
<td>119</td>
<td>13</td>
<td>4</td>
</tr>
<tr>
<td></td>
<td>92 h 25%</td>
<td>76</td>
<td>9</td>
<td>7</td>
</tr>
<tr>
<td>B: selected</td>
<td>136 h 13%</td>
<td>119</td>
<td>13</td>
<td>4</td>
</tr>
</tbody>
</table>

In figure IV.5 the selected (L₅-L₉₅ ≤ 4 dB) 5 minute equivalent immission sound levels Leq,5min are plotted as a function of wind direction (left) and of wind velocity (right) at 10 m height, for both location A (above) and B (below). The KNMI wind velocity data (used for location B) were given as integer values of the wind velocity. Also the wind velocity at 10 m and 2 m height on location A are plotted (in IV.5A and IV.5B, respectively), and the local wind velocity (influenced by trees) at 10 m on location B (IV.5C). The immission level data points are separated in two classes where the atmosphere was stable or neutral, according to observations of wind velocity and cloud cover at Eelde. Eelde is the nearest KNMI site for these observations, but it is 40 km to the west, so not all observations will be valid for our area.

In figure IV.5B a grey line is plotted connecting calculated sound levels with sound power levels according to table IV.1 (the lowest value at 2.5 m/s is extrapolated [Van den Berg et al 2002]), implicitly assuming a fixed logarithmic wind profile according to formula (III.2). If this line is compressed in the direction of the abscissa with a factor 2.6, the result is a (black) line coinciding with the maximum one hour values (Lₑq,₁₉h). Apparently for data points on this line the sound emission corresponds to a wind velocity at hub height that is 2.6 times higher than expected. In figure IV.6 this is given for one hour periods: all 5 minute measurement periods.
that satisfied the $L_5$-$L_{95}$-criterion, with at least 4 periods per hour, were taken together in consecutive hourly periods and the resulting $L_{eq,T}$ ($T = 20$ to 60 minutes) was calculated. The resulting 83 $L_{eq,T}$-values are plotted against the average wind velocity $V_{10}$. Also plotted in figure IV.6 are the expected immission levels assuming a logarithmic wind profile calculated from (III.4), with $f_{log} = (V_{98}/V_{10})_{log} = 1.4$ (for $f_{xx}$; see text above equation III.4); the immission levels assuming a stable wind profile with $m = 0.41$, so $f_{stable} = 2.5 = 1.8 \cdot f_{log}$; the maximum immission levels assuming $f_{max} = 3.7$
= 2.6·\text{f}_{\text{log}}, \text{ in agreement with a wind profile (III.2) with } m = 0.57. \text{ The best fit of all data points (L}_{\text{eq,T}} \text{ in figure IV.6 is } L_{\text{eq}} = 32\cdot \log(V_{10}) + 22 \text{ dB (correlation coefficient 0.80) with } 1 < V_{10} < 5.5 \text{ m/s. This agrees within 0.5 dB with the expected level according to the stable wind profile. The best fit of all 5 minute data-points in figure IV.5B yields the same result. Thus on location A the highest one hour averaged hub height wind velocities at night are 2.6 times the expected values according to the logarithmic wind profile in formula (III.4). As a consequence, sound levels at (in night-time) frequently occurring wind velocities of 3 and 4 m/s are 15 dB higher than expected, 15 dB being the vertical distance between the expected and highest one-hour immission levels at 3- 4 m/s (upper and lower lines in figures 5B and 6).}

\begin{figure}[h]
\centering
\includegraphics[width=\textwidth]{figureIV6.png}
\caption{selected measured sound levels L_{eq,T} (T = 20 – 60 min) at location A with best fit; and expected sound levels according to a logarithmic wind profile (v_{98}/v_{10} = \text{f}_{\text{log}} = 1.4), a stable wind profile (v_{98}/v_{10} = 1.8\text{f}_{\text{log}}) and with the maximum wind speed ratio (v_{98}/v_{10} = 2.6\text{f}_{\text{log}})}
\end{figure}
The same lines as in figure IV.5B, but valid for location B, are plotted in figure IV.5D; immission levels here exceed the calculated levels, even if calculated on the basis of a 2.6 higher wind velocity at hub height. An explanation may be that a lower ambient sound level is necessary compared to location A to allow wind turbine sound to be dominant at location B (as selected with the $L_5 - L_{95}$ -criterion), implying a lower near ground wind velocity and thus a higher stability. It may also be caused by an underestimate of actual sound level in the calculation model for long distances, at least for night conditions (this issue will be addressed in section IV.10).

As is clear from the wind velocity at 2 m height plotted in figure IV.5B, there is only a very light wind near the ground even when the turbines rotate at high power. This implies that in a quiet area with low vegetation the ambient sound level may be very low. The contrast between the turbine sound and the ambient sound is therefore at night higher than in daytime.

Although at most times the wind turbine sound dominates the sound levels in figure IV.5, it is possible that at low sound levels, i.e. at low rotational speeds and low wind velocities, the $L_5-L_{95}$-criterion is met while the sound level is not entirely determined by the wind turbines. This is certainly the case at levels close to 20 dB(A), the sound level meter noise floor. The long term night-time ambient background level, expressed as the 95-percentile ($L_{95}$) of all measured night-time sound levels on location B, was 23 dB(A) at 3 m/s ($V_{10}$) and increasing with 3.3 dB/m·s$^{-1}$ up to $V_{10} = 8$ m/s [Van den Berg et al 2002]. Comparing this predominantly non-turbine background level with the sound levels in figure IV.5B and 5D, it is clear that the lowest sound levels may not be determined by the wind turbines, but by other ambient sounds (and instrument noise). This wind velocity dependent, non-turbine background sound level $L_{95}$ is, however, insignificant with respect to the highest measured levels. Thus, the high sound levels do not include a significant amount of ambient sound not coming from the wind turbines. This has also been verified in a number of evenings and nights by personal observation.
IV.8 Comparison of emission and immission sound levels

From the 30 measurements of the equivalent sound level $L_{eq,T}$ (with $T$ typically 5 minutes) measured at distance $R$ from the turbine hub ($R$ typically $100\sqrt{2}$ m), a relation between sound power level $L_w$ and rotational speed $N$ of a turbine could be determined: see formula (IV.3). This relation can be compared with the measured immission sound level $L_{imm,T}$ ($T = 5$ minutes) at location A, 400 m from the wind farm (closest turbine), in 22 cases where the rotational speed was known. The best logarithmic fit for the data points of the immission sound level $L_{imm}$ as a function of rotational speed $N$ is:

$$L_{imm} = 57.6\cdot\log(N) - 30.6\ \text{dB(A)}$$

(IV.4)

with a correlation coefficient of 0.92 and a standard deviation of 1.5 dB with respect to the fit. Both relations from formulae (IV.3) and (IV.4) and the datapoints are given in figure IV.7. The difference between both relations is $L_w - L_{imm} = 9.5\cdot\log(N) + 46.0\ \text{dB}$. For the range 14 – 20 rpm, where both series have data points, the average difference is 57.9 dB, the maximum deviation from this average is 0.8 dB (14 rpm: 57.1 dB(A); 20 rpm: 58.6 dB(A); see lower part of figure IV.7). It can be shown by calculation that about half of this deviation can be explained by the variation of sound power spectrum with increasing speed $N$.

The sound immission level can be calculated with formula (IV.1). For location A, assuming all turbines have the same sound power $L_w$, this leads to $L_w - L_{imm} = 58.0\ \text{dB}$. This is independent of sound power level or rotational speed, as it is calculated with a constant spectrum averaged over several turbine conditions, i.e. turbine speeds. The measured difference (57.9 dB) matches very closely the calculated difference (58.0 dB).

The variation in sound immission level at a specific wind velocity $V_{10}$ in figures IV.5B and IV.5D is thus seen to correspond to a variation in rotational speed $N$, which in turn is related to a variation in wind velocity.
at hub height, not to a variation in $V_{10}$. At location A, $N$ can be calculated from the measured immission level with the help of formula (IV.4) or its inverse form: $N = 3.4 \cdot 10^{\frac{L_{\text{imm}}}{57.6}}$.

**Figure IV.7:**

turbine sound power levels $L_w$ measured near wind turbines and immission levels $L_{\text{imm}}$ measured at 400 m from wind farm: averages differ 57.9 dB; (below) increase of difference $L_w - L_{\text{imm}}$ with rotational speed ($L_w$ data points taken from figure IV.3)

### IV.9 Atmospheric stability and Pasquill class

In figure IV.5 measurement data have been separated in two sets according to atmospheric stability in Pasquill classes, supplied by KNMI from their measurement site Eelde, 40 km to the west of our measurement site. Although the degree of stability will not always be the same for Eelde and our measurement location, the locations will correlate to a high degree in view of the relatively small distance between them. For night-time conditions ‘stable’ refers to Pasquill classes E and F (lightly to very stable) and corresponds to $V_{10} \leq 5$ m/s and cloud coverage $C \leq 50\%$ or $V_{10} \leq 3.5$ m/s and $C \leq 75\%$, ‘neutral’ (class D) corresponding to all other situations. Although from figure IV.5 it is clear that the very highest sound levels at an easterly wind ($\approx 80^\circ$) do indeed occur in stable conditions, it is also
clear that in neutral conditions too the sound level is higher than expected for most of the time, the expected values corresponding to the grey lines in figures IV.5B and D, derived from daytime conditions. According to this study the sound production, and thus wind velocity at 100 m height is at night often higher than expected, in a stable, but also in a neutral atmosphere. On the other hand, even in stable conditions sound levels may be lower than expected (i.e. below the grey lines), although this occurs rarely. It may be concluded from these measurements that a logarithmic wind profile based only on surface roughness does not apply to the nighttime atmosphere in our measurements, not in a stable atmosphere and not always in a neutral atmosphere when determined from Pasquill classes.

**IV.10 Additional measurements**

In several nights in the period that the SMWS was measuring at location A, manual measurements were performed at a number of locations in the area between 0.6 and 2.3 km west of the wind farm. The locations are plotted in figure IV.2. Most locations were close to dwellings, but two (locations U and X) were in open fields. Locations P and Q are close and at the same distance from the western row of turbines and can be considered equal with respect to the turbines (Q was chosen instead of P as P was at the verge of a garden with a loud bird chorus in the early morning). The surface of most of the area is covered with grass and low crops, with trees at some places.

For these measurements one or more logging sound level meters (accuracy type 1 or 2) were used simultaneously, storing a broad band A-weighted sound pressure level every second. Before and after measurement the meters were calibrated with a 94.0 dB, 1000 Hz calibration source, and as a result measurement accuracy due to the instruments is within 0.2 dB. On every location the microphone was in a 10 cm spherical foam wind screen approximately 1.2 m above the surface. There were no reflections of the wind turbine sound to the microphone, except via the ground.
### IV.10.1 Measured and calculated immission sound levels

Figure IV.8 gives a simultaneous registration from just before midnight on May 17, 2002, till noon on May 18, of the equivalent sound pressure levels per 5 minutes at locations A (from the SWMS), P/Q and U (from the manual meters) at distances to the westernmost row of turbines of 400, 750 and 1050 m, respectively. In the night hours the sound of the turbines was dominant at each of these locations, apart from an occasional bird or car. Also plotted in figure IV.8 are the wind velocity at 2 and 10 m heights at location A.

![Figure IV.8: measured sound immission level (Leq.5min) at locations A, P/Q, and U, and wind velocities at A with an eastnortheasterly wind](image)

A short decrease in wind velocity at around 2:00 is apparently accompanied by a similar decrease in wind velocity at hub height, as the sound level varies much in the same way. However, the registrations show that the sound level increases from 0:30 until 6:00 while the 10-m wind velocity does not show a net increase in this period. In fact the sound level at location A at 3:00 implies a rotational speed of 21 rpm, which is just below maximum (22 rpm), even though the wind velocity at 10 m height is...
only 4.5 m/s and at 2 m height is less than 1 m/s. Only occasionally there are other sounds until the dawn chorus of birds just after 4:00 and after that the near-ground wind picks up.

In figure IV.9 the 5-minute equivalent sound levels at P/Q and U relative to the sound level at A are plotted. The advantage of taking the sound level at A as a reference value is that it is not necessary to know the exact sound power level of the turbines themselves. The level differences are 3.5 and 6.5 dB, respectively, with a variation of ±1 dB. The variations must be due to differences in sound propagation mostly, because other disturbances (such as one at 23:55 at P) are rare.

![Figure IV.9](image)

**Figure IV.9: difference between simultaneously measured broad band A-weighted immission levels at locations U and P/Q and at location A**

Comparable simultaneous measurements have been made in the night of June 2 - 3 and of June 17 - 18, 2002. In Appendix C the registrations are given, as well as the level differences between the distant locations P through T, V and X and the reference location A. The measured and calculated decrease in sound level with distance, relative to location A, as well as the discrepancy between both, are given in table IV.4 and figure IV.10. In all cases the wind was easterly (60° – 100°), that is: from the
wind farm to the measurement location. Also there was little near-ground wind and low background sound levels from other sources.

The calculated differences have been determined with equation IV.1 and the Dutch national model [VROM 1999]. The measured differences in table IV.4 are the difference in the equivalent sound level at a location minus the same at location A over the given measurement time T; only very few of the $L_{eq,5min}$ values were omitted from this $L_{eq,T}$ because they were apparently disturbed by another sound. To minimize influence of possible disturbing sounds the median of all $L_{eq,5min}$ values can be used, as this value gives the prevailing difference and is thus less sensitive to the influence of disturbances; this, however, yields the same results within 0.5 dB.

The discrepancies between measured and calculated levels are small, especially considering the large distances involved: -0.2 to 1.5 dB. One may conclude that the calculation model is quite satisfactory in this relatively simple situation (a high sound source above flat ground).

Table IV.4: measured and calculated differences in sound level $L_{eq,T}$ at locations R - T and at location A, when wind blows from the wind farm

<table>
<thead>
<tr>
<th>location</th>
<th>R</th>
<th>P/Q</th>
<th>U</th>
<th>V</th>
<th>S</th>
<th>X</th>
<th>T</th>
</tr>
</thead>
<tbody>
<tr>
<td>distance to western row wind farm (m)</td>
<td>600</td>
<td>750</td>
<td>1000</td>
<td>1100</td>
<td>1250</td>
<td>1900</td>
<td>2250</td>
</tr>
<tr>
<td>date of measurement (in 2002)</td>
<td>June 2/3</td>
<td>May 17/18, June 2/3 +18</td>
<td>May 17/18</td>
<td>June 18</td>
<td>June 2/3</td>
<td>June 18</td>
<td>June 2/3</td>
</tr>
<tr>
<td>measurement time T (min.)</td>
<td>200</td>
<td>295 +200+115</td>
<td>120</td>
<td>140</td>
<td>190</td>
<td>85</td>
<td>195</td>
</tr>
<tr>
<td>measured difference</td>
<td>-3.5</td>
<td>-3.8 *</td>
<td>-6.4</td>
<td>-9.1</td>
<td>-8.5</td>
<td>-12.1</td>
<td>-1.3</td>
</tr>
<tr>
<td>calculated difference</td>
<td>-4.5</td>
<td>-4.1</td>
<td>-6.6</td>
<td>-10.6</td>
<td>-8.3</td>
<td>-13.1</td>
<td>-14.2</td>
</tr>
<tr>
<td>discrepancy calculation - measurement</td>
<td>-1.0</td>
<td>-0.3</td>
<td>-0.2</td>
<td>-1.5</td>
<td>0.2</td>
<td>-1.0</td>
<td>-12.9</td>
</tr>
</tbody>
</table>

*: measurement time weighted logarithmic average of resp. 3.5, 3.6 and 4.6 dB
In figure IV.10 a line is plotted corresponding to \(-20 \cdot \log(R/R_a)\), where \(R_a\) is the distance from A to the western turbine row. This decrease corresponds to spherical divergence from a point source only, with no attenuation due to absorption. It is clear that, with the exception of location T (see next section), the measured decrease is close to this spherical divergence: the measured values at the locations P/Q, U, S and X are 1.4 to 1.7 dB above the plotted line, at the more northern locations R and V they are 0 to 0.3 dB below the line. Approximately the same is true for the calculated levels: the calculated values at the locations P/Q, U, S and X are 0.4 to 1.6 dB above the plotted line, at the more northern locations R and V they are 1.0 to 1.8 dB below the line.

There are two counteracting causes explaining this apparently ‘almost spherical’ attenuation. The first is that the wind farm cannot be considered a point source. Due to its large dimension (3 km from south to north, see figure IV.2) normal to the shortest distance from location A and locations further west, the geometrical divergence should be between cylindrical and
spherical divergence, that is: proportional to \(-X \cdot \log(R/R_A)\), with \(10 < X < 20\). Secondly one expects a decrease due to absorption ('excess attenuation') above the decrease due to geometrical divergence: for the Rhede turbines calculation shows that this excess attenuation is expected to be 1.7 dB per km.

**IV.10.2 Immission level increase due to inversion layer?**

In the night of June 2 to 3, 2002, high sound levels were measured at the most distant measurement location T, 2250 m from the wind farm. The immission sound level varied between approximately 40 and 45 dB(A) and was more variable than at the other locations (see Appendix C). The resident close to this measurement location could hear the wind farm well, at 22:30 hours describing it as: “The sound changes from ‘an endless train’ to a more pulsating sound; the sound grows louder and sharper. At the background is a kind of humming, comparable to the sound of a welding transformer”. The sound was audible indoors.

In our research we have not met this phenomenon again. However, mr. Flight living near another wind farm south of the Rhede wind farm observed the same phenomenon: on a location appr. 750 m from the closest turbine, where at night he usually measured an immission level of 42 to 44 dB(A), he measured a level of 50 to 52 dB(A) in the night of September 24, 2002. It was clear that the sound came from the nearest wind farm, but also from a second, more distant wind farm that usually was not audible here. Again, the atmosphere was stable and there was a weak near-ground easterly wind, blowing from the wind farm to the observer.

This may be a result of strong refraction of sound below an inversion layer. This inversion layer must be at or above the rotor to have the highest effect, so at or above 130 m (= hub height + blade length).

Suppose the turbines in the Rhede wind farm each have a sound power level \(L_W\) at a certain wind velocity. If we substitute the entire farm by one single turbine at the site of the turbine closest to location T (nr. 12), it can be calculated that the sound level of that single turbine must be \(L_W + 9.4\) to produce the same immission level at T as the entire wind farm.
Considering only spherical spreading, this immission level is \( L_{\text{imm}} = L_w + 9.4 - 10 \cdot \log(4\pi \cdot 2250^2) = L_w - 68.6 \). Now the sound waves will be refracted downwards at the inversion layer and we assume that all sound propagates below the inversion layer. At large distances (>> height inversion layer) this is equivalent to sound spreading cylindrically from a vertical line source. To simulate this we replace the substitute single turbine, which was modelled as a point source at hub height, by a vertical line source from the ground up to the inversion layer height (130 m). If the sound power levels of both point and line source are equal, the line source must have a sound power level of \( L_w' = L_w + 9.4 - 10 \cdot \log(130) = L_w - 11.7 \) dB/m. If again the sound level decreases by geometrical (now: cylindrical) spreading only, the sound immission level at 2250 m from this line source is \( L_{\text{imm}}' = L_w - 11.7 - 10 \cdot \log(2\pi \cdot 2250) = L_w - 54.6 \) dB. Comparison of the immission level due to a point source (\( L_w - 68.6 \)) and a line source (\( L_w - 54.6 \)) shows that the line source causes a 14 dB higher immission level. This simple calculation shows that the rise in level caused by a simplified high inversion layer is close to the observed increase (13 dB): the higher level is a result of the sound being 'trapped' below the inversion layer. However, more observations and data are needed to verify this hypothesis.

**IV.11 Conclusion**

Sound immission measurements have been made at 400 m (location A) and 1500 m (location B) from the wind farm Rhede with 17 tall (98 m hub height), variable speed wind turbines. It is customary in wind turbine noise assessment to calculate immission sound levels assuming wind velocities based on wind velocities \( V_{10} \) at reference height (10 m) and a logarithmic wind profile. Our study shows that the immission sound level may, at the same wind velocity \( V_{10} \) at 10 m height, be significantly higher in nighttime than in daytime. A 'stable' wind profile predicts a wind velocity \( V_h \) at hub height 1.8 times higher than expected and agrees excellently with the average measured night-time sound immission levels. Wind velocity at hub height may still be higher: at low wind velocities \( V_{10} \) up to 4 m/s, the wind velocity \( v_h \) is at night up to 2.6 times higher than expected.
Thus, the logarithmic wind profile, depending only on surface roughness and not on atmospheric stability, is not a good predictor for wind profiles at night. Especially for tall wind turbines, estimates of the wind regime at hub height based on the wind velocity distribution at 10 m, will lead to an underestimate of the immission sound level at night: at low wind velocities ($V_{10} < 5$ m/s) the actual sound level will be higher than expected for a significant proportion of time. This is not only the case for a stable atmosphere, but also -to a lesser degree- for a neutral atmosphere.

The change in wind profile at night also results in lower ambient background levels then expected: at night the wind velocity near the ground may be lower than expected from the velocity at 10 m and a logarithmic wind profile, resulting in low levels of wind induced sound from vegetation. The contrast between wind turbine and ambient sound levels is therefore at night more pronounced.

Measured immission sound levels at 400 m from the nearest wind turbine almost perfectly match (average difference: 0.1 dB) sound levels calculated from measured emission levels near the turbines. From this it may be concluded that both the emission and immission sound levels could be determined accurately, even though the emission measurements were not fully in agreement with the standard method. As both levels can be related through a propagation model, it may not be necessary to measure both: the immission measurements can be used to assess immission as well as emission sound levels. At greater distances the calculated level may underestimate the measured level, but considering the distances involved (up to 2 km) the discrepancy is small: 1.5 dB or less.

In one night the sound level at a distant location (over 2 km from the wind farm) was much higher than expected, perhaps because of an inversion layer adding more downward refracted sound. It apparently is a rare occurrence at the Rhede wind farm, and could be more significant where high inversion layers occur more often.
V.1 Effects of atmospheric stability

Atmospheric stability is not only relevant for wind turbine sound levels, as we saw in the preceding chapter, but also for the character of the sound. In conditions where the atmosphere is stable, distant wind turbines can produce a beating or thumping sound that is not apparent in daytime.

The magnitude of the effects of increasing stability depends on wind turbine properties such as speed, diameter and height. We will use the dimensions of the wind turbines in the Rhede wind farm, that are typical for a modern variable speed 2 MW wind turbine: hub height 100 m, blade length 35 m and blade tip speed increasing with wind velocity up to a maximum value of $\Omega \cdot R = 81 \text{ m/s}$ (at 22 rpm). Here a speed of 20 rpm (70 m/s) will be used as this was typical for situations where at the Rhede wind farm a clear beating sound was heard.

We will assume the optimum angle of attack $\alpha$ is 4°. The change in trailing edge (TE) sound pressure level $\Delta \text{SPL}_{\text{TE}}$ with the angle of attack from this optimum up to 10° can be approximated by $\Delta \text{SPL}_{\text{TE}}(\alpha) = 1.5 \cdot \alpha - 1.2 \text{ dB}$ or $d(\Delta \text{SPL}_{\text{TE}})/d\alpha = 1.5$ (see appendix B, equation B.8). When the pitch angle is constant, the change in angle of attack due to a variation $dV$ in wind velocity is $d\alpha = 0.84 \cdot dV$ (see appendix B, equation B.9).

To calculate vertical wind velocity gradients the simple engineering formula (III.1) will be used: $V_h = V_{\text{ref}}(h/h_{\text{ref}})^m$ (see section III.2). In the text below we will use a value $m = 0.15$ for a daytime atmosphere (unstable – neutral), $m = 0.4$ for a stable, and $m = 0.65$ for a very stable atmosphere (see table III.1). These values will be used for altitudes between 10 and 120 m.

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1 A value $m = 0.65$ is not obvious from table III.1, but is chosen as a relatively high value that is exceeded for a small part of the time (see figures VI.6 and VI.16, and section VI.6)
There are now three factors influencing blade swish level when the atmosphere becomes more stable: a) the higher wind velocity gradient, b) the higher wind direction gradient, and c) the relative absence of large scale turbulence.

\(a. \text{ Wind velocity gradient.} \) Rotational speed is determined by a rotor averaged wind velocity, which here is assumed to be the induced wind velocity at hub height (equation III.5). The free, unobstructed wind at height \( h \) is denoted by \( V_h \), the induced wind speed at the blade by \( V_{h,b} \). With increasing atmospheric stability the difference in wind velocity between the upper and lower part of the rotor increases. As in a complete rotation the pitch angle is constant the change in angle of attack due to a change in induced wind velocity is 
\[
\Delta \alpha = 0.82 \cdot dV_{h,b}
\]
which can be expressed in a change of the free wind velocity by 
\[
\Delta \alpha = 0.82 \cdot (2/3) \cdot dV_h = 0.55 \cdot dV_h
\] (see equation III.5).

Suppose that the free wind velocity at hub height is \( V_{100} = 14 \) m/s, corresponding to \( V_{10} = 9.8 \) m/s in a neutral atmosphere in flat open grass land (roughness length 5 cm). Then in daytime \( (m = 0.15) \) the free wind velocity at the height of the lowest point of the rotor would be \( V_{65} = 13.1 \) m/s, at the height of the highest point \( V_{135} = 14.6 \) m/s (corresponding to velocities at the blade of \( V_{65,b} = 8.7 \) m/s and \( V_{135,b} = 9.7 \) m/s, respectively). The difference of 1.0 m/s between the low tip and hub height wind velocities causes a change in angle of attack on the blade of \( \Delta \alpha = 0.55^\circ \). Between the high tip and hub height the change is smaller and of opposite sign: -0.3°. In a stable atmosphere \( (m = 0.4) \), at the same wind velocity at hub height, \( V_{65} \) is 11.8 m/s causing a change in angle of attack at the lower tip relative to hub height of 1.2° (at the high tip: \( V_{135} = 15.8 \) m/s, \( \Delta \alpha = -1.0^\circ \)). When the atmosphere is very stable \( (m = 0.65) \), wind velocity \( V_{65} = 10.5 \) m/s and the angle of attack on the low altitude tip deviates 1.9° from the angle at hub height (at the high tip: \( V_{135} = 17.0 \) m/s, \( \Delta \alpha = -1.7^\circ \)). In fact when the lower tip passes the tower there is a greater mismatch between optimum and actual angle of attack \( \alpha \) because there was already a change in angle of attack related to the wind velocity deficit in front of the tower. For a daytime atmosphere and with respect to the situation at hub
height, the change in \( \alpha \) associated to a blade swish level of \( 2 \pm 1 \) dB is estimated as \( 1.8 \pm 1.1^\circ \) (see appendix B.3), part of which (\( 0.55^\circ \)) is due to the wind profile and the rest to the tower. The increase in \( \alpha \) due to the stability related wind profile change must be added to this daytime change in \( \alpha \). Thus, the change in angle of attack when the lower tip passes the mast is \( 1.8 \pm 1.1^\circ \) in daytime (unstable to neutral atmosphere), increasing to \( 2.5 \pm 1.1^\circ \) in a stable atmosphere and to \( 3.2 \pm 1.1^\circ \) in a very stable atmosphere. The associated change in TE sound level is \( 3.8 \pm 1.7 \) dB for a stable and \( 4.8 \pm 1.7 \) dB for a very stable atmosphere (compared to \( 2 \pm 1 \) dB in daytime), which is the increase when the blade passes the tower. The corresponding total A-weighted sound level will be somewhat less as trailing edge sound is not the only sound source (but it is the dominant source; see section V.2.3).

At the high tip the change in angle of attack is smaller and of opposite sign with respect to the low tip, and also there is no (sudden) tower induced change to add to the wind gradient dependent change. The change in angle of attack at the high tip in a very stable atmosphere (-1.7°) is comparable to the change at the low tip in daytime, and this change is more gradual than for the low tip. This in fact lowers the sound emission from the high tip (with approximately 2 dB), most so when the high blade is vertical so just before and just after the low blade passes the tower, thereby in fact increasing the variation in swish sound level even more.

Thus we find that, for \( v_{100} = 14 \) m/s, the 1-2 dB daytime blade swish level increases to approximately 5 dB in a very stable atmosphere. The effect is stronger when wind velocity increases, up to the point where friction turbulence overrides stability and the atmosphere becomes neutral. The increase in trailing edge sound level will be accompanied by a lower peak frequency (see appendix B, equation B.2). For \( \Delta \alpha = 5^\circ \) the shift is one octave.

\( b. \) Wind direction gradient. In a stable atmosphere air masses at different altitudes are only coupled by small scale turbulence and are therefore relatively independent. Apart from a higher velocity gradient a higher wind direction gradient is also possible, and with increasing height the wind
direction may change significantly. This wind direction shear will change the angle of attack with height. Assuming the wind at hub height to be normal to the rotor, the angle of attack will decrease below and increase above hub height (or vice versa). This effect, however, is small: if we suppose a change in wind direction of 20° over the rotor height at an induced wind velocity of 10 m/s, the change in angle of attack between extreme tip positions at 20 rpm is only 0.25°, which is negligible relative to the wind velocity shear.

c. Less turbulence. In a stable atmosphere turbines in a wind farm can run almost synchronously because the absence of large scale turbulence leads to less variation superimposed on the constant (average) wind velocity at each turbine. In unstable conditions the average wind velocity at the turbines will be equal, but instantaneous local wind velocities will differ because of the presence of large, turbulent eddies at the scale of the inter-turbine distance. In a stable atmosphere the turbulence scale decreases with a factor up to 10, relative to the neutral atmosphere and even more relative to an unstable atmosphere [Garratt 1992]. In stable conditions turbines in a wind farm therefore experience a more similar wind and as a consequence their instantaneous speeds are more nearly equal. This is confirmed by long term measurements by Nanahara et al. [2004] who analysed coherence of wind velocities between different locations in two coastal areas. At night wind velocities at different locations were found to change more coherently than they did at daytime [Nanahara 2004]. The difference between night and day was not very strong, probably because time of day on its own is not a sufficient indicator for stability. The decay of coherence was strongly correlated with turbulence intensity, which in turn is closely correlated to stability.
Thus several turbines can be nearly synchronous: sometimes two or more turbines are in phase and the blade passing pulses coincide, then they go out of phase again. Synchronicity here refers to the sound pulses from the

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1 In a coastal location atmospheric stability also depends on wind direction as landwards stability is a diurnal, but seawards a seasonal phenomenon. Also, a fixed duration for all nights in a year does not coincide with the time that the surface cools (between sundown and sunrise), which is a prerequisite for stability.
different turbines as observed at the location of the observer: pulses synchronise when they arrive simultaneously. This is determined by differences in phase (rotor position) between turbines and in propagation distances of the sound from the turbines. Phase differences between turbine rotors occur because turbines are not connected and because of differences in actual performance. The place where synchronicity is observed will change when the phase difference between turbines changes. With exact synchronicity there would be a fixed interference pattern, with synchronicity at fixed spots. However, because of near-synchronicity, synchronous arrival of pulses will change over time and place and an observer will hear coinciding pulses for part of the time only.

Near a wind farm the variation in sound level will depend on the distances of the wind turbines relative to the observer: the level increase due to several turbines will reach higher levels when more turbines are at approximately equal distances and thus contribute equal immission levels. The increase in level variation, or beating, is thus at well-audible frequencies and has a repetition rate equal to the blade passing frequency.

A second effect of the decrease in turbulence strength is that in-flow turbulent sound level also decreases. The resulting decrease in sound level at frequencies below that of TE noise lowers the minimum in the temporal variations, thereby increasing modulation depth. The higher infrasound level due to extra blade loading is not perceptible because of the high hearing threshold at the very low blade passing frequency and its harmonics.

Thus, theoretically it can be concluded that in stable conditions (low ambient sound level, high turbine sound power and higher modulation or swish level) wind turbine sound can be heard at greater distances where it is of lower frequency due to absorption and the frequency shift of swish sound. It will thus be a louder and more low frequency ‘thumping’ sound and less the swishing sound that is observed close to a daytime wind turbine.
V.2 Measurement results

V.2.1 Locations

In the summer of 2002 and of 2004 wind turbine sound has been recorded in and near the Rhede wind farm (see section IV.1 for a specification of the turbines and a map of the area). In this chapter measurement results will be used from two locations: R and P (see figure IV.2). Location R is close to a dwelling west of the turbines, 625 m from the nearest turbine. The microphone position was at 4 m height and close to the house, but with no reflections except from the ground. Location P, 870 m south of R, was 1.5 m above a paved terrace in front of the façade of a dwelling at 750 m distance from the nearest turbine (in fact this is a short distance from the location P in chapter IV, which was not in front of the façade). The entire area is quiet, flat, agricultural land with some trees close to the dwellings. There is little traffic and there are no significant permanent human sound sources.

A third dwelling Z is in Boazum in the northern part of the Netherlands, 280 m west of a single, two-speed turbine (45 m hub height, 23 m blade length, 20/26 rpm). The area is again quiet, flat and agricultural, with some trees close to the dwelling. The immission measurement point is at 1.5 m height above gravel near dwelling Z. This measurement site is included here to show that the influence of stability on blade swish levels occurs also with smaller and single turbines. At all locations near dwellings the microphone was fitted in a 9 cm diameter foam wind screen.

Table V.1 gives an overview of measurement (start) time and date, of observed turbine speed and of wind velocity and direction, for situations of which results will be given below. The wind velocity at hub height $V_{\text{hub}}$ has been determined from turbine rotation speed $N$ or sound power level $L_w$ (figure III.3, the relation $V_{\text{hub}} - N$ follows from \[\text{Kerkers 1999}\] and \[\text{Van den Berg 2002}\]). The wind velocity $V_{10}$ was continuously measured at or near location A, except for location Z, where data from several meteorological stations were used showing that the wind was similar and nearly constant throughout the night of the measurement in the entire northern part of the Netherlands. In all cases there were no significant variations in wind velocity at the time of measurement. Wind velocity at
the microphone was lower than \(V_{10}\) because of the low microphone height and shelter provided by trees nearby. Wind direction is given in degrees relative to north and clockwise (90° is east).

The spectra near a turbine were measured with the microphone just above a hard surface at ground level 100 m downwind of a turbine in compliance with IEC 61400 [IEC 1998] as much as possible (non-compliance did not lead to differences in result; for reasons of non-compliance, see section II.4). The levels presented here are broad band immission levels: measured \(\text{L}_{eq}\) minus 6 dB correction for coherent reflection against the hard surface [IEC 1998]. The presented levels near the dwellings are also broad band immission levels: measured \(\text{L}_{eq}\) minus 3 dB correction for incoherent reflection at the façade for dwelling P, or measured \(\text{L}_{eq}\) without any correction for dwellings R and Z.

Table V.1: overview of measurement locations and times and of turbine speed and wind

<table>
<thead>
<tr>
<th>Location</th>
<th>measurement date</th>
<th>measurement time</th>
<th>turbine speed (rpm)</th>
<th>wind velocity (V_{10}) (m/s)</th>
<th>wind velocity (V_{hub}) (m/s)</th>
<th>wind direction (° north)</th>
</tr>
</thead>
<tbody>
<tr>
<td>Dwelling P</td>
<td>June 3, 2002</td>
<td>00:45</td>
<td>20</td>
<td>5</td>
<td>14</td>
<td>100</td>
</tr>
<tr>
<td>Turbine 7</td>
<td>June 3, 2002</td>
<td>06:30</td>
<td>19</td>
<td>5</td>
<td>15</td>
<td>100</td>
</tr>
<tr>
<td>Turbine 1</td>
<td>June 3, 2002</td>
<td>06:45</td>
<td>19</td>
<td>5</td>
<td>15</td>
<td>100</td>
</tr>
<tr>
<td>Dwelling R</td>
<td>Sep.9, 2004</td>
<td>23:07</td>
<td>18</td>
<td>4</td>
<td>14</td>
<td>80</td>
</tr>
<tr>
<td>Turbine 16</td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>Dwelling Z</td>
<td>Oct.18, 2003</td>
<td>01:43</td>
<td>26</td>
<td>3</td>
<td>6</td>
<td>60</td>
</tr>
</tbody>
</table>

At dwelling P at the time of measurement the beat in the turbine sound was very pronounced. In the other measurements (dwellings R and Z) the beating was not as loud. The measurements near turbine 16 and dwelling R at 23:07 on September 9 were performed simultaneously.

V.2.2 Frequency response of instruments

For the Rhede measurements in this chapter sound was recorded on a TASCAM DA-1 DAT-recorder with a precision 1" Sennheiser MKH 20
P48 microphone. The sound was then sampled in 1-second intervals on a Larson Davis 2800 frequency analyser. From 1 to 10 000 Hz the frequency response of the DAT-recorder and LD2800 analyser have been determined with a pure tone electrical signal as input. The LD2800 response is flat (±1 dB) for all frequencies. The DAT-recorder is a first order high pass filter with a corner frequency of 2 Hz. The frequency response of the microphone was of most influence and has been determined relative to a B&K ½" microphone type 4189 with a known frequency response [B&K 1995]. Equivalent spectral sound levels with both microphones in the same sound field (10 cm mutual distance) were compared. For frequencies of 2 Hz and above the entire measurement chain is within 3 dB equivalent to a series of two high pass filters with corner frequencies of $f_1 = 4 \text{ Hz}$ and $f_2 = 9 \text{ Hz}$, or a transfer function equal to $-10 \log[1+(f_1/f)^2] -10 \log[1+(f_2/f)^2]$. For frequencies below 2 Hz this leads to high signal reductions ($<-40 \text{ dB}$) and consequentially low signal to (system) noise ratios. Therefore values at frequencies $<2 \text{ Hz}$ are not presented.

For the Boazum measurements sound was recorded on a Sharp MD-MT99 minidisc recorder with a 1" Sennheiser ME62 microphone. The frequency response of this measurement chain is not known, but is assumed to be flat in the usual audio frequency range. Simultaneous measurement of the broad band A-weighted sound level were done with a precision (type 1) 01dB sound level meter. Absolute precision is not required here as the minidisc recorded spectra are only used to demonstrate relative spectral levels. Because of the ATRAC time coding of a signal, a minidisc recording does not accurately follow a level change in a time interval $<11.6 \text{ ms}$. This is insignificant in the present case as the ‘fast’ response time of a sound level meter is much slower (125 ms).

V.2.3 Measured emission and immission spectra

Recordings were made at evening, night or early morning. On June 3, 2002, sound was recorded at dwelling P at around midnight and early in the morning near two turbines (numbers 1 and 7 in figure IV.1). At P at these times a distinct beat was audible in the wind turbine sound. In figure V.1, 1/3 octave band spectra of the recorded sound at P and at both
turbines have been plotted. In each figure A, B and C, 200 sound pressure spectra sampled in one-second intervals, as well as the energy averaged spectrum of the 200 samples have been plotted. The standard deviation of 1/3 octave band levels is typically 7 dB at very low frequencies, decreasing to approx. 1 dB at 1 kHz. The correlation coefficient $\rho$ between all 200 unweighted 1/3 octave band levels and the overall A-weighted sound level has also been plotted for each 1/3 octave band frequency.

For frequencies below approximately 10 Hz the sound is dominated by the thickness sound associated with the blade passing frequency and harmonics. In the rest of the infrasound region and upwards, in-flow turbulence is the dominant sound producing mechanism. Gradually, at frequencies above 100 Hz, trailing edge sound becomes the most dominant source, declining at high frequencies of one to several kHz. Trailing edge sound is more pronounced at turbine 1 (T1) compared to turbine 7 (T7), causing a hump near 1000 Hz in the T1 spectra. At very high frequencies (> 2 kHz) sometimes spectral levels are influenced by birds’ sounds.

It is clear from the spectra that most energy is found at lower frequencies. However, most of this sound is not perceptible. To assess the infrasound level relevant to human perception it can be expressed as a G-weighted level [ISO 1995], With G-weighting sound above the infrasound range is suppressed. The average infrasound perception threshold is 95 dB(G) [Jakobsen 2004]. The measured G-weighted levels are 15-20 dB below this threshold: 80.5 and 81.1 dB(G) near turbines 1 and 7 respectively, and 76.4 dB(G) at the façade.

The correlations show that variations in total A-weighted level near the turbines are correlated with the 1/3 octave band levels with frequencies from 400 through 3150 Hz (where $\rho > 0.4$), which is trailing edge sound. This is one octave lower (200 - 1600 Hz) for the sound at the façade: the higher frequencies were better absorbed during propagation through the atmosphere.
Figure V.1:
left axis:
200 consecutive, unweighted and 1 second spaced 1/3 octave band spectra (thin lines), and averaged spectrum (thick line) of sound pressure level $L_p$ near turbines 1 (A) and 7 (B) and near dwelling P (C);
right axis:
coefficient of correlation (line with markers) at each 1/3 octave band frequency between all 200 1/3 octave band levels and overall $A$-weighted level.
The façade spectra in figure V.1C show a local minimum at 50-63 Hz, followed by a local maximum at 80-100 Hz.\textsuperscript{1} This is caused by interference between the direct sound wave and the wave reflected by the façade at 1.5 m from the microphone: for wave lengths of approximately 6 m (55 Hz) this leads to destructive interference, for wave lengths of 3 m (110 Hz) to constructive interference.

In figure V.2A the three average spectra at the same locations as in figure V.1A-C have been plotted, but now for a total measurement time of 9.5 (façade), 5 (T7) and 6 (T1) minutes. For each of these measurement periods the average of the 5\% of samples with the highest broad band A-weighted sound level (i.e. the equivalent spectral level of the \textit{L}_{A5} percentile) has also been plotted, as well as the 5\% of samples with the lowest broad band level (\textit{L}_{A95}). The range in A-weighted broad band level can be defined as the difference between the highest and lowest value: \textit{R}_{bb} = \textit{L}_{A_{\text{max}}} - \textit{L}_{A_{\text{min}}}\. Similarly the range per 1/3 octave or octave band \textit{R}_{f} can be defined by the difference in spectral levels corresponding to \textit{L}_{A_{\text{max}}} and \textit{L}_{A_{\text{min}}}\. The difference between \textit{L}_{A5} and \textit{L}_{A95} is a more stable value, avoiding possibly incidental extreme values, especially when spectral data are used. \textit{R}_{bb,90} is defined as the difference in level between the 5\% highest and the 5\% lowest broad band sound levels: \textit{R}_{bb,90} = \textit{L}_{A5} - \textit{L}_{A95}. For spectral data, \textit{R}_{f,90} is the difference between spectral levels associated with \textit{L}_{A5} and \textit{L}_{A95}. Values of \textit{R}_{f,90} are plotted in the lower part of figure V.2A (here octave band levels have been used to avoid the somewhat ‘jumpy’ behaviour of the 1/3 octave band levels). Close to turbines 1 and 7 \textit{R}_{bb} is 4.8 and 4.1 dB, respectively. \textit{R}_{bb,90} is 3.2 and 2.6 dB, which is almost the same as \textit{R}_{f,90} (3.2 and 3.0 dB) at 1000–4000 Hz. Further away, at the façade, \textit{R}_{bb} is comparable to the near turbine values: 4.9 dB. \textit{R}_{bb,90} at the façade is 3.3 dB and again almost the same as maximum \textit{R}_{f,90} (3.5 dB) at 1000 Hz.

Also, close to the turbine there is a low frequency maximum in \textit{R}_{f,90} at 2 (or 8) Hz that is also present at the façade, indicating that the modulation of trailing edge sound is correlated in time with the infrasound caused by the blade movement.

\textsuperscript{1} In an FFT spectrum minima are at 57 and 170 Hz, maxima at 110 and 220 Hz.\textsuperscript{2}
Figure V.2: upper panels of A, B, C: thick lines: 1/3 octave band $L_{eq}$ near windturbines and dwellings; dotted lines: $L_{eq}$ of all samples with resp. 5% highest (thin dotted lines) and 5% lowest (thick dotted lines) values of broad band $L_A$;

lower panels A, B, C: difference between $L_{eq}$ of 5- and 95-percentile octave band levels = dynamic range $R_{f(90)}$
Figure V.2B presents similar plots for the average spectra and the \( L_{A5} \) and \( L_{A95} \) spectra at dwelling R and near turbine T16, simultaneously over a period of 16 minutes. Close to the turbine the broadband \( R_{bb} \) is 6.2 dB and \( R_{bb,90} \) is 3.7 dB; octave band \( R_{f,90} \) is highest (5.1 dB) at 1000 Hz. Near R broad band \( R_{bb,90} \) is also 3.7 dB, and octave band \( R_{f,90} \) is highest (4.0 dB) at 500 Hz. The \( R_{bb} \) ranges are 2.3–2.5 dB higher than the 90% ranges \( R_{bb,90} \).

In the measurements at this time and place (dwelling R) the infrasound level was lower than in the previous measurements at dwelling P where beating was more pronounced. G-weighted sound level during the 16 minutes at R was 70.4 dB(G), and at T16 77.1 dB(G).

Finally figure V.2C gives average spectra over a period of 16 minutes at dwelling Z. \( R_{f,90} \) is now highest (4.8 dB) at 1 kHz, and broadband \( R_{bb,90} \) is 4.3 dB (\( R_{bb} = 5.9 \) dB). The turbine near Z is smaller and lower, but rotates faster than the Rhede turbines; for a hub height wind velocity of 6 m/s the expected calculated increase in trailing edge sound for the lower tip relative to the day time situation is 2.0 ± 0.8 dB for a stable, and 2.9 ± 0.8 dB for a very stable atmosphere. For this turbine a peak trailing edge sound level is expected (according to equation B.2 in appendix B) at a frequency of 1550/\( b \) Hz ≈ 400 – 800 Hz.

In all cases above the measured sound includes ambient background sound. Ambient background sound level could not be determined separately at the same locations because the wind turbine(s) could not be stopped (see section II.4). However, at audible frequencies it could be ascertained by ear that wind turbine sound was dominant. At infrasound frequencies this could not be ascertained. But if significant ambient sound were present, subtracting it from the measured levels would lead to lower (infrasound) sound levels, which would not change the conclusion, based on the G-weighted level, that measured infrasound must be considered inaudible.

A 25 second part of the 16 min period that corresponds with the spectra in figure V.2B is shown in figure V.3. The broad band level \( L_{A} \) changes with time at T16 and R, showing a more or less regular variation with a period of approximately 1 s (= 1/\( f_{b} \)). Note that the level differences at R are of the
same magnitude as close to the turbine, but the fluctuations at R consist of narrow peaks in comparison to the broader near-turbine fluctuations.

**V.2.4 Beats caused by interaction of several wind turbines**

In the previous section we saw that measured variations in broad band sound level ($R_{bb}$) were 4 to 6 dB. In figure V.4 a registration is given of the sound pressure level every 50 msec over a 180 seconds period, taken from a DAT-recording on a summer night (June 3rd, 0:40 h) on a terrace of dwelling P at 750 m west of the westernmost row of wind turbines (this sound includes the reflection on the façade). In this night stable conditions prevailed ($m = 0.45$ from the wind velocities in table V.1). Turbines 12 and 11 are closest at 710 and 750 m, followed by turbines 9 and 14 at 880 and 910 m. Other turbines are more than 1 km distant and have an at least 4 dB lower immission level than the closest turbine has.

In figure V.4 there is a slow variation of the 'base line' (minimum levels) probably caused by variations in wind velocity and atmospheric sound transmission. There is furthermore a variation in dynamic range: a small difference between subsequent maximum and minimum levels of less than 2 dB is alternated by larger differences.

The expanded part of the sequence in figure V.4 (lower panel) begins when the turbine sound is noisy and constant within 2 dB. After some time (at $t = 155$ s) regular pulses\(^1\) appear with a maximum height of 3 dB, followed by a short period with louder (5 dB) and steeper (rise time up to 23 dB/s)

\(^1\) the term 'pulse' is used to indicate a short, upward variation in sound level
pulses. The pulse frequency is equal to the blade passing frequency. Then (t > 175 s) the pulses become weaker and there is a light increase in wind velocity.

This was one of the nights where a distinct beat was audible: a period with a distinct beat alternating with a period with a weaker or no beat, repeated more or less during the entire night. This pattern is compatible with a complex of three pulse trains with slightly different repetition frequencies of ca. 1 Hz. When the pulses are out of phase (around 150 s in figure V.4), the variations are 1 dB or less. When 2 of them are in phase (around 160 s) pulse height is doubled (+3 dB), and tripled (+5 dB, 170 s) when all three are in phase. The rotational speed of the turbines at the time was 20 rpm, so the repetition rate of blades passing a mast was 1 Hz.

The low number of pulse trains, compared to 17 turbines, is compatible with the fact that only a few turbines dominate the sound immission at this location. The calculated immission level is predominantly caused by two wind turbines (numbers 11 and 12: see figure IV.2, contributing 35% of the A-weighted sound energy), less by two others (9 and 14; 21%), so only 4 turbines contribute more than half of the sound immission energy.
In figure V.5 the equivalent 1/3 octave band spectrum at the façade of P has been plotted for the period of the beat (165 < t < 175 s in figure 6, spectra sampled at a rate of 20 s⁻¹), as well as the equivalent spectrum associated with the 5% highest (Lₐ₅ = 52.3 dB(A)) and the 5% lowest (Lₐ₉₅ = 47.7 dB(A)) broad band levels within this 10 s period, and the difference between both. As in the similar spectra in figure 4 we see that the beat corresponds to an increase at frequencies where trailing edge sound dominates: the sound pulses correspond to variations in 1/3 octave band levels at frequencies between 200 and 1250 Hz and are highest at 800 Hz. In figure V.5 also the equivalent 1/3 octave band levels are plotted for the period after beating where the wind was picking up slightly (t > 175 s in figure 6). Here spectral levels above 400 Hz are the same or slightly lower as on average at the time of beating, but at lower frequencies down to 80 Hz (related to in-flow turbulence) levels now are 1 to 2 dB higher. The increase in the ‘more wind’ spectrum at high frequencies (> 2000 Hz) is probably from rustling tree leaves.

Figure V.6 shows sound power spectra for a period with a distinct beat (150 < t < 175 s in figure 6), and a period with a weak or no beat (130 < t < 150 s). Each spectrum is an FFT of 0.2 Hz line width from broad band A-weighted immission sound pressure level values. The frequencies are therefore modulation, not sound frequencies. The spectra show that distinct beating is associated with higher total A-weighted levels at the blade passing frequency and its harmonics (k-fₘ with k = 1, 2, 3, ...). As has been shown above, the higher
level is related to the frequency range of trailing edge sound. Infrasound frequencies linked to thickness sound are negligible in total A-weighted sound levels. When beating is weaker but there is more wind \((t > 175 \text{ s})\), the level of the odd harmonics (base frequency \(k = 1\), and \(k = 3\)) is lower than during ‘beat’, whereas the first two even harmonics \((k = 2, 4)\) are equally loud, indicating more distorted (less sinusoidal) and lower level pulses. It is important to realize that the periodic variation as represented in figure V.6 is the result from a wind farm, not from a single turbine.

![Figure V.6: sound power spectrum of A-weighted broad band immission sound level at façade of dwelling P when beating is distinctly or not audible and with slightly increased wind speed. The ordinate spans 20 dB.](image)

In the long term measurements near the Rhede wind farm (see Chapter IV) average and percentile sound levels were determined over 5 minute periods. Periods where wind turbine sound was dominant could be selected with a criterion \((R_{bb,90} \leq 4 \text{ dB})\) implying a fairly constant source with less than 4 dB variation for 90% of the time. The statistical distribution of the values of \(R_{bb,90} = L_{A5} - L_{A95} \leq 4 \text{ dB}\) has been plotted in 1 dB intervals in figure V.7 for the two long term measurement locations A and B (see map in figure IV.1). Relative to dwellings P and R, location A (400 m from nearest turbine) is closer to the turbines, while location B (1500 m) is further away. Total measurement times –with levels in compliance with the criterion- were 110 and 135 hours, respectively. Figure V.7 shows that the criterion value \(R_{bb,90}\) (cut off at 4 dB) at both locations peaks at 2.5 dB.
Also plotted in figure V.7 is the value of $L_{A_{\text{max}}} - L_{A_{\text{eq}}}$ within 5 minute periods (while $R_{bb,90} \leq 4$ dB), peaking at 3.5 dB at both locations. Finally, the difference between maximum and minimum level within 5 minute periods, $R_{bb} = L_{A_{\text{max}}} - L_{A_{\text{min}}}$, peaks at 4.5 dB (location A) and 5.5 dB (B).

![Figure V.7: statistical distribution of level differences (in 1 dB-classes) between high and low sound levels within 5 minute periods at 400 m (left) and 1500 m (right) from the nearest wind turbine](image)

Where $R_{bb} > 7$ dB, the distributions are influenced by louder (non-turbine) sounds, such as from birds, causing a tail in the distributions at high levels. If we assume approximately symmetrical distributions without high level tails, the maximum range $L_{A_{\text{max}}} - L_{A_{\text{min}}} = R_{bb}$ due to the wind farm is 8.5 dB (location A) to 9.5 dB (B). This is 4 dB more than the prevailing difference at both locations.

### V.2.5 Summary of results

In table V.2 the level variations due to blade swish as determined in the previous sections have been summarised. Some values not presented in the text have been added. The ranges are presented as $R_{bb}$ and $R_{bb,90}$. The

---

1 in table in [Van den Berg 2005a] level variations close to the turbines were also given (as shown in figures V.2A-B); these values ($R_{bb} = 4.8$ dB close to turbine T1, 4.1 dB at T7 and 6.0 dB at T16) are not presented here as in fact these variations are not caused by the mechanism given in section V.1, but by other phenomena (see section II.2)
latter is of course a lower value as it leaves out high and low excursions occurring less than 10% of the time. The time interval over which these level differences occur differ: from several up to 16 minutes for the short term measurements, where wind conditions can be presumed constant, up to over 100 hours at locations A and B.

Table V.2: level variation in wind turbine sound due to blade swish, in dB

<table>
<thead>
<tr>
<th>location</th>
<th>Reference</th>
<th>atmospheric condition</th>
<th>(R_{bb})</th>
<th>(L_{A_{\text{max}}} - L_{A_{\text{min}}})</th>
<th>(R_{bb, 90})</th>
<th>(L_{A5} - L_{A95})</th>
</tr>
</thead>
<tbody>
<tr>
<td>Single turbine</td>
<td>Section V.1a</td>
<td>neutral</td>
<td>2 ± 1</td>
<td>(L_{A_{\text{max}}} - L_{A_{\text{min}}})</td>
<td>(R_{bb, 90})</td>
<td>(L_{A5} - L_{A95})</td>
</tr>
<tr>
<td>Single turbine</td>
<td>Section V.1a</td>
<td>stable</td>
<td>3.8 ± 1.7</td>
<td>(L_{A_{\text{max}}} - L_{A_{\text{min}}})</td>
<td>(R_{bb, 90})</td>
<td>(L_{A5} - L_{A95})</td>
</tr>
<tr>
<td>Single turbine</td>
<td>Section V.1a</td>
<td>very stable</td>
<td>4.8 ± 1.7</td>
<td>(L_{A_{\text{max}}} - L_{A_{\text{min}}})</td>
<td>(R_{bb, 90})</td>
<td>(L_{A5} - L_{A95})</td>
</tr>
<tr>
<td>N equidistant turbines</td>
<td>(very) stable</td>
<td>single +</td>
<td>(\text{10} \cdot \log N)</td>
<td>(L_{A_{\text{max}}} - L_{A_{\text{min}}})</td>
<td>(R_{bb, 90})</td>
<td>(L_{A5} - L_{A95})</td>
</tr>
</tbody>
</table>

Measured results

| Single turbine | dwelling Z | [ETSU 1996] | unspecified | \(< 3\) | \(5.9^{3}\) | \(4.3\) |
| Solar | dwelling P | section | stable | | | |
| Solar | dwelling P | section | long term, stable | \(4.5 \text{ (most frequent)}\) | \(8.5 \text{ (maximum)}\) | \(5.5 \text{ (most frequent)}\) | \(9.5 \text{ (maximum)}\) |

notes: 1) hub height 100 m, rotor diameter 70 m, 20 rpm; 2) probably neutral; 3) for this turbine \((H = 45 \text{ m, } D = 46 \text{ m, } 26 \text{ rpm, } V_r = 12 \text{ m/s})\) \(R_{bb} \leq 3.3 \text{ dB}\) was calculated

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V.3 Perception of wind turbine sound

In a review of literature on wind turbine sound Pedersen concluded that wind turbine noise was not studied in sufficient detail to be able to draw general conclusions, but that the available studies indicated that at relatively low levels wind turbine sound was more annoying than other sources of community noise such as traffic [Pedersen 2003]. In a field study by Pedersen and Persson Waye [2004] 8 of 40 respondents living in dwellings with (calculated) maximum outdoor immission levels of 37.5 - 40.0 dB(A) were very annoyed by the sound, and at levels above 40 dB(A) 9 of 25 respondents were very annoyed. The correlation between sound level (in 2.5 dB classes) and annoyance was significant (p < 0.001). In this field study annoyance was correlated to descriptions of the sound characteristics, most strongly to swishing with a correlation coefficient of 0.72 [Pedersen et al 2004]. A high degree of annoyance is not expected at levels below 40 dB(A), unless the sound has special features such as a low-frequency components or an intermittent character [WHO 2000]. Psychoacoustic characteristics of wind turbine sound have been investigated by Persson-Waye and Öhrström in a laboratory setting with naive listeners (students not used to wind turbine sound): the most annoying sound recorded from five different turbines were described as ‘swishing’, ‘lapping’ and ‘whistling’, the least annoying as ‘grinding’ and ‘low frequency’ [Persson Waye et al 2002]. People living close to wind turbines, interviewed by Pedersen et al. [2004], felt irritated because of the intrusion of the wind turbines in their homes and gardens, especially the swishing sound, the blinking shadows and constant rotation.

Our experience at distances of approx. 700 to 1500 m from the Rhede wind farm, with the turbines rotating at high speed in a clear night and pronounced beating audible, is that the sound resembles distant pile driving. When asked to describe the sound of the turbines in this wind farm, a resident compares it to the surf on a rocky coast. A resident living further away from the wind farm (1200 m) likens the sound to an ‘endless train’. Another resident near a set of smaller wind turbines, described the sound as that of a racing rowing boat (where rowers simultaneously draw, also creating a periodic swish). On the website of MAIWAG, a group of
citizens from villages near four wind farms in the south of Cumbria (UK),
the sound is described as ‘an old boot in a tumble dryer’, and also as
‘Whumph! Whumph! Whumph!’ (see text box in section III.4). Several
residents near single wind turbines remarked that the sound often changed
to clapping, thumping or beating when night falls: ‘like a washing
machine’. It is common in all descriptions that there is noise (‘like a nearby
motorway’, ‘a B747 constantly taking off’) with a periodic fluctuation
superimposed. In all cases the sound acquires this more striking character
late in the afternoon or at night, especially in clear nights and downwind
from a turbine.
Part of the relatively high annoyance level and the characterisation of wind
turbine sound as lapping, swishing, clapping or beating may be explained
by the increased fluctuation of the sound. Our results in table V.2 show
that in a stable atmosphere measured fluctuation levels are 4 to 6 dB for
single turbines, and in long term measurements (over many 5 minute
periods) near the Rhede wind farm fluctuation levels of approx. 5 dB are
common but may reach values up to 9 dB.
The level difference associated with an amplitude modulation (AM) factor
mf is:
\[
\Delta L = 20 \cdot \log((1+mf)/(1-mf))
\]
(V.2a)
The modulation factor mf is the change in sound pressure amplitude due to
modulation, relative to the average amplitude. For \( \Delta L < 9 \text{ dB} \) a good
approximation (±5%) is:
\[
mf = 0.055 \cdot \Delta L
\]
(V.2b)
Now when \( \Delta L \) rises from 3 dB, presumably a maximum value for a
daytime (unstable or neutral) atmosphere, to 6 dB, mf rises from 17% to
33%. For a maximum value of \( \Delta L = 9 \text{ dB} \), mf is 50%.
Fluctuations are perceived as such when the modulation frequencies are
less than 20 Hz. Human sensitivity for fluctuations is highest at \( f_{mod} = 4 \)
Hz, which is the frequency typical for rhythm in music and speech
[Zwicker et al 1999], and for frequencies of the modulated sound close to 1
kHZ. For wind turbines we found that a typical modulation frequency is 1 Hz, modulating the trailing edge sound that itself is at frequencies of 500 – 1000 Hz. So human sensitivity for wind turbine sound fluctuations is relatively high.

Fluctuation strength can be expressed in a percentage relative to the highest perceptible fluctuation strength (100%) or as an absolute value in the unit vacil [Zwicker et al 1999]. The reference value for the absolute fluctuation strength is 1 vacil, equalling a 60 dB, 1 kHz tone, 100% amplitude-modulated at 4 Hz [Zwicker et al 1999]. For an AM pure tone as well as AM broad band noise, absolute fluctuations strength is zero until $\Delta L \approx 3$ dB, then increases approximately linearly with modulation depth up to a value of 1 vacil. For a broad band noise level $L_A$ the fluctuation strength $F_{bb}$ can be written as [Zwicker et al 1999]:

$$F_{bb} = \frac{5.8(1.25\cdot mf-0.25)(0.05\cdot L_A - 1)}{(f_{mod}/5 \text{ Hz})^2 + (4 \text{ Hz}/f_{mod}) + 1.5} \text{ vacil}$$

(V.3a)

With typical values for wind turbine noise of $f_{mod} = 1$ Hz and $L_A = 40$ dB(A), this can be written as $F_{bb} = 1.31\cdot(mf-0.2)$ vacil or, when $\Delta L < 9$ dB:

$$F_{bb} = 0.072\cdot(\Delta L - 3.6) \text{ vacil}$$

(V.3b)

When $\Delta L$ increases from 3 to 5 dB, $F_{bb}$ increases from negligible to 0.1 vacil. For the high fluctuation levels found at locations A and B ($\Delta L = 8$ to 9 dB), $F_{bb}$ is 0.3 to 0.4 vacil.

It can be concluded that, in a stable atmosphere, the fluctuations in modern wind turbine sound can be readily perceived. As yet it is not clear how this relates to possible annoyance. However, the sound can be likened to the rhythmic beat of music: pleasant when the music is appreciated, but distinctly intrusive when the music is unwanted.

The hypothesis that these fluctuations are important, is supported by descriptions of the character of wind turbine sound as ‘lapping’, ‘swishing’, ‘clapping’, ‘beating’ or ‘like the surf’. Those who visit a wind
turbine in daytime will usually not hear this and probably not realise that the sound can be rather different in conditions that do not occur in daytime. This may add to the frustration of residents: “Being highly affected by the wind turbines was hard to explain to people who have not had the experiences themselves and the informants felt that they were not being believed” [Pedersen et al 2004]. Persson-Waye et al [2002] observed that, from five recorded different turbine sounds “the more annoying noises were also paid attention to for a longer time”. This supported the hypothesis that awareness of the noise and possibly the degree of annoyance depended on the content (or intrusive character) of the sound.

Fluctuations with peak levels of 3 – 9 dB above a constant level may have effects on sleep quality. The Dutch Health Council [2004] states that "at a given L_{night} value, the most unfavourable situation in terms of a particular direct biological effect of night-time noise is not, as might be supposed, one characterised by a few loud noise events per night. Rather, the worst scenario involves a number of noise events all of which are roughly 5 dB(A) above the threshold for the effect in question.” For transportation noise (road, rail, air traffic) the threshold for motility (movement), a direct biological effect having a negative impact on sleep quality, is a sound exposure level per sound event of SEL = 40 dB(A) in the bedroom [Health Council 2004]. The pulses in figure V.4 have SEL-values up to 50 dB(A), but were measured on the façade. With an open window facing the wind turbines indoor SEL-values may exceed the threshold level. In other situations this of course depends on distance to and sound power of the turbines and on the attenuation between façade and bedroom. It is not clear whether the constant and relatively rapid repetition of wind turbine sound beats will have more or less effect on sleep quality, compared to vehicle or airplane passages. Pedersen and Persson Waye [2004] found that at dwellings where the (outdoor) sound level due to wind turbines exceeded 35 dB(A), 16% of 128 respondents reported sleep disturbance by this sound, of whom all but two slept with a window open in summer.
V.4 Conclusion

Atmospheric stability has a significant effect on the character of wind turbine sound. The change in wind profile causes a change in angle of attack on the turbine blades. This increases the thickness (infra)sound level as well as the level of trailing edge (TE) sound, especially when a blade passes the tower. TE sound is modulated at the blade passing frequency, but it is a high frequency sound, well audible and indeed the most dominant component of wind turbine noise. The periodic increase in sound level dubbed blade swish, is a well known phenomenon. Less well known is the fact that increasing atmospheric stability creates greater changes in the angle of attack over the rotor plane that add up with the change near the tower. This results in a thicker turbulent TE boundary layer, in turn causing a higher swish level and a shift to somewhat lower frequencies. It can be shown theoretically that for a modern, tall wind turbine in flat, open land the angle of attack at the blade tip passing the tower changes with approx. 2° in daytime, but this value increases with 2° when the atmosphere becomes very stable. The calculated rise in sound level during swish then increases from 2 dB to 5 dB. This value is confirmed by measurements at single turbines in the Rhede wind farm where maximum sound levels rise 4 to 6 dB above minimum sound levels within short periods of time.

Added to this, atmospheric stability involves a decrease in large scale turbulence. Large fluctuations in wind velocity (at the scale of a turbine) vanish, and the coherence in wind velocity over distances as great as or larger than the size of an entire wind farm increases. As a result turbines in the farm are exposed to a more constant wind and rotate at a more similar speed with less fluctuations. Because of the near-synchronicity, blade swishes may arrive simultaneously for a period of time and increase swish level. The phase difference between turbines determines where this amplification occurs: whether the swish pulses will coincide at a location depends on this phase difference and the propagation time of the sound. In an area where two or more turbines are comparably loud the place where this amplification occurs will sweep over the area with a velocity determined by the difference in rotational frequency. The magnitude of this effect thus depends on stability, but also on the number of wind turbines.
and the distances to the observer. This effect is in contrast to what was expected, as it seemed reasonable to suppose that turbines would behave independently and thus the blade swish pulses from several turbines would arrive at random, resulting in an even more constant level than from one turbine. Also, within a wind farm the effect may not be noticed, since comparable positions in relation to two or more turbines are less easily realised at close distances and the position relative to a turbine rotor is quite different.

Sound level differences $L_{A_{\max}}-L_{A_{\min}}$ (corresponding to swish pulse heights) within 5 minute periods over long measurement periods near the Rhede wind farm show that level changes of approximately 5 dB occur for an appreciable amount of time and may less often be as high as 8 or 9 dB. This level difference did not decrease with distance (from 400 m to 1500 m). The added 3-5 dB, relative to a single turbine, is in agreement with simultaneously arriving pulses from two or three approximately equally loud turbines.

The increase in blade swish level creates a new percept, fluctuating sound, that is absent or weak in neutral or unstable atmospheric conditions. Blade passing frequency is now an important parameter as a modulation frequency (not as an infrasound frequency). Human perception is most sensitive to modulation frequencies close to 4 Hz of sound with a frequency of approximately 1 kHz. The hypothesis that fluctuations are important is supported by descriptions given by naïve listeners as well as residents: turbines sound like ‘lapping’, ‘swishing’, ‘clapping’, ‘beating’ or ‘like the surf’. It is not clear to what degree this fluctuating character determines the relatively high annoyance caused by wind turbine sound and to a deterioration of sleep quality. Further research is necessary into the perception and annoyance of wind turbine sound, with correct assumptions on the level and character of the sound. Also the sound exposure level of fluctuations in the sound in the bedroom must be investigated to be able to assess the effects on sleep quality.
VI STRONG WINDS BLOW UPON TALL TURBINES: wind statistics below 200 m altitude

VI.1 Atmospheric stability in wind energy research

In the European Wind Atlas model (‘Wind Atlas Analysis and Application Program’ or WAsP) [Troen et al 1989] wind energy available at hub height is calculated from wind velocities at lower heights. The Atlas states that “modifications of the logarithmic wind profile are often neglected in connection with wind energy, the justification being the relative unimportance of the low wind velocity range. The present model treats stability modifications as small perturbations to a basic neutral state.” With the increase of wind turbine heights this quote is now an understatement. In recent years atmospheric stability is receiving gradually more attention as a determinant in wind energy potential, as demonstrated by a growing number of articles on stability related wind profiles in different types of environments such as Danish offshore sites [Motta et al 2005], the Baltic Sea [Smedman et al 1996], a Spanish plateau [Pérez et al 2005] or the American Midwest [Smith et al 2002]. Recently Archer and Jakobsen [2003] showed that wind energy potential at 80 m altitude in the contiguous US ‘may be substantially greater than previously estimated’ because atmospheric stability was not taken into account: on average 80-m wind velocities appear to be 1.3 – 1.7 m/s higher than assumed from 10-m extrapolated wind velocities in a neutral atmosphere.

VI.2 The Cabauw site and available data

To investigate the effect of atmospheric stability on wind, and thence on energy and sound production, data from the meteorological research station of the KNMI (Royal Netherlands Meteorological Institute) at Cabauw in the western part of the Netherlands were kindly provided by dr Bosveld of the KNMI. The site is in open pasture for at least 400 m in all directions. Farther to the west the landscape is open, to the distant east are trees and low houses. More site information is given in [KNMI 2005, Van Ulden et al 1996]. The site is considered representative for the flat western and
northern parts of the Netherlands. These in turn are part of the low-lying plain stretching from France to Sweden. Meteorological data are available as half hour averages over several years. Here data of the year 1987 are used. Wind velocity and direction are measured at 10, 20, 40, 80, 140 and 200 m altitude. Cabauw data are related to Greenwich Mean Time (GMT); in the Netherlands the highest elevation of the sun is at approximately 12:40 Dutch winter time, which is 20 minutes before 12:00 GMT.

An indirect measure for stability is Pasquill class, derived from cloud cover, wind velocity and position of sun (above or below horizon). Classes range from A (very unstable: less than 50% clouding, weak or moderate wind, sun up) to F (moderately to very stable: less than 75% clouding, weak or moderate wind, sun down). Pasquill class values have been estimated routinely at Dutch meteorological stations [KNMI 1972].

**VI.3  Reference conditions**

To relate the meteorological situation to wind turbine performance, an 80 m hub height wind turbine with three 40 m long blades will be used as reference for a modern 2 to 3 MW, variable speed wind turbine. To calculate electrical power and sound power level, specifications of the 78 m tall Vestas V80 – 2MW wind turbine will be used. For this turbine cut-in
(hub height) wind velocity is 4 m/s, and highest operational wind velocity 25 m/s.

Most data presented here will refer to wind velocity at the usual observation height of 10 m and at 80 m hub height. Wind shear will be presented for this height range as well as the range 40 to 140 m where the rotor is. The meteorological situation is as measured in Cabauw in 1987, with a roughness height of 2 cm. The year will be divided in meteorological seasons, with spring, summer, autumn and winter beginning on the first day of April, July, October and January, respectively.

We will consider four classes of wind velocity derived from Pasquill classes A to F and shown in table 1: unstable, neutral, stable and very stable. In table VI.1 (the same as table III.1, but written slightly different to show boundaries between stability classes in terms of m) this is also given in terms of the shear exponent, but this is tentative as there is no fixed relation between Pasquill classification and shear exponent or stability function Ψ. This classification is in agreement with that in chapter III, though there typical mid-class values of m were given, not values at the boundaries between classes. In our reference situation ‘very stable’ (m > 0.4) corresponds to a Monin-Obukhov length 0 < L < 100 m, ‘stable’ (0.25 < m < 0.4) refers to 100 m < L < 400 m, near neutral to |L| > 400 m.

This is somewhat different from the Monin-Obukhov length based classification used by Motta et al [2005] for a coastal/marine environment. Motta et al qualified 0 < L < 200 m as very stable, 200 m < L < 1000 m as stable and |L| > 1000 m as near-neutral, so they considered a wider range of conditions as (very) stable when compared to table 1.

<table>
<thead>
<tr>
<th>Pasquill class</th>
<th>name</th>
<th>shear exponent</th>
</tr>
</thead>
<tbody>
<tr>
<td>A – B</td>
<td>(very – moderately) unstable</td>
<td>m ≤ 0.21</td>
</tr>
<tr>
<td>C</td>
<td>near neutral</td>
<td>0.21 &lt; m ≤ 0.25</td>
</tr>
<tr>
<td>D – E</td>
<td>(slightly – moderately) stable</td>
<td>0.25 &lt; m ≤ 0.4</td>
</tr>
<tr>
<td>F</td>
<td>very stable</td>
<td>0.4 &lt; m</td>
</tr>
</tbody>
</table>
VI.4 Results: wind shear and stability

VI.4.1 Wind velocity shear

In figure VI.2 the average wind velocities at altitudes of 10 m to 200 m are plotted versus time of day. Plotted are averages per half hour of all appropriate half hours in 1987. As figure VI.2 shows, the wind velocity at 10 m follows the popular notion that wind picks up after sunrise and abates after sundown. This is obviously a ‘near-ground’ notion as the reverse is true at altitudes above 80 m. Figure VI.2 helps to explain why this is so: after sunrise low altitude winds are coupled to high altitude winds due to the vertical air movements caused by the developing thermal turbulence. As a result low altitude winds are accelerated by high altitude winds that in turn are slowed down. At sunset this process is reversed. In figure VI.2 also the wind velocity $V_{80}$ is plotted as calculated from the measured wind velocity $V_{10}$ with equation III.3 ($z_0 = 2$ cm, equivalent to equation III.1 with $m = 0.14$), as well as the shear exponent $m$ calculated with equation III.4. The logarithmically extrapolated $V_{80}$ approximates actual $V_{80}$ in daytime when the shear exponent has values close to 0.14. However, the prediction is very poor at night time, when $m$ rises to a value of 0.3, indicating a stable atmosphere.

![Figure VI.2: solid lines, bottom to top: 1987 wind velocity per clock hour at heights 10 to 200 m; dotted line: logarithmically extrapolated $V_{80}$; $+$: shear exponent $m_{10,80}$]
For the hourly progress of wind velocities large deviations from the average wind profile occur. This is illustrated in figure VI.3 for a week in winter and a week in summer with measured $V_{10}$ values and measured as well as logarithmically extrapolated $V_{80}$ values. In the winter week in January 1987 ground and air were cold for a long time (below freezing point) with very little insolation. Temperature varied from night to day (diurnal minimum to maximum) with 7 °C on the first day and 5 °C or less on the next days, and the atmosphere was close to neutral with measured $V_{80}$ more or less equal to the extrapolated $V_{80}$. In the summer week in July 1987 there was little clouding after the first two days; insolation was strong in daytime, and nights were 10 to 14 °C cooler than days, resulting in a stable to very stable night time atmosphere. Here, night time wind velocity was rather higher than predicted with the logarithmic wind profile.

In figure VI.4 wind velocities per half hour are again plotted for different heights, as in figure VI.2, but now averaged per clock half hour and per meteorological season. In spring and summer differences between night
and day seem more pronounced than in autumn or winter. In fall and winter wind velocities are on average higher.

Figure VI.4: wind velocity per hour GMT at heights of 10, 20, 40, 80, 140 and 200 m (bottom to top; 80 m is bold) in the meteorological seasons in 1987
In figure VI.5 the frequency distribution is plotted of the half-hourly wind velocities at five different heights. Also plotted is the distribution of wind velocity at 80 m as calculated from the 10-m wind velocity with the logarithmic wind profile (equation III.3, \( m = 0.14 \)). Wind velocity at 80 m has a value of \( 7 \pm 2 \) m/s for 50% of the time. For the logarithmically extrapolated wind velocity at 80 m this is \( 4.5 \pm 2 \) m/s.

In figure VI.6 the prevalence of the shear exponent in the four meteorological seasons is plotted, determined from the half-hourly 10-m and 80-m wind velocities. It shows that, relative to autumn and winter, a neutral or mildly stable atmosphere occurs less often in spring and summer, whereas an unstable as well as—in summer— a very stable atmosphere occurs more often. As summer nights are short this means that a relatively high percentage of summer night hours has a stable atmosphere.

**Figure VI.5**: distribution of measured wind velocities at 10, 40, 80, 140 and 200 m; dashed line: \( V_{80} \) extrapolated from \( V_{10} \)

**Figure VI.6**: distribution of shear exponent per meteorological season, determined from \( V_{80}/V_{10} \)
VI.4.2 Shear and ground heat flux

Figure VI.7 shows how the shear exponent depends on the total heat flow to the ground for two different height ranges: 10 – 80 m in the left panel, 40 – 140 m in the right panel. The shear exponent is calculated from the wind velocity ratio with equation III.1. The heat flow at Cabauw is determined from temperature measurements at different heights, independent of wind velocity. Total heat flow is the sum of net radiation, latent and sensible heat flow, and positive when incoming flow dominates. For heat flows above approximately 200 W/m² the shear exponent m is between 0 and 0.21, corresponding to an unstable atmosphere, as expected. For low or negative (ground cooling) heat flows the range for m increases, extending from -1 up to +1.7. These values include conditions with very low wind velocities. If low wind velocities at 80 m height (V_{80} < 4 m/s, occurring for 19.7% of the time) are excluded, m_{10,80} varies (with very few exceptions) between 0 and 0.6, and m_{40,140} varies between -0.1 and +0.8. A negative exponent means wind velocity decreases with height. The data show that below 80 m this occurs in situations with little wind (V_{80} < 4 m/s), but at greater heights also at higher wind velocities. In fact, V_{140} was lower than V_{80} for 7.5% of all hours in 1987, of which almost half (3.1%)
when $V_{80}$ was over 4 m/s. Such a decrease of wind velocity with height occurs at the top of a ‘low level jet’ or nocturnal maximum; it occurs at night when kinetic energy of low altitude air is transferred to higher altitudes.

For $V_{80} > 4$ m/s both shear exponents ($m_{10,80}$ and $m_{40,140}$) are fairly strongly correlated (correlation coefficient 0.85), showing that generally there is no appreciable change between both altitude ranges. For low wind velocities ($V_{80} < 4$ m/s) both shear exponents are less highly correlated (correlation coefficient 0.62).

### VI.4.3 Wind direction shear

When stability sets in the decoupling of layers of air also affects wind direction. The higher altitude wind more readily follows geostrophic wind and therefore can change direction when stability sets in, while lower altitude winds are still influenced by the surface following the earth’s rotation. In the left panel of figure VI.8 the change in wind direction at 80 m relative to 10 m is plotted as a function of the shear exponent as a measure of stability. A positive change means a clockwise change (veering wind) at increasing altitude. The right panel shows the wind direction change from 40 to 140 m as a function of the shear exponent determined from the wind velocities at these heights. In both cases the prevailing change from $m = 0$ to $m = 0.5$ is $30^\circ$, but with considerable variation.

![Figure VI.8: wind direction change between 10 and 80 m (left) and 40 and 140 m (right) vs. shear exponent $m$ between same heights for $V_{80} > 4$ m/s](image-url)
VI.4.4 Prevalence of stability

In figure VI.9 the percentages are given that the atmosphere is very stable, stable, neutral and unstable respectively (as defined in table VI.1) for 1987 as a whole and per meteorological season. Prevalence is given for heights from 10 and 80 m (upper panel figure VI.9) and for heights from 40 to 140 m (lower panel). The upper panel is in fact a summation over the four ranges of the shear exponent indicated in figure VI.6. It appears that in autumn the atmosphere is most often stable, and least often unstable. In spring the opposite is true: instability occurs more often than stability. Overall the atmosphere up to 80 m is unstable (m < 0.21) for 47% of the time and stable (m > 0.25) for 43% of the time. At higher altitudes (40 to 140 m) percentages are almost the same: 44% and 47%, respectively. This means that for most of the daytime hours the atmosphere is unstable, and for most of the night time hours stable. For the rest (9 to 10%) of the time the atmosphere is near neutral.

Climatological observations can put the Cabauw data in national perspective. In figure VI.10 the prevalence of Pasquill classes E and F (corresponding to approximately m > 0.33) are given as observed at 12 meteorological stations all over the Netherlands over the period 1940 - 1970 [KNMI 1972], ordered according to yearly prevalence. Three of the
dunes on the North Sea coast, Vlissingen is at the Westerschelde estuary and Den Helder is on a peninsula between the North Sea and the Wadden Sea. At Den Helder a stable atmosphere occurs for only 8% of the time per year, whereas at both other coastal stations this is 13% to 16% and at the other landward stations 15% to 20% of the time. At Cabauw a value of $m > 0.33$ occurs for 27% of the time.

![Figure VI.10: prevalence of observed stability (Pasquill classes E and F) per season and per year at 12 different Dutch stations over 30 years (data from [KNMI])](image)

**VI.5. Results: effects on wind turbine performance**

**VI.5.1 Effect on power production**

The effect of atmospheric stability can be investigated by applying the Cabauw data to a reference wind turbine, the Vestas V80-2MW [Vestas 2003, Jorgensen 2002]. This turbine has an ‘Optispeed’ sound reduction possibility to reduce sound power level (by adapting the speed of the rotor and generator). We will present data for the highest (‘105.1dB(A)’) and lowest (‘101.0dB(A)’) sound power curve. To calculate the electric power $P_{80}$ as a function of wind velocity $V_h$ at hub height the factory ‘105.1dB(A)’ highest power (‘hp’) curve is approximated with a fourth power polynome:
\[ P_{h,\text{hp}} = 0.0885 \cdot V_h^4 - 8.35 \cdot V_h^3 + 186 \cdot V_h^2 - 1273 \cdot V_h + 2897 \text{ kW} \quad (\text{VI.1a}) \]

which is valid for \(4 < V_h < 14.3\) m/s. In figure VI.11 this fitted curve is plotted as diamonds on top of the manufacturer’s specification [Vestas 2003]. For higher wind velocities (>14.3 m/s; 2% of time at Cabauw) electric power is constant at 2000 kW, for lower wind velocities (< 4 m/s; 20% of time) electric power is set to zero.

A fourth power relation is used as this is convenient to fit the power curve at 12 m/s where maximum power is approached. For lower wind velocities \((V_h < 11\) m/s) the power curve can be fitted with a third power \((P_h = 1.3 \cdot V_h^3)\) in agreement with the physical relation between wind power and wind velocity.

![V80-2.0 MW power curves](image)

**Figure VI.11:** lines: Vestas V80 power curves vs. hub height wind speed; diamonds: best fit to 105.1 dB(A) curve (figure adapted from [Vestas 2003])

Electric power can thus be calculated from real wind velocities as measured each half hour at 80 m height, or from 80-m wind velocities logarithmically extrapolated from wind velocity at 10 m height. The result is plotted in figure VI.12 as an average power versus time of day \(P_{80,\text{hp}}\) (the power averages are over all hours in 1987 at each clock hour). Actual power production appears to be more constant than estimated with extrapolations from 10-m wind velocities. When using a logarithmic extrapolation, daytime power production is overestimated, while night time power production is underestimated. The all year average is plotted with large symbols at the right side of the graph in figure VI.12: 598 kW when based on measured wind velocity or a 30% annual load factor, 495 kW
when based on extrapolated wind velocity or a 25% load factor. In figure VI.12 also the wind power is plotted when the turbine operates in the lowest ‘101.0dB(A)’ power curve (‘lp’) where the best fit is:

\[ P_{h,lp} = 0.089 \cdot V_h^4 + 0.265 \cdot V_h^3 + 43 \cdot V_h^2 - 326 \cdot V_h + 749 \text{ kW} \]  

(VI.1b)

The year average is now 569 kW, corresponding to a 28% annual load factor. The 4 dB lower sound level setting thus means that yearly power production has decreased to a factor 0.94.

In the calculations it was implicitly assumed that the wind velocity gradient over the rotor was the same as at the time the power production was determined as a function of hub height wind velocity. In stable conditions however, the higher wind gradient causes a non-optimal angle of attack at the blade tips when the tips travel far below and above the hub. This will involve some loss, which is not determined here.

**VI.5.2 Effect on sound production**

Figure VI.13 shows ‘theoretical’ sound power levels for the Vestas turbine [Vestas 2003, Jorgensen 2002]; in fact for \( V_h < 8 \) m/s measured levels are somewhat lower, for \( V_h > 8 \) m/s somewhat higher [Jorgensen 2002]. To calculate the sound power level \( L_W \) as a function of hub height wind velocity \( V_h \) the factory ‘105.1dB(A)’ high power curve is approximated with a fourth power polynomial:

\[ L_{W,hp} = -0.0023 \cdot V_h^4 + 0.146 \cdot V_h^3 - 2.82 \cdot V_h^2 + 22.6 \cdot V_h + 39.5 \text{ dB(A)} \]  

(VI.2a)
for $4 < V_h < 12$ m/s and $L_{W,lp} = 107$ dB(A) for $V_h > 12$ m/s. In figure VI.14 the result per clock hour is plotted when using actual and extrapolated (from 10 m) wind velocities. Averaged over all 1987 the sound power level in daytime is overestimated by 0.5 dB, but at night underestimated by 2 dB. In the ‘101.0dB(A)’ low power curve setting the best fourth power polynomial fit is (in figure VI.13 plotted as diamonds over the Vestas curve):

$$L_{W,lp} = -0.022 \cdot V_h^4 + 0.78 \cdot V_h^3 - 10 \cdot V_h^2 + 55.3 \cdot V_h - 12.3 \text{ dB(A)}$$  \hspace{1cm} (VI.2b)

for $4 < V_h < 12$ m/s and $L_{W,lp} = 105$ dB(A) for $V_h > 12$ m/s. The sound power levels in this setting are, for $6 < V_h < 12$ m/s, on average 3 dB lower than in the high power setting.
The differences between actual and logarithmically predicted sound power levels can be bigger than the over one year hourly averaged values in figure VI.14 show. This is illustrated in figure VI.15 for two days each in January and July 1987 (also shown in figure VI.3) where actual and predicted half-hour sound power levels are plotted as a function of 10-m wind velocity. On both winter days actual sound power agrees within 1 dB with the predicted sound power for wind velocities $V_{10} > 5.5$ m/s; at lower 10-m wind velocities actual levels are rather higher for most of the time. On both summer days the 10-m wind velocities are lower than in winter, and sound power level now is more often higher than predicted and can reach near maximum levels even at very low (2.5 m/s) 10-m wind velocities (when at ground level people will probably feel no wind at all). In these conditions residents in a quiet area will perceive the highest contrast: hardly or no wind induced sound in vegetation, while the turbine(s) are rotating at almost top speed. In these conditions also an increased fluctuation strength of the turbine sound will occur (see chapter V), making the sound more conspicuous.

![Figure VI.15: half-hourly progress of actual (grey diamonds) and logarithmically predicted (black dots) sound power level plotted vs. 10-m wind speed over 48 hours; left: January 13-14; right: July 2-3](image-url)
VI.6 Other onshore results

Values of wind shear have been reported by various authors, showing similar results. Pérez et al [2005] measured wind velocities up to 500 m above an 840 m altitude plateau north of Valladolid, Spain, for every hour over sixteen months. The shear exponent, calculated from the wind velocity at 40 m and 220 m, varied from 0.05 to 0.95, but was more usual between 0.1 and 0.7. High shear exponents occurred more often than in Cabauw: $m > 0.48$ for 50% of the time. This is likely the result of the more southern position: insolation is higher, causing bigger temperature differences between day and night, and the atmosphere above the plateau is probably drier causing less reflection of outward infrared radiation at night. There was a distinct seasonal pattern, with little day-night differences in January, and very pronounced differences in July.

Smith et al [2002] used data from wind turbine sites in the US Midwest over periods of 1.5 to 2.5 years and calculated shear exponents for wind velocities between a low altitude of 25 - 40 m and a high altitude of 40 – 123 m. At four sites the hourly averaged night time (22:00 – 6:00) shear exponent ranged from 0.26 to 0.44, in daytime from 0.09 to 0.19. The fifth station (Ft. Davis, Texas) was exceptional with a day and night time wind shear below 0.17 and a very low day time wind shear ($m = 0.05$).

Archer et al [2003] investigated wind velocities at 10 m and 80 m from over 1300 meteorological stations in the continental USA. No shear statistics are given, but for 10 stations the ratio $V_{80}/V_{10}$ is plotted versus time of day. At all these stations the ratio is $1.4 \pm 0.2$ in most of the daytime and $2.1 \pm 0.3$ in most of the night time. Using equation III.4, it follows that the shear exponent has a value of $0.15 \pm 0.07$ and $0.35 \pm 0.07$, respectively.

At the 2005 Berlin Conference on Wind Turbine Noise two presentations added to these wind shear data, now (also) from a noise perspective. Harders et al [2005] showed hourly wind velocity averaged over the year 2000 at altitudes between 10 and 98 m from the Lindenberg Observatory near Berlin. The results are very much like those in figure VI.2, with a wind velocity ratio $V_{80}/V_{10} = 1.3$ at noon, increasing to 1.9 in night time.
hours. This corresponds to an average shear exponent of 0.13 and 0.3, respectively.

Botha [2005] presented results from 8 to 12 months measurements at sites in two flat Australian areas and two sites in more complex (non flat) New Zealand terrain. On the Australian sites the average day time wind velocity ratio $V_{80}/V_{10}$ was 1.5, in night time 1.7 and 1.8. This corresponds to shear exponents of to 0.19 and 0.26 to 0.28, respectively. In the hilly New Zealand areas the average wind velocity ratio was between 1.2 and 1.25 in day as well as night time, from which the shear exponent can be calculated as 0.1.

From the measurements at the Rhede wind farm the shear exponent could be calculated from the 10-m and 100-m wind velocity, the latter determined from the sound level and the relation between sound power level and hub height (100 m) wind velocity. This was done for all (892) five minute periods when wind turbine sound was dominant between 23:00 and 04:00 hours within the measurement period (May and June; location A in figure IV.2). From the Cabauw data the same period and time was selected and all values of the half-hour shear exponent $m_{10,80}$ were determined. For both locations the resulting frequency distributions of the shear exponent are plotted in figure VI.16. The distributions are rather similar and show that a stable atmosphere ($m > 0.25$) occurred for over 95% of the time in night time hours (23 – 4 o’clock) in spring (May – June) at Cabauw as well as at Rhede.

![stability in May-June 23 - 04 hours](image)

**Figure VI.16: frequency distribution of the shear exponent at Cabauw and in the measurement period near the Rhede wind farm in the same period of time**
VI.7 Conclusion

Results from various landward areas show that the shear exponent in the lower atmospheric boundary layer (< 200 m) in daytime is 0.1 to 0.2, corresponding to a wind velocity ratio \( V_{80}/V_{10} \) of 1.25 to 1.5. The associated wind profile is comparable to the profile predicted by the well-known logarithmic wind profile for low roughness lengths (low vegetation). At night the situation is quite different and in various landward areas the shear exponent has a much wider range with values up to 1, but more usually between 0.25 and 0.7. Near the Rhede wind farm the same range of wind shear occurred, showing that the site indeed was suitable to study the effect of atmospheric stability on wind turbine performance and representative for many other locations. A shear exponent 0.25 < \( m < 0.7 \) means that the ratio \( V_{80}/V_{10} \) varies between 1.7 and 4.3. High altitude wind velocities are thus (much) higher than expected from logarithmic extrapolation of 10-m wind velocities.

A high wind shear at night is very common and must be regarded a standard feature of the night time atmosphere in the temperate zone and over land. In fact the atmosphere is neutral for only a small part (approximately 10%) of the time. For the rest it is either stable (sun down) or unstable (sun up).

As far as wind power concerns, the underestimate of high altitude night time wind velocity has been compensated somewhat by the overestimate of high altitude daytime wind velocity. This may partly explain why, until recently, atmospheric stability was not recognized as an important determinant for wind power.

To assess wind turbine electrical and sound power production the use of a neutral wind profile should be abandoned as it yields data that are not consistent with reality.
VII THINKING OF SOLUTIONS: measures to mitigate night time wind turbine noise

VII.1 Meeting noise limits

Sound from modern wind turbines is predominantly the result of turbulence on the blades; reduction of this source is the topic of dedicated research, such as the SIROCCO (Silent rotors by acoustic optimisation) program which seeks to improve the design of the wind turbine blade; in the near future a reduction of approximately 2 dB might be achieved [Schepers et al 2005]. Sound reduction by reducing blade speed is an option already available in modern turbines.

In this chapter we will deal with the (‘added’) sound produced by a wind turbine due to increased atmospheric stability. To address this problem two types of mitigation measures can be explored:

1. reduce the sound level down to the pertinent (legal) limit for environmental noise;
2. reduce the level variations due to blade swish/beating.

The first measure of course must be pursued as it is a legal obligation. The need for reduction depends on the type of limit. E.g., in Germany the limit applies to the maximum sound immission level (the level produced at nominal maximum power), regardless of wind velocity as such. In many countries the limit is based on the wind velocity related background ambient sound level ($L_{95}$ or $L_{90}$). In the UK and elsewhere the limit is a constant at low 10-m wind velocities and 5 dB above ambient background level ($L_{90} + 5$ dB) at higher 10-m wind velocities. In the Netherlands the standard limit is a reference curve constructed from a constant value at low 10-m wind velocities and a wind velocity dependent part at higher 10-m wind velocities (see figure VII.1). For wind farms over 15 MW other limit values may apply, and local authorities may enforce other limits in ‘non-standard’ local conditions.
In assessments of wind turbine noise immission the effect of atmospheric stability has usually been disregarded and the 10-m wind velocity was erroneously used for all atmospheric conditions. In that case high sound levels only occur at high wind velocities and this can be accommodated by limit values as in figure VII.1. In reality however these limits are not always met as high immission sound levels already occur at a lower 10-m wind velocity. This implies that an extra effort to reduce the immission level may be necessary.

In hilly and certainly in mountainous terrain this change in wind profile may be influenced or even overridden by relief related changes. For example: in a valley a down flowing (decelerating) wind may enhance the effect of stability, whereas an up flowing (accelerating) wind may compensate the effect of stability. Furthermore the wind profile as well as the temperature profile will simultaneously influence the propagation paths of sound. Combined effects are therefore complex and, though readily understood qualitatively, not easily predicted quantitatively.

The second measure is worth considering when the noise limit incorporates a penalty for a sound having a distinctive (impulsive or fluctuating) character. In that case either the sound immission level should be reduced by a value equal to the penalty (usually 5 dB) or the sound character must change.

**VII.2 Reduction of sound level**

When the sound immission level is limited to a value depending on the 10-m wind velocity or the (supposedly 10-m wind velocity dependent) ambient sound level, the problem is that hub height wind velocity is not uniquely related to 10-m wind velocity and the sound emission as well as immission level can have a range of levels depending on atmospheric...
stability. The turbine thus operates at hub height wind velocity, but must be controlled by a 10-m based wind velocity. To decrease the sound level from a given turbine the speed of rotation can be decreased, either by directly changing blade pitch or indirectly by changing the mechanical load (torque) on the rotor. This implies a lower efficiency at the turbine as the tip speed ratio $\Omega R/V_o$ will decrease and deviate from its value optimized for produced power. It is necessary to find a new optimum that also takes noise production into account.

VII.2.1 Wind velocity controlled sound emission

As a result of opposition to wind farm proposals in the relatively densely populated central province of Utrecht in the Netherlands all proposals were cancelled but one. The exception is in Houten (incidentally 8 km east of Cabauw; see previous chapter), where the local authorities want to stimulate wind energy by allowing the constructing of several 3 MW turbines, at the same time ensuring that residents will not be seriously annoyed. Atmospheric stability is taken into account by not accepting the usual logarithmic relation between 10-m and hub height wind velocity. The official permission will require that the immission sound level at specified locations must not exceed the background level of all existing ambient sound. Of course ambient sound level depends on wind velocity if the wind is sufficiently strong, but in this area it also depends on wind direction as that determines audibility of distant sources: a motorway to the west, the town to the north-east and relatively quiet agricultural land to the south-east. So the ambient background level, measured as $L_{95}$, must be measured in a number of conditions: as a function of wind velocity (1 m/s classes), wind direction (4 quadrants) and time of day (day, evening, night). These values equal the limit values for the immission level $L_{imm}$, and from this it can be calculated what the maximum allowable sound power level $L_{W_{max}}$ per turbine is at every condition, presuming all (or perhaps a selection of) turbines produce. It is advisable to determine wind characteristics and turbine performance over a period of at least five minutes, as wind velocity variations are relatively strong at frequencies above approximately 3 mHz (inverse of 5 min) and weak at lower frequencies down to the order of 0.1 mHz (inverse of several hours) [Wagner et al 1996]. On the other hand it is
desirable to adapt to changing conditions, so averaging over 5 minutes seems a good choice.

Control will thus be achieved in a number of steps:
1. measure wind direction $D_{10}$ and wind velocity $V_{10}$ in open land over a 5-minute period; from this determine the ambient back ground level from the previously established relation $L_{95}(D_{10}, V_{10})$.
2. determine the limit value for the sound power level $L_{W_{\text{max}}}$ from the previously established relation $L_{\text{imm}}(L_w)$; the limit value is determined by $L_{\text{imm}} = L_{95}$.
3. determine the actual sound power level $L_{W,5\text{min}}$ from wind turbine performance (electric power or speed);
4. if $L_{W,5\text{min}} > L_{W_{\text{max}}}$ (equivalent to $L_{\text{imm}} > L_{95}$) the control system must decrease sound power level for the next period; if $L_{W,5\text{min}} < L_{W_{\text{max}}}$ the reverse applies (until maximum speed is attained).

The pro’s of this control system are that it is straightforward, simple, easy to implement and directly related to existing Dutch noise limits. However, it is based on the assumption that $L_{95}$ depends on three parameters only: wind velocity, wind direction and diurnal period (day, evening, night). In reality background level will also vary within a diurnal period (e.g. traffic: nights are very quiet at around 4 AM and most busy just before 7 AM), and it will depend on the day of the week (e.g. Sunday mornings are quieter than weekday mornings), the season (vegetation, holidays), the degree of atmospheric stability (no wind in low vegetation in stable conditions, even when 10-m wind velocity is several m/s) and other weather conditions such as rain. Also sound immission from distant sources will differ with weather conditions.

Measurements show that indeed 10-m wind velocity is not a precise predictor of ambient sound level. These measurements were performed from June 9 through June 20, 2005 at two locations: wind velocity was measured at 10-m height in open terrain, at least 250 m from any obstacles over 1 m height (trees lining the busy and broad Amsterdam-Rhine Canal to the northeast) and over 1000 m from obstacles in any other direction; the
sound level was measured close to a farm next to the canal (see figure VII.2). Total measurement time was 220 hours.

Some results are plotted in figure VII.3: $L_{95}$ per 5-minute period as a function of wind velocity, separately for two opposite wind directions (left and right panel) and two periods (black and blue markers). The periods are night (23 PM – 7 AM) and day (7 AM – 7 PM), the wind directions southeast (90° - 180° relative to north) and northwest (270° - 360°), where respectively the lowest and highest ambient levels were expected. The northwest data total 675 5-minute periods or 26% of all measurement time, the southeast data cover 511 periods or 19% of the measurement time.
The values of $L_{95,5\text{min}}$ are calculated from all (300) 1-second samples of the sound pressure level within each 5-minute period, wind velocity is the average value of all 1-second samples of the wind velocity. To determine a long-term background level an appropriate selection (wind direction, period) of all measured 1-second sound levels can be aggregated in 1 m/s wind velocity classes (0-1 m/s, 1-2 m/s, etc.). In figure VII.3 these aggregated values (connected by lines to assist visibility) are plotted for day and night separately. It is clear that in many cases the 5-minute period values of $L_{95}$ are higher, in less cases lower than the long-term value. This means that if the immission limit is based on the measured long-term background sound level, then in a significant amount of time the actual background level will not be equal to the previously established long-term background level. In many instances the actual value of $L_{95}$ is higher than the long-term background level $L_{95,\text{lt}}$, which would allow for more wind turbine sound at that time.\(^1\)

VII.3.2 Ambient sound level controlled sound emission

An alternative to a wind velocity controlled emission level is to measure the ambient sound level itself and thus determine the actual limit value directly. If the limit is $L_{95}$, then the immission level must be $L_{\text{imm}} \leq L_{95}$. To achieve this the background ambient sound level can be determined by measurement (e.g. in 5-minute intervals) and compared to the immission level calculated from the actual turbine performance. If the immission level $L_{\text{imm}}$ would exactly equal the ambient background level $L_{95}$ without turbine sound, it would attain its maximum value $L_{\text{imm,max}} = L_{95}$. Then background sound level including turbine sound would be $L_{95+\text{wt}} = \log\text{sum}(L_{\text{imm,max}} + L_{95}) = L_{\text{imm,max}} + 3$ dB or $L_{\text{imm,max}} = L_{95+\text{wt}} - 3$ dB. If the calculated immission level exceeds the measured ambient level $L_{95+\text{wt}} - 3$, turbine sound apparently dominates the background level and the turbine should slow down.

\(^1\) perhaps for this reason the approach in the British ETSU-R-07 guideline [ETSU 1996] is to not use the long-term $L_{A90,\text{lt}}$, but an average of 10 minute $L_{A90,10\text{min}}$ values; this odd statistical construction can be viewed as an inefficient compromise that effectively allows excess of an appropriate limit in half of the time and a too severe limit in the other half.
This type of control can also be achieved in several steps. Again assuming 5-minute measurement periods, these are:

1. determine the actual sound power level $L_{W,5\text{min}}$ (integrated over 5 minutes) from turbine power production or speed.
2. determine $L_{\text{imm}}$ from the previously established relation $L_{\text{imm}}(L_W)$.
3. measure actual background level $L_{95+\text{wt},5\text{min}}$ at a location where the limit applies;
4. if $L_{\text{imm}} > L_{95+\text{wt},5\text{min}} - 3 \text{ dB}$, then $L_{W,5\text{min}} > L_{W\text{max}}$ and the control system must decrease sound power level for the next 5-minute period, if $L_{W,5\text{min}} < L_{W\text{max}}$ the reverse must happen (until maximum speed is attained).

Here it is assumed that the microphone is on a location where immission level must not exceed the ambient background level. If a measurement location is chosen further away from the turbine(s), the immission sound level will decrease with a factor $\Delta L_{\text{imm}}$ at constant $L_W$, whereas $L_{95}$ will not change (assuming that 5-minute ambient background sound does not depend on location). In this case a correction must be applied to the measured $L_{95+\text{wt}}$ ($L_{\text{imm, max}} = L_{95+\text{wt}} - 10\log(1+10^{-0.1\Delta L_{\text{imm}}})$) to determine what sound power level is acceptable. An advantage of a more distant measurement location is that it is less influenced by the turbine sound. A similar approach may be used if the limit is not $L_{95}$ itself, but $L_{95} + 5 \text{ dB}$. In that case, is it not possible to determine $L_{95}$ from measurements at a location where this limit applies, as the turbine sound is allowed to be twice as intense as background sound itself. In that case a measurement location may be chosen where, e.g., $\Delta L_{\text{imm}} = 5 \text{ dB}$.

An apparent drawback of this sound based control is that measured ambient sound may be contaminated by local sounds, that is: from a source close to the microphone, increasing only the local ambient sound level. Also, figure VII.3 suggests that there are significant variations in $L_{95,5\text{min}}$, which could imply large control imposed power excursions if these variations occur in short time.

The first drawback can be solved by using two or more microphones far enough apart not to be both influenced by a local source. The limit value is
then either $L_{95,5\text{min}}$ determined from all measured sound levels within the previous 5-minute period, or the lowest value of $L_{95,5\text{min}}$ from each microphone location. It must be borne in mind that the value of $L_{95,5\text{min}}$ is not sensitive to sounds of short duration. Sounds from birds or passing vehicles or airplanes do not influence a measured $L_{95,5\text{min}}$ significantly, except when they are present for most of the time within the 5 minute period.

With regard to the second point: large variations in either wind velocity or background sound level are rare, as is shown in figure VII.4 where the difference is plotted between consecutive 5-minute values of $L_9$ and average free 10-m wind velocity. The change in wind velocity averaged over consecutive periods of 5 is less than 0.5 m/s in 72% of the time, and less than 1.5 m/s in 99% of the time. The change in background sound level over consecutive periods of 5 minutes is less than 2.5 dB in 88% of the time and less than 3.5 dB in 94% of the time. So, if the adjustment of sound power level is in steps no larger than 3 dB, most changes can be dealt with in a single step. This also holds when a longer averaging period of 15 minutes is chosen: the change in background sound level over

![Figure VII.4: frequency distributions of changes per 5 and per 15 minutes of average wind velocity and background ambient sound level in classes of one unit (dB or m/s)](image-url)
consecutive periods of 15 minutes is less than 2.5 dB in 89% of the time and less than 3.5 dB in 96% of the time. The frequency of changes between 5-minute periods that are 10 minutes apart (that is: with two 5-minute periods in between) is very similar to the distributions in figure VII.4. This means that when there is a change of 3 dB for two consecutive periods, it is unlikely a similar change occurs within the next one or two periods.

VII.4 Reduction of fluctuations in sound level

The level variation due to blade swish increases when the atmosphere becomes more stable because the angle of attack on the blade changes. As a result the turbulent layer at the trailing edge of the blade becomes thicker and produces more sound. In a wind farm the increased level variations from two or more turbines may coincide to produce still higher fluctuations. The increase of blade swish, or rather: blade beating, may be lessened by adapting the blade pitch angle, the increase due to coincidence (also) by desynchronizing turbines.

VII.4.1 Pitch angle

When a blade rotates in a vertical plane the optimum blade pitch angle $\alpha$ is determined by the ratio of the wind velocity and the rotational speed of the blade. As the rotational speed is a function of radial distance (from the hub), blade pitch changes over the blade length and is lowest at the tip. As the wind velocity closer to the ground is usually lower, the wind velocity at the low tip (where the tip passes the tower) is lower than at the high tip. As a result the angle of attack changes within a rotation if blade pitch is kept constant. For a 100 m hub height and 70 m diameter turbine at 20 rpm this change (relative to hub height) is about $0.5^\circ$ at the lower tip in an unstable atmosphere, increasing to almost $2^\circ$ in a very stable atmosphere (see section V.1). Added to this is a further change (of the order of $2^\circ$) in the angle of attack in front of the tower due to the fact that the tower is an obstacle slowing down air passing the tower. At the high tip the change in angle of attack is $-0.3^\circ$ (unstable) to $-1.7^\circ$ (very stable).
The optimum angle of attack of the incoming air at every position of the rotating blade can be realized by adapting the blade pitch angle to the local wind velocity. Pitch must then increase for a blade going upward and decrease on the downward flight. Such a continuous change in blade pitch is common in helicopter technology. If the effect of stability on the wind profile would be compensated by pitch control, blade swish due to the presence of the tower would still be left. This residual blade swish can be eliminated by an extra decrease in blade pitch close to the tower. If the variations in angle of attack can be reduced to 1° or less, blade swish will cause variations less than 2 dB which are not perceived as fluctuating sound.

VII.4.2 Rotor tilt

If the rotor is tilted backwards, a blade element will move forward on the downward stroke and backward on the upward stroke, thus having a varying velocity component in the direction of the wind. As a result the angle of attack will change while the blade rotates because the flow angle will depend on blade position. If the tilt angle changes from zero to \( \theta \), the flow angle at the low tip increases from \( \varphi \) to \( \varphi' \) (see figure III.2). From geometrical considerations (see figure VII.5) of a blade segment tilted around a horizontal axis, it follows that \( C \cdot \sin \varphi + r \cdot \tan \theta = r \cdot \tan(\theta + \gamma) \), where \( \gamma = \arctan(C \sin \varphi / r) \). This leads to:

\[
sin \varphi' = S \cdot (\tan[\theta + \arctan(\sin \varphi / S)] - \tan \theta) \quad \text{(VII.1)}
\]

where \( S = r / C \) is the ratio of radius \( r \) and blade width (or chord length) \( C \) at radius \( r \). For small blade pitch angles and blade slenderness \( S \) between 10 and 40 the...
increase of blade pitch with tilt (from 0 to $\theta$) can be approximated with:

$$
\Delta \phi = \varphi' - \varphi = 1.1 \cdot \varphi \cdot \theta^2 \quad \text{(angles in radians)} 
$$

(VII.2a)

For values of $\varphi$, $S$ and $\theta$ in the range $\varphi \leq 10^\circ$, $30 \leq S \leq 50$ and $\theta \leq 20^\circ$, the standard deviation of the constant 1.1 is 0.01. With angles expressed in degrees, equation VII.2a reads:

$$
\Delta \phi = 33 \cdot 10^{-5} \varphi \cdot \theta^2 \quad \text{(angles in degrees)} 
$$

(VII.2b)

This means that for a tilt angle of $2^\circ$ and a $6^\circ$ blade pitch (tip rotational speed 70 m/s, induced wind velocity 10 m/s, angle of attack $2^\circ$), the change in angle of attack (relative to a vertical rotor with zero tilt) is negligible ($0.008^\circ$). Rotor tilt could now compensate a $1^\circ$ change in angle of attack at the low tip when the tilt angle is $22^\circ$. In this case the horizontal distance between the low tip and the turbine tower increases with approximately 15 m. This will in turn lead to a smaller change in angle of attack as at this distance the velocity deficit due to the presence of the tower is lower. For higher values of the blade pitch angle (ceteris paribus implying lower values of the angle of attack) increasing the tilt angle has a bigger effect. A substantial tilt however has major disadvantages as it decreases the rotor surface normal to the wind and induces a flow component parallel to the rotor surface which again changes the inflow angle. It therefore does not seem an efficient way to reduce the fluctuation level.

**VII.4.3 Desynchronization of turbines**

When the atmosphere becomes stable, large scale turbulence becomes weaker and wind velocity is more coherent over larger distances. The result is that different turbines in a wind farm are exposed to a wind with less variations, and near-synchronization of the turbines may lead to coincidence of blade beats from two or more turbines for an observer near the wind farm, and thus higher pulse levels (see section V.2.4). To desynchronize the turbines in this situation, the random variation induced by atmospheric turbulence (such as occurs in an unstable and neutral atmosphere) can be simulated by small and random fluctuations of the blade pitch angle or the electric load of each turbine separately.
In an unstable atmosphere turbulence strength peaks at a non-dimensional frequency $n = \frac{f z}{V} \approx 0.01$, where $V$ is the mean wind velocity and $z$ is height (this is according to custom in acoustics; in atmospheric physics traditionally $f$ is non-dimensional and $n$ physical frequency). At $z = 100$ m and $V = 10$ m/s this corresponds to a physical frequency $f = nV/z = 1$ mHz. At higher frequencies the turbulence spectral power density decreases with $f^{-5/3}$. When atmospheric instability decreases, the maximum shifts to a higher frequency and wind velocity fluctuations in the non-dimensional frequency range of 0.01 to 1 tend to vanish. So, to simulate atmospheric turbulence the blade pitch setting of each turbine (or the load imposed by the generator) must be fed independently with a signal corresponding to noise such as pink ($f^{-1}$) or brown ($f^{-2}$) noise, in the range of appr. 1 to 100 mHz. The (total) amplitude of this signal must be determined from local conditions, but is of the order of 1°.

**VII.5 Conclusion**

Wind turbine noise has shown to be a complex phenomenon. In the future quieter blades will be available, reducing sound emission by some 2 dB. The only presently available effective measures to decrease the sound impact of modern turbines are to create more distance or to slow down the rotor.

In existing turbines the sound immission level can be decreased by controlling the sound emission, which in turn is decreased by slowing down the rotor speed. When the limit is a single maximum sound immission level, this in fact dictates minimum distance for a given turbine and there is no further legal obligation to control.

In other cases the control strategy will depend on whether the legally enforced limit is a 10-m wind velocity or an ambient background sound level dependent limit. The 10-m wind velocity or the background sound level act as the control system input, blade pitch and/or load on the rotor is the controlled parameter. In both cases a suitable place must be chosen to measure the input parameter. For background sound level as input it is probably necessary to use two or more inputs to minimize the influence of local (near-microphone) sounds. It may however be the best strategy in
relatively quiet areas as it controls an important impact parameter: the level above background or intrusiveness of the wind turbine sound. 

Controlling sound emission requires a new strategy in wind turbine control: in the present situation there is usually more room for sound in daytime and in very windy nights, but less in quiet nights.

A clear characteristic of night time wind turbine noise is its beating character. Even if the sound emission level does not change, annoyance may decrease by eliminating the rhythm due to the blades passing the tower. Again, a lower rotational speed will help as this reduces the overall level including the pulse level. A better solution is to continuously change the blade pitch, adapting the angle of attack to local conditions in each rotation. This will also be an advantage from an energetic point of view as it optimizes lift at every rotor angle, and it will decrease the extra mechanical load on the blades accompanying the sound pulses.

When the impulsive character of the sound is heightened because of the interaction of several turbines in a wind farm, this may be eliminated by adding small random variations to the blade pitch, mimicking the random variations imposed by atmospheric turbulence in daytime when this effect does not occur.
Figure VIII.0: foam wind screens
VIII RUMBLING WIND: wind induced sound in a screened microphone

VIII.1 Overview of microphone noise research

It is commonly known that a wind screen over a microphone reduces ‘wind noise’ that apparently results from the air flow around the microphone. An explanation for this phenomenon has been addressed by several authors. According to a dimensional analysis by Strasberg [1988] the pressure within a spherical or cylindrical wind screen with diameter D in a flow with velocity V, depends on Strouhal number $\text{Sr} = \frac{fD}{V}$, Reynolds number $\text{Re} = \frac{DV}{\nu}$ and Mach number $\text{M} = \frac{V}{c}$ (where $\nu$ is the kinematic viscosity of air and $c$ the velocity of sound). Writing the rms pressure in a relatively narrow frequency band centered at frequency $f$ as $p_f$, and in dimensionless form by division with $\rho V^2$, Strasberg found: $\frac{p_f}{\rho V^2} = \text{function}(\text{Sr}, \text{Re}, \text{M})$.

Comparison with measured 1/3 octave band levels from four authors on 2.5 - 25 cm diameter wind screens, in air velocities ranging from 6 to 23 m/s yielded a definite expression for 1/3 octave frequency band:

$$20 \cdot \log_{10} \left( \frac{p_{1/3}}{\rho V^2} \right) = -23 \cdot \log_{10} \left( \frac{f_mD}{V} \right) - 81$$

(VIII.1)

where $f_m$ is the middle frequency of the 1/3 octave band. The data points agreed within appr. 3 dB with equation VIII.1 for $0.1 < \frac{fD}{V} < 5$, except for one of the fourteen data series where measured values diverged at $\frac{fD}{V} > 2$. Equation VIII.1 can also be written in acoustical terms by expressing the rms pressure as a sound pressure level relative to 20 $\mu$Pa:

$$L_{1/3} = 40 \cdot \log_{10} \left( \frac{V}{V_o} \right) - 23 \cdot \log_{10} \left( \frac{f_mD}{V} \right) + 15$$

(VIII.2)

Here $V_o$ is a reference velocity of 1 m/s and $\rho = 1.23$ kg/m$^3$ is used (air density at 1 bar and 10 °C). Equation VIII.2 is slightly different from the expression given by Strasberg because SI-units are used and terms in logarithms have been non-dimensionalized.

Morgan and Raspet pointed out that all measurements reported by Strasberg were made in low turbulence flows, such as wind tunnel flow [Morgan et al 1992]. Strasberg’s result thus referred to the wake created by a wind screen and excluded atmospheric turbulence (as Strasberg had
noted himself in his concluding remarks [Strasberg 1988]). Outdoors, however, the flow is turbulent, and induced pressure variations are expected to depend on meteorological parameters also. Morgan & Raspet applied Bernoulli’s principle by decomposing the wind velocity \( U \) in a constant time-averaged velocity \( V \) and a fluctuation velocity \( u \) with a time average \( u = 0 \), to obtain the rms pressure fluctuation \( p = \rho Vu \) [Morgan et al 1992] (in this chapter italics are used to denote the rms value \( x \) of a variable \( x : x = \sqrt{\langle x^2 \rangle} \)). This method can be compared to Strasberg’s model for a microphone in turbulent water flow [Strasberg 1979]. Measurements in wind velocities of 3 – 13 m/s at 30.5 m and 1.5 m height for different screen diameters (90 and 180 mm) and screen pore sizes (10, 20, 40 and 80 ppi) yielded:

\[
p = \alpha \rho (Vu)^k
\]  

(VIII.3)

with \( \alpha \) ranging from 0.16 to 0.26 and \( k \) from 1.0 to 1.3 [Morgan et al 1992]. For some measurements Morgan et al showed spectra over almost the same frequency range where equation VIII.1 is valid (0.1 < fD/V < 5). The spectra have a positive slope up to 3 Hz, possibly due to a non-linear instrumental frequency response. At higher values the slope is roughly comparable to what Strasberg found, but values of 20\log_{10}(p_{1/3}/\rho V^2) are generally 8 – 20 dB higher as predicted by equation VIII.1, implying that atmospheric turbulence dominated expected wake turbulence.

Zheng and Tan tried to solve this problem analytically [Zheng et al 2003]. Their analysis applies to low frequency variations, so the velocity variation \( u \) is uniform over the wind screen. Zheng & Tan state that this assumption seems to be valid for a low screen number \( D/\lambda, (< 0.3) \), the ratio between screen diameter and wavelength. Ignoring viscous effects (i.e. infinite Reynolds number), and calculating the pressure variation \( p(0) \) at the center of a spherical wind screen caused by pressure variations at the surface induced by a wind velocity \( U = V + u \), they found \( p(0) = -\frac{1}{2} \rho Vu \) or:

\[
p(0) = \frac{1}{2} \rho Vu
\]  

(VIII.4)

Comparison with equation VIII.3 shows that now \( \alpha = 0.5 \) and \( k = 1 \).
Finally, in this overview, Boersma [1997] found that sound spectra due to wind measured at 1.5 m above flat, open grassland were in good agreement with Strasberg’s results. However, Boersma used 95 percentile levels ($L_{95}$) which he estimated to be 6 to 13 dB lower than equivalent sound levels in the range considered ($30 < L_{95} < 70$ dB) [Boersma 1997], but he did not apply a level correction. So, in fact he found that his wind related spectra had slopes comparable to Strasberg’s, but with a $6 - 13$ dB higher value, not unlike the Morgan & Raspet spectra.

So, from literature we conclude that air turbulence creates pressure fluctuations especially at low frequencies, but the origin -wake or atmospheric turbulence- has not been definitely resolved.

In this chapter we will try to estimate the level of pressure variations due to atmospheric turbulence, \textit{i.e.} the 'sound' pressure level taken from a sound level meter caused by turbulence on the microphone wind screen. First we will describe the spectral distribution of atmospheric turbulence and the effect this turbulence has on a screened microphone. Then we will turn to measured spectra related to wind, obtained by the author as well as by others. Finally the results will be discussed.

\textbf{VII.2 Atmospheric turbulence}

A wind borne eddy that is large relative to the microphone wind screen (hence the change of wind velocity is nearly the same all over the wind screen) can be regarded as a change in magnitude and/or direction of the wind velocity [Zheng \textit{et al} 2003]. The change in the magnitude of the velocity causes a change in pressure; the change in direction is irrelevant for a spherical wind screen as nothing changes relative to the sphere. As we saw in the previous section, when the velocity $U$ is written as a constant (average) wind velocity $V$ and a fluctuating part $u$, and similarly $P = P_{\text{average}} + p$, the relation between the rms microphone pressure fluctuation $p$ and the rms wind velocity fluctuation $u$ is $p = \alpha p V u$. For inviscid flow $\alpha = 0.5$. For finite Reynolds numbers ($\text{Re}/10^4 \approx 0.5 - 15$ for wind screens of 4 - 20 cm and wind velocities of $2 - 12$ m/s), screening is better [Zheng \textit{et al} 2003], and $\alpha \leq 0.5$; Morgan & Raspet [1992] found $\alpha = 0.16 - 0.26$. The
pressure level due to atmospheric turbulence can be expressed as a sound pressure level $L_{at}$ (with reference pressure $p_{ref} = 20 \, \mu Pa$):

$$L_{at}(u) = 20 \cdot \log_{10}(\alpha p V u / p_{ref})$$  \hspace{1cm} (VIII.5)

which is frequency dependent because of $u$.

**VIII.2.1 Turbulence spectra**

Turbulent velocity fluctuations $v$ and $w$ also exist perpendicular to the average wind velocity, in the vertical ($w$) as well as horizontal ($v$) direction, and are of the same order of magnitude as in the longitudinal direction [Jensen et al. 1982]. Zheng & Tan [2003] showed that the effect of these fluctuations on the pressure at the microphone can be neglected in a first order approximation, as it scales with $v^2$ and $w^2$ and is therefore second order compared to the effect of the component $u$ in line with the average wind velocity $V$ that scales as $Vu$.

Atmospheric turbulence is treated in many papers and textbooks (such as [Jensen et al. 1982, Zhang et al. 2001]), also in reference to acoustics (see, e.g., [Wilson et al. 1994]). Here a short elucidation will be presented, leading to our topic of interest: turbulence spectra.

Atmospheric turbulence is created by friction and by thermal convection. Turbulence due to friction is a result of wind shear: at the surface the wind velocity is zero whereas at high altitudes the geostrophic wind is not influenced by the surface but a result of large scale pressure differences as well as Coriolis forces resulting from earth's rotation. In between, in the atmospheric boundary layer wind velocity increases with height $z$, equation III.2 is valid and for convenience repeated here:

$$V = (u^*/\kappa) \cdot [\ln(z/z_o) - \Psi]$$  \hspace{1cm} (VIII.6)

For $-1 < \zeta < 1$, $\Psi(\zeta)$ is of the same order of magnitude as the logarithmic term in equation VIII.6 ($2 < \ln(z/z_o) < 6$ for $1 < z < 5$ m, $1 < z_o < 10$ cm). Hence, at the same height and roughness length, $V$ may still change appreciably due to (in)stability.
The friction created by wind shear produces eddies over a range of frequencies and lengths, their size determined by \( z \) and \( V \). These eddies break up in ever smaller eddies and kinetic turbulent energy is cascaded to smaller sizes at higher frequencies, until the eddies reach the Kolmogorov size \( \eta_s \) (\( \approx 1 \) mm) and dissipate into heat by viscous friction. It has been shown by Kolmogorov that for this energy cascade, in the so-called inertial subrange of the turbulent spectrum, the frequency dependency follows the well known 'law of 5/3': the spectrum falls with \( f^{-5/3} \).

It is customary in atmospheric physics to express turbulence frequency in dimensionless form \( n \), with \( n = f z / V \) (in fact \( n \) and \( f \) are usually interchanged, but we will use \( f \) for dimensional frequency, as is usual in acoustics). The seminal Kansas measurements showed that the squared longitudinal velocity fluctuation \( u_f^2 \) per unit frequency in a neutral atmosphere depends on frequency as [Kaimal et al 1972]:

\[
f u_f^2 / u_*^2 = 105 n (1 + 33 n)^{-5/3}
\]  

(VIII.7)

The experimentally determined constants in this equation, the non-dimensional turbulent energy spectrum, are not exact, but are close to values determined by others [Garrat 1992, Zhang et al 2001]. For \( n << 1 \), the right-hand side approximates 105\( n \), which, with \( n = f z / V \) and equation VIII.6, leads to \( u_f^2 = 105 u_*^2 z / V = 105 \kappa^2 z V \cdot [\ln(z/z_o) - \Psi]^{-2} \). Applying this to VIII.5, the induced pressure level per unit of frequency appears to be independent of frequency, but increases with wind velocity (\( \sim 30 \cdot \log V \)). For \( n >> 1 \) the right-hand side of equation VIII.7 reduces to \( 3.2 \cdot (33 n)^{-2/3} \), leading to \( u_f^2 = 0.3 \cdot u_*^2 \cdot (V / z)^{2/3} \cdot f^{-5/3} \), which describes the inertial subrange. The frequency where the wind velocity spectrum VIII.7 has a maximum is \( n_{max} = 0.05 \) or \( f_{max} = 0.05 V / z \). As sound measurement are usually at heights \( 1 < z < 5 \) m, \( f_{max} \) is less than 1 Hz for wind velocities \( V < 20 \) m/s.

When insolation increases the surface temperature, the atmosphere changes from neutral to unstable and eddies are created by thermal differences with sizes up to the boundary layer height with an order of magnitude of 1 km. Turbulent kinetic energy production then shifts to lower frequencies. In contrast in a stable atmosphere, where surface temperature decreases because of surface cooling, eddy production at low frequencies
(corresponding to large eddy diameters) is damped and the spectral maximum shifts to a higher frequency up to appr. \( n = 0.5 \) for a very stable atmosphere. As low-altitude wind velocities \( (z < 5 \text{ m}) \) in a stable atmosphere are restricted to relatively low values (for higher wind velocities, stability is disrupted and the atmosphere becomes neutral), the spectral maximum may shift up to \( 0.5V/z \approx 3 \text{ Hz} \). The inertial subrange thus expands or shrinks at its lower boundary, but its frequency dependency follows the ‘law of 5/3’.

**VIII.2.2 Effect on microphone in wind screen**

The spectrum of longitudinal atmospheric turbulence in the inertial subrange was described in the previous section with the (squared) rms value of velocity variation per unit frequency \( u_f^2 = 0.3 \cdot u^* \cdot (V/z)^{2/3} \cdot f^{-5/3} \). It is convenient to integrate this over a frequency range \( f_1 - f_2 \) to obtain a 1/3-octave band level \( (f_m = 2^{1/6} f_2 = 2^{1/6} \cdot f_1) \) with centre frequency \( f_m: u_{1/3}^2 = 0.046 \cdot u^* \cdot (f_m \cdot z/V)^{2/3} = [0.215 \cdot u^* \cdot (f_m \cdot z/V)^{1/3}]^2 \). Substituting \( u^* \) from equation VIII.6 and applying the result to equation VIII.5 for 1/3 octave band levels \( L_{at,1/3}(f_m) = 20 \cdot \log(\alpha \rho V u_{1/3}/p_{ref}), \) yields:

\[
L_{at,1/3}(f) = 40 \cdot \log(V/V_o) - 6.67 \cdot \log(zf/V) - 20 \cdot \log[\ln(z/z_o)/\Psi] + C \quad (VIII.8)
\]

Here the frequency index \( m \) as well as the logarithm index 10 have been dropped, as will be done in the rest of the text. In equation VIII.8 \( C = 20 \cdot \log(0.215 \cdot u^* \cdot V_o^2/p_{ref}) = 62.4 \text{ dB} \) for \( \kappa = 0.4, \alpha = 0.25, \rho = 1.23 \text{ kg/m}^3 \) and pressure level is taken re \( p_{ref} = 20 \text{ µPa} \). For octave band levels \( L_{at,1/1}(f) \) the constant \( C \) in the right hand side of VIII.8 is 67.2 dB.

Equation VIII.7 does not apply to frequencies where eddies are smaller than the wind screen. The contribution of small eddies will decrease proportional to the ratio of eddy size \( (\ell^2, \text{ where } \ell \text{ is the eddy length scale and } f = V/\ell) \) and wind screen surface \( \pi D^2 \). When this ratio decreases more eddies will simultaneously be present at the screen surface and resulting pressure fluctuations at the surface will more effectively cancel one another in the interior of the wind screen. The pressure variation in the wind screen centre resulting from one eddy is proportional to the size of
the eddy relative to the screen surface, \( \frac{\ell^2}{D^2} \), but also the screen centre pressure resulting from the random contributions of all \( N \) eddies on the screen surface is proportional to \( \sqrt{N} \), where \( N \approx \frac{D^2}{\ell^2} \). The resulting screen centre pressure is thus proportional to individual eddy pressure \( p_f \) and \( \left( \frac{\ell^2}{D^2} \right) \sqrt{\left( \frac{D^2}{\ell^2} \right)} = \ell/D = V/\ell \). Consequently a factor \(-20\cdot\log(fD/V)\) must be added to the resulting rms pressure level.

In wind noise reduction measured by Morgan there is a change in frequency dependency at screen number \( D/\ell \approx 1/3 \) ([Morgan 1993], see also [Zheng et al 2003]). We therefore expect at sufficiently high frequencies the pressure level at the microphone to decrease proportional to \( 20\cdot\log(D/\ell) \), relative to the level in equation (VIII.8), and this decrease must vanish when \( D/\ell = Df/V < 1/3 \), i.e. below the cut-off frequency \( f_c = V/(3D) \). As the change will be gradual, a smooth transition can be added to equation VIII.8:

\[
L_{at,1/3}(f) = 40\cdot\log(V/V_0) - 6.67\cdot\log(zf/V) - 20\cdot\log[\ln(z/z_0)\cdot\Psi] + 
-10\cdot\log(1+(f/f_c)^2) + C 
\] (VIII.9a)

With usual screen diameters 5 – 25 cm and wind velocities 1 - 20 m/s, the cut-off frequency is in the range of 1 to 100 Hz. With the common 10 cm diameter wind screen \( f_c \) will usually be in the infrasound region. Equation VIII.9a can be rewritten with Strouhal number \( Sr = fD/V \) as independent variable of a ‘meteorologically reduced’ 1/3 octave band level \( L_{red} \):

\[
L_{red,1/3} = L_{at,1/3} - 40\cdot\log(V/V_0) + 20\cdot\log[(z/D)^{1/3}\cdot(\ln(z/z_0)\cdot\Psi)] = 
-6.67\cdot\log(Sr) - 10\cdot\log[1+(3Sr)^2] + C 
\] (VIII.9b)

The levels according to equation VIII.9 have been plotted in figure VIII.1 for different wind velocities and with \( z = 20\cdot D = 40\cdot z_0 = 2 \) m, \( \Psi = 0 \). For \( f < 0.5\cdot f_c \) the term before \( C \) is less than 1 dB and equation VIII.9a reduces to equation VIII.8. For frequencies \( f >> f_c \) the term before \( C \) in equation VIII.9b reduces to \(-20\cdot\log(3Sr)\) and equation VIII.9b can be written as:

\[
L_{red,1/3} = -26.67\cdot\log(Sr) + C - 9.5 \] (VIII.10a)

This can be rewritten in a aerodynamic terms as:

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\[ L_{p,1/3} = 20 \cdot \log(p/\rho V^2) - 26.67 \cdot \log(Sr) + F(z) + C_p \]  \hspace{1cm} (VIII.10b)

where and \( F(z) = -20 \cdot \log\left(\frac{z}{D}\right)^{1/3} \cdot (\ln\left(\frac{z}{z_o}\right) - \Psi) \) and \( C_p = 20 \cdot \log(0.215 \kappa \alpha) - 9.5 = -43 \text{ dB} \). For \( F(z) = -20 \text{ dB} \) (e.g. a 10 cm diameter wind screen at a \( z = 2 \text{ m}, z_o = 5 \text{ cm} \) and \( \Psi = 0 \)) the right hand side of equation VIII.10b is \(-26.67 \cdot \log(Sr) - 63\). Comparing this with Strasberg’s result (equation VIII.1 and gray lines in figure VIII.1) we see that the frequency dependency is slightly different, and levels are \(13 - 19 \text{ dB} \) higher \((0.5 < Sr < 20)\), which is of the order of what we found in the measurements by Boersma and Raspet et al (see section VIII.1). The change in slope, visible at Strouhal number\( Df_c/V = 1/3 \) in figure VIII.1, is a feature not explained by the earlier authors.

### VIII.2.3 Frequency regions

From the theory above it can now be concluded that the wind induced pressure level on a (screened) microphone stretches over four successive frequency regions:

1. at very low frequencies (less than a few Hz) the turbulence spectrum is in the energy-producing subrange; 1/3 octave band pressure level \( L_{at,1/3} \) is independent of frequency (white noise), but increases with wind velocity;

2. at frequencies up to \( f_c = 0.3V/D \), which is usually in the infrasound region, the turbulence spectrum is in the inertial subrange, \( L_{at,1/3} \sim 46.7 \cdot \log V \) and \( -6.7 \cdot \log f_c \);

3. at higher frequencies, but still in the inertial subrange, eddies average out over the wind screen more effectively at increasing frequency.

Figure VIII.1: black lines: calculated 1/3 octave band levels \( L_{at,1/3} \) due to atmospheric turbulence at wind velocities of (bottom to top) 2, 4, 6, 8 and 10 m/s; \( F(z) = -18 \text{ dB} \); gray lines: levels at same wind velocities according to Strasberg
(L_{at,1/3} \sim -26.7\cdot \log f), but pressure level increases faster with wind velocity (L_{at,1/3} \sim 66.7\cdot \log V);

iv. at frequencies beyond 0.1V/\eta_\infty (see [Plate 2000, p. 585]) atmospheric turbulence enters the dissipation range and turbulence vanishes. This is in the range \( Sr = fD/V > 0.1D/\eta_\infty \approx 100[D/m] = D/cm. \)

The inertial subrange (ii and iii) is of most interest here, as it is within the commonly used range of acoustic frequency and level.

### VIII.2.4 Wind induced broad band A-weighted pressure level

In figure VIII.2 1/3-octave band levels according to equation VIII.9 are plotted for different wind velocities for \( z = 50\cdot z_o = 20\cdot D = 2 \text{ m} \) (or \( F(z) = -20.5 \text{ dB} \) with \( \Psi = 0 \)). Also levels are plotted after A-weighting to show the relevance to most acoustic measurements, where wind induced noise may be a disturbance added to an A-weighted sound level. At the frequency where turbulent eddies enter the dissipation subrange (\( f \approx 0.1V/\eta_\infty \)), no data are plotted as the turbulent velocity spectrum falls very steeply and induced pressure levels are considered negligible. A-weighted pressure levels \( L_{at,A} \) can be calculated by summing over all 1/3-octave bands. The wind velocity dependency can then be determined from the best fit of \( L_{at,A} \) vs. \( V \):

\[
L_{at,A} = 69.4\cdot \log (V/V_o) - 26.7\cdot \log (D/\ell_o) + F(z) + C - 74.8 \quad \text{(VIII.11a)}
\]

[Figure VIII.2: calculated linear (dashed) and A-weighted (solid lines) 1/3-octave pressure levels due to atmospheric turbulence on a screened microphone with \( F(z) + C = 42 \text{ dB}, D = 0.1 \text{ m} \) and wind speeds 2, 4, 6, 8, 10 m/s (black, bottom to top); bold grey lines: 1/3 octave band levels according to Strasberg for 10 m/s]
where $\ell_o = 1 \text{ m}$ is a reference length. Equation VIII.11a has the same structure as VIII.10a, but a rather higher slope with $\log V$ because higher frequencies (with lower A-weighting) are progressively important, and a much smaller constant term as a result of A-weighting. The slope decreases with wind screen diameter and is 65.5 dB when $D = 1.25 \text{ cm}$ (unscreened $\frac{1}{2}''$ microphone), but is constant within 1 dB for $5 < D/\text{cm} < 50$. Equation VIII.11a is not very sensitive for the cut-off at $f = 0.1V/\nu$: if spectral levels are integrated over all frequencies, total level does not increase significantly at high wind velocities, and with less than 3 dB at low wind velocities. It will be noted that the slope with wind velocity is slightly higher than for individual spectral levels for $f > f_C$ (66.7 dB, see equation VIII.10a, due to lower A-weighting at the increasingly higher frequencies.

If we put $G(z) = F(z) - 6.7 \cdot \log(D/\ell_o) + 14 = -20 \cdot \log[0.2(z/\ell_o)^{1/3} \cdot (\ln(z/\ell_o - \Psi)]$, and use 10D for convenience, equation VIII.11a becomes:

$$L_{at, A} = 69.4 \cdot \log(V/V_o) - 20 \cdot \log(10D/\ell_o) + G(z) + C - 68.8 \quad (\text{VIII.11b})$$

Now for $z_o = 2.5 - 6 \text{ cm}$ and $\Psi = 0$, $G(2 \text{ m}) = 0 \pm 1 \text{ dB}$. This means that for a 10 cm wind screen and measurement over a flat area with a low vegetation cover in neutral conditions $L_{at, A} \approx 69.4 \cdot \log(V/V_o) - 6.4 \text{ dB(A)}$.

Figure VIII.3 is a plot of equation VIII.11 with $G(z) = 0$, $C = 62 \text{ dB}$. Also plotted in figure VIII.3 is the relation according to Strasberg, obtained by A-weighting and integrating equation VIII.2 over $f$. 

**Figure VIII.3:** calculated A-weighted broad band pressure level caused by atmospheric or wake turbulence with $G(z) + C = 62.4 \text{ dB}$ and $D = 0.1 \text{ m}$. 

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VIII.3  Comparison with experimental results

VIII.3.1 Measured spectral pressure levels

Several authors have performed measurements to determine spectral levels due to wind, including wind induced sound pressure fluctuations. We will use data from Larsson and Israelsson [1982], Jakobsen and Andersen [1983] and Boersma [1997] from screened as well as unscreened microphones. Table VIII.1 gives an overview of measurement parameters. None of the authors give the degree of stability, but in Jakobsen’s data $\Psi \leq 0$ (night), in Boersma’s $\Psi \geq 0$ (summer’s day). Jakobsen mentions roughness height of the location (a golf course), Boersma grass height ($\approx 10$ cm), Larsson only mentions measurement height over grass at either 1.25 or 4 m, without specifying which height applies to a measurement result. To prevent using spectra at large values of $|\Psi|$ no data at low wind velocities ($< 2$ m/s at microphone) are used. This is also recommendable as at low wind velocity sound not related to wind is more likely to dominate. We preferably use $L_{eq}$ data. However, these are not available from Boersma. Boersma used 95 percentile levels ($L_{95}$), but we have $L_{50}$ values from the original data. Though Boersma quotes $L_{Aeq} \approx L_{A50}$, we will use $L_{Aeq} \approx L_{A50} + 3$, in agreement with long term data on wind noise [Van den Berg 2004b] and assume this to be valid for every frequency band. If measurements yielded octave band levels, 4.8 dB was subtracted to obtain the 1/3 octave band level at the same frequency.

Also $L_{eq}$ values are presented from measurements made by the author at several locations; at one location (Zernike) for the purpose of wind noise measurements, and otherwise (Horsterwold, Kwelder) selected for having little other noise. Here also the degree of atmospheric stability is unknown, as at the time of measurement it was not known to be a relevant factor. The ‘Zernike’ measurements were done at the university grounds (latitude $53^\circ 14'43''$, longitude $6^\circ 31'48''$) with both the microphone (in a spherical foam screen of 2.5, 3.8 or 9.5 cm diameter) and the wind meter at 1.2 or 2.5 m over grass at least several hundred meters from trees, and an estimated roughness height of 5 cm. They were performed in daytime in December 2003 and August 2004 with a fair wind under heavy clouding.
The ‘Kwelder’ measurements were made in daytime or evening in July and August of 1996 at an open area at the Dutch coast (latitude 53°25'46", longitude 6°32'40"), consisting of level land overgrown with grass and low weeds and close to tidal water. Sound measurements were taken at a height of 1.5 m at times when no sound could be heard but wind-related sound and distant birds. The microphone was fitted with a spherical 9.5 cm diameter foam wind screen. Wind velocity at microphone height at 1.5 m was estimated from measured wind velocity at 5 m height with equation VIII.6, $z_0$ estimated as 2 cm. Finally the ‘Horsterwold’ measurements were made in December 2001 in an open space with grass and reeds (latitude 52°18'3", longitude 5°29'38") between 5 to 10 m high trees at a distance of approximately 30 m but further in the windward direction, in a mostly clouded night. Wind velocity and sound were measured at 2 m height, the wind screen was a 9 cm diameter foam cylinder. Due to the differences in vegetation, roughness length here was difficult to estimate, and was determined by fitting measurement results to the expected level (resulting in 60 cm and a more limited range of values of $\Psi$ to fit).

At very low frequencies in our Zernike measurements the 1/3 octave band levels were corrected for non-linear response. The frequency response of the B&K ½" microphone type 4189 is specified by Brüel & Kjaer [B&K 1995] and is effectively a high pass filter with a corner frequency of 2.6 Hz. The response of the Larson Davis type 2800 frequency analyser is flat (±1 dB) for all frequencies.

To plot spectra we calculate the reduced pressure level $L_{\text{red,1/3}}$, leaving only the screen diameter based Strouhal number $Sr = fD/V$ as the independent variable. Octave band pressure levels $L_{\text{red,1/1}}$ are substituted by $L_{\text{red,1/3}} + 4.8$. As atmospheric stability is as yet unknown, the stability function is set to zero. If wind velocity was not measured at microphone height, the logarithmic wind profile (equation (VIII.6 with $\Psi = 0$, or III.3) is used to determine $V_{\text{mic}}$ from the wind velocity at height $h$.

Linear spectra of 1/3-octave levels are plotted in the left part of figure VIII.4 for the unscreened microphones. Also plotted is the spectrum according to Larsson et al [1982], valid for the inertial subrange. Due to
the small size of the unscreened microphone (1.25 cm) part of the spectrum lies in the dissipation range at frequencies \( f > 0.1V/\eta \approx 100V/m \), corresponding to \( Sr > 100D/m = 1.25 \).

In figure VIII.4B spectra are plotted from screened microphones, from the data from Larsson, Jakobsen and Boersma. As these spectra were determined with a range of screen diameters, the change from the inertial to the dissipation subrange extends over a range of non-dimensional frequencies (Strouhal numbers). Finally figure VIII.4C shows spectra from the Horsterwold, Zernike and Kwelder measurements. In all figures spectra deviate from the predicted spectrum at high Strouhal numbers because either the lower measurement range of the sound level meter is reached or

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<th>Table VIII.1: wind induced noise measurement characteristics</th>
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notes: 1: # of measurements  2: estimated; 3: fitted; 4: no = unscreened; 5: observations of unknown length; 6: 1/1 or 1/3 octave band
Figure VIII.4:
reduced 1/3 octave band pressure levels at different wind velocities (in legend: $V$ in m/s), bold line is predicted spectrum;

A: unscreened microphone, from Larsson (open symbols) and Boersma (black symbols);

B: screened microphone, from Larsson (open symbols), Jakobsen (grey) and Boersma (black symbols);

C: screened microphone, measurements in Horsterwold (open symbols), Kweider (grey) and Zernike (black symbols).
Figure VIII.5: same as figure VIII.4, but after fitting with stability function

Figure VIII.6: values of the stability function $\Psi$ found by fitting reduced spectral levels $L_{red}$ with theoretical spectrum, for measurements in day or night time, and for unscreened microphones in daytime
ambient sound dominates the wind-induced pressure level. Also, at these high Strouhal numbers most values are in the dissipation range where the present model is not valid.

In figure VIII.4 atmospheric stability has not been taken into account yet (in fact $\Psi = 0$ was used), due to lack of data to determine $\Psi$. In stable conditions ($\Psi < 0$) $L_{\text{red}}$ will be higher, in unstable conditions ($\Psi > 0$) lower, causing the plotted spectra to shift vertically if the proper value $\Psi \neq 0$ is applied.

If wind velocity at microphone height is deduced from wind velocity at another height, the shift is more complex, as stability then also affects the term $40 \cdot \log(V/V_0)$ as well as the ordinate value $S_r = fD/V$. The approach taken here is to vary $\Psi$ to obtain a best fit to the theoretical value of the $L_{\text{red}}$ levels at non-dimensional frequencies in the inertial subrange. The fitted spectra are plotted in figure VIII.5. The values of $\Psi$ that gave the best fits are plotted in figure VIII.6, categorized in daytime and night time measurements (where one would expect $\Psi \geq 0$ and $\Psi \leq 0$, respectively). Measurements with unscreened microphones are indicated separately, and are in daytime for Boersma’s measurements and probably also for Larsson’s, so one would expect $\Psi \geq 0$.

### VIII.3.2 Measured broad band pressure levels

Several authors give a relation between broad band A-weighted sound pressure level $L_A$ and wind velocity [Boersma 1997, Larsson et al 1982, Jakobsen et al 1983]. According to Boersma $L_A \sim 22.6 \cdot \log(v)$ (with $v$ measured at 2 m height, $L_A$ at 1.5 m), to Larsson $L_A = 4.4 \cdot v + 27.5$ ($v$ and $L_A$ measured at the same height), to Jakobsen $L_A = 6.8 \cdot v - 2.6$ ($v$ measured at 10 m, $L_A$ at 1.5 m). However, as Boersma clearly shows, most of the A-weighted sound is due to ambient wind induced sound, especially at low wind velocities. So we cannot use these relations for just sound induced by wind on the microphone.

A practical situation where the influence of wind on the microphone + wind screen could be investigated directly offered itself when on May 28, 2000 a storm occurred during our 'Wieringerwaard' measurements. The
microphone, in a 9 cm foam cylinder, and a wind meter were both placed at a height of 4.6 m, 2 m apart, in front of a big farmer's shed 5 m to the west of the microphone (latitude $52^\circ 48' 41''$, longitude $4^\circ 52' 23'''$). A second, 'free wind' windmeter at 10 m height was placed further away to measure undisturbed wind. Around the measurement location were fields with potato plants of 20 - 30 cm height. As it was May, an unstable atmosphere is expected in daytime, leaning to neutral when the wind velocity increases.

Some measurement results are given in figure VIII.7 (all values are 10 minute averages of samples measured at a rate of 1 s$^{-1}$). In the left part of the figure the 'free' wind velocity $v_{10}$ is seen to increase to 20 m/s (72 km/h) in the course of the day after a relatively quiet night. The wind velocity $v_{mic}$ near the microphone increased at practically the same rate between 6 and 12 o'clock, but then abruptly falls from 13 m/s to 2 m/s and thereafter remains at a low value even while the 'free' wind velocity is still increasing. Up to 12 o'clock the sound level (equivalent A-weighted level per 10 minutes) increases in proportion to the wind velocity reaching a maximum of 84 dB(A), but then falls abruptly to 50 dB(A) at the same time the local wind velocity collapses. In this morning the unobstructed wind began in the east and gradually turns south. When at 12 o’clock the wind passes behind the shed, the microphone is suddenly taken out of the wind. There is no reason that the sound reaching the microphone changes significantly during this change, but due to the sudden wind velocity reduction the measure sound pressure level drops to 50 dB(A). After that the sound pressure level increases again as long as the storm is gaining strength. The measured pressure level above 60 dB(A) is pure wind-induced ‘pseudo’ sound, that is: sound resulting from moving air, not from airborne sound.

In the right part of figure VIII.7 the A-weighted equivalent (pseudo-) sound pressure level per 10 minutes over the same period as in the left part of figure 7, is plotted as a function of wind velocity at the microphone. There is an obvious direct correlation between pressure level and wind velocity at higher wind velocities ($V \geq 6$ m/s) in contrast to the levels at
lower wind velocities. Again, the stability factor $\Psi$ is not known, but in daytime and in strong winds it must be small and positive. The lines in figure VIII.7 show the calculated pressure levels for plausible values $0 < \Psi < 1$ (with $z_0 = 20$ cm), encompassing the measured values.

**VIII.3.3 Screen reduction**

For two of our Zernike summer measurements (see table VIII.1) with place and atmospheric conditions unchanged within the measurement period, the difference between $1/3$ octave band pressure levels measured with an approximately spherical 2.4 cm wind screen and a spherical 9.5 cm wind screen are plotted in figure VIII.8. Also...
plotted is the calculated screening effect based on equation VIII.9a, with only both term before C differing between both measurements. It appears that the measured screening effect is on average approximately 1 dB higher than the calculated level. It is not clear why the difference in screening is negative at frequencies below 2 Hz. For a somewhat smaller wind screen (18 mm < D < 24 mm) the average screening effect would agree better with the calculated effect.

VIII.4 Discussion

The model developed in this paper starts with the assumption that wind induced ‘sound’ pressure levels on a microphone are caused by atmospheric turbulence. Then, at low non-dimensional frequencies (Sr << 0.3) spectral levels are determined entirely by atmospheric turbulence. In this frequency range a wind screen has no effect. At higher frequencies, where pressure fluctuations tend to cancel one another more effectively as their scale decreases relative to the wind screen diameter, a wind screen acts as a first order low pass filter for turbulent fluctuations. In this frequency range (Sr > 0.3) a wind screen diminishes the effect of turbulence, and better so if it is bigger.

Wind induced pressure levels are determined not just by wind velocity and screen diameter, but also by two factors that are relevant for the production of turbulence: atmospheric instability and surface roughness. The stability factor Ψ and roughness height z₀ are determinants for thermal and frictional turbulence, respectively. These determinants are usually not taken into account with respect to wind induced noise and are consequently not reported. Atmospheric stability therefore had to be estimated by varying the value of Ψ until a best fit was obtained of measured spectra to the calculated spectrum. Roughness length, when unknown, was assumed to be comparable to vegetation height.

The values of Ψ that resulted in the best fits are shown in figure VIII.6. They can also be compared to values obtained from long term measurements at the Cabauw measurement site of the Royal Netherlands Meteorological Institute (KNMI). The Cabauw site is in open, flat land west of the central part of the Netherlands (see Chapter VI) and may be considered representative for locations in comparable terrain in the north.
and central parts of the Netherlands (Boersma’s and our measurements), Denmark (Jakobsen et al) and the Swedish Uppsala plain (Larsson et al). The KNMI provided us with a data file containing 30 minute averages of the Monin-Obukhov length $L$ over one year (1987). From this the dimensionless height $\zeta = z/L$ can be calculated and then the stability factor $\Psi$ (see text below equation VIII.6). In figure VIII.9 the frequency distribution is shown of all 17520 ($= 2 \cdot 24 \cdot 365$) values of $\Psi$, for two altitudes: 2 m and 5 m. Also the frequency distribution is shown of the 42 values of $\Psi$ resulting from our fitting procedure. The distribution of our fitted values resemble the distribution of actually occurring values (in 1987) and thus seems plausible.  

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Figure VIII.9: frequency distributions of stability factor $\Psi$ at 2 m and 5 m height, based on 1/2 hour observations over 1987 and resulting from fitted spectra
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Two constants are not known accurately: $\alpha$, assumed to have a value 0.25, and the ratio of screen diameter and eddy size at the corner frequency, where 3 was used. If the $Sr$-related slopes are as in equation VIII.9b, the best fit of all data points in figure VIII.5 at $Sr < 2.5$ is a line $L_{red,1/3} = -6.7 \cdot \log(Sr) - 10 \cdot \log[1 + (3.8 \cdot Sr)^2] + 62.0$. This fit is within 2.2 dB of the calculated value (equation VIII.9b). It follows that the ratio $\ell/D$ (3.8) where screen averaging over eddies sets in may be greater than assumed (viz. 3), and the constant term may be somewhat smaller, which could be a result of a lower value of $\alpha$ than assumed (0.24 instead of 0.25). For $2.5 < Sr < 16$ the best fit is on average 2.1 dB above the calculated value. The standard deviation of the measured 1/3 Strouhal octave band levels is less than 3.5 dB at $Sr < 2.5$ and up to 7 dB at $2.5 < Sr < 16$.  

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VIII.5 Applications

As microphone wind noise appears to be closely correlated to atmospheric turbulence, acoustic measurements can alternatively be used to measure turbulence spectra or turbulence strength, especially in the inertial subrange. This provides a new way to determine (e.g.) friction velocity or atmospheric stability. As the measured signal decreases above the corner frequency \( f_c = V/(3D) \) this frequency is best chosen high, which can be achieved with a small, bare microphone.

The present model can be used to distinguish wind induced noise from other wind related sound. An application is the measurement of wind turbine sound or (without an operating wind turbine) ambient background sound in relatively strong winds. If the measurement is on a wind exposed site it is probable that at high wind velocities wind induced noise influences or even dominates either wind turbine sound or proper ambient sound. A measured level can now be corrected for wind induced sound with a calculated wind noise level. In less exposed sites it is usually not clear in what degree the measured levels are influenced by wind induced noise. To calculate wind induced noise levels additional measurements are necessary to determine roughness height and atmospheric stability. Stability can be estimated from wind velocity measurements on two heights, using equation VIII.6. Roughness height can be estimated from tabulated values or from wind velocity measurement at two heights in a neutral atmosphere, at times when the logarithmic wind profile is valid (equation VIII.6 with \( \Psi = 0 \)). In neutral and stable conditions wind induced noise levels are not very sensitive to errors in roughness height: with an error of a factor of 2 in \( z_o = 10 \) cm, the level changes less than 2 dB if microphone height is 3 m or more.

VIII.6 Conclusion

Measured spectra, reduced with a term for wind velocity and turbulence strength, coincide well with calculated values for unscreened as well as screened microphones in the range where the theoretical model (equation VIII.9) is valid. To test the model more thoroughly, measurements should
include a determination of roughness length and atmospheric stability, in addition to the usual measurement of wind velocity and measurement height.

The model shows that to avoid high wind induced pressure levels, measurements are best performed at low wind velocity and with a large diameter wind screen, which is common knowledge in acoustics. The overall reduction $\Delta L_A$ from a bigger wind screen relative to a smaller one is determined by the ratio of the screen diameters $D_1$ and $D_2$: $\Delta L_A = 20 \cdot \log(D_2/D_1)$ (from equation VIII.11b, $D > 5$ cm). A wind screen does not reduce noise from atmospheric turbulence at frequencies $f < V/(3D)$.

The model also shows that, to reduce wind induced sound, it helps to measure over a low roughness surface and at night (stable atmosphere), as both factors help to reduce turbulence, even if the (average) wind velocity on the microphone does not change. With reduced turbulence, wind induced pressure levels will finally reach the level given by Strasberg (equation VIII.1 or VIII.2), where turbulence is the result of the wake caused by the wind screen.

One might be tempted to think that a higher measurement altitude would also help to reduce wind noise (as this would make $G(z)$ in equation VIII.11b more negative, thus reducing $L_{at,A}$). However, in practice increasing altitude will lead to higher wind velocities, especially so in a stable atmosphere, and the first term in equation VIII.11b would more then compensate the decrease in $G(z)$. It is therefore preferable to measure at low altitude if less wind noise is desired.
IX GENERAL CONCLUSIONS

The research aims formulated in the introductory chapter (section 1.6) have been addressed separately in the previous chapters. In this chapter we present an overview of all results. The results are presented in a logical order, which is not entirely in the sequence of the previous chapters.

IX.1 Effect of atmospheric stability on wind turbine sound

It is customary in wind turbine noise assessment to calculate the sound level on neighbouring premises by assuming hub height wind velocities predicted using a logarithmic wind profile. This wind profile depends only on surface roughness and is valid in a neutral atmosphere. However, it is not a predictor for wind profiles in either an unstable or stable atmosphere. Especially in a stable atmosphere a wind profile can be very different from the logarithmic, neutral profile and the hub height wind velocity is higher than predicted by the neutral profile. As more wind at hub height makes a variable speed wind turbine rotate at a higher speed, the sound power level may be significantly higher in a stable atmosphere at the same wind 10-m velocity $V_{10}$ (which usually occurs when the sun is down and no strong near-ground wind is present) than in an unstable atmosphere (usually when the sun is up). This is especially relevant for modern, that is: tall and variable speed, wind turbines.

A stability dependent wind profile predicts the wind velocity at hub height more accurately. When a correct wind profile is used, calculated immission sound levels agree with measured night-time sound immission levels.

Sound immission measurements have been made at distances up to 2 km from the Rhede wind farm containing seventeen 98 m hub height, variable speed wind turbines, and at 280 m from a single 45 m hub height, two speed wind turbine at Boazum. Measured immission sound levels at 400 m west of the Rhede wind farm almost perfectly match (average difference: 0.1 dB) sound levels calculated from measured emission levels near the turbines. At distances up to 2 km the calculated level may underestimate
the measured level, but the discrepancy is small: 1.5 dB or less.\textsuperscript{1} Thus, from the measurements both the emission and immission sound levels could be determined accurately. As both levels can be related through a propagation model, it may not be necessary to measure both: immission measurements can be used to assess immission as well as emission sound levels of an entire wind farm.

The level of aerodynamic wind turbine noise depends on the angle of attack: the angle between the blade and the incoming air flow. Increasing atmospheric stability also creates greater changes in the angle of attack over each rotation, resulting in stronger turbine sound fluctuations. It can be shown theoretically for a modern turbine rotating at high speed that, when the atmosphere becomes very stable, the fluctuation in turbine sound level increases to approximately 5 dB. This value is confirmed by measurements at a single wind turbine where the maximum sound level periodically rises 4 to 6 dB above the minimum sound level within short periods of time. At some distance from a wind farm the fluctuations from two or more turbines may arrive simultaneously for a period of time and increase the fluctuation level further at the observer's position up to approximately 9 dB. This effect develops in a stable atmosphere because the spatial coherence in wind velocity over distances at the size of an entire wind farm increases. As a result turbines in the farm are exposed to a more constant wind and rotate almost synchronously. Because of this near-synchronicity, the fluctuations in sound level will for some time coincide at some locations, causing an amplification of the fluctuation. The place where such an amplification occurs will sweep over the area with a velocity determined by the difference in rotational frequency. The magnitude of this effect thus depends on stability, but also on the number of wind turbines and their distances to the observer.

Blade passing frequency is the parameter determining the modulation frequency of wind turbine sound. Human perception is most sensitive to

\textsuperscript{1} In one night the sound level at over 2 km from the wind farm was much higher than calculated, probably because of an inversion layer adding more downward refracted sound. This apparently rare occurrence at the Rhede wind farm could be more significant where high inversion layers occur more often.
modulation frequencies close to 4 Hz and the modulated sound has a frequency of approximately 1000 Hz. The hypothesis that fluctuations are important is supported by descriptions given by naïve listeners as well as residents: turbines sound like ‘lapping’, ‘swishing’, ‘clapping’, ‘beating’ or ‘like the surf’. It is probable that this fluctuating character is responsible the relatively high annoyance caused by wind turbine sound and a deterioration of sleep quality.

Atmospheric stability also affects the energy yield of wind turbines: relative to the ‘standard’ (neutral) atmosphere, a stable atmosphere increases the yield, especially for modern tall turbines. The reverse is true for an unstable atmosphere, though to a lesser degree. Perhaps atmospheric stability was not recognized as an important determinant for wind power as the underestimated night time yield is compensated partly by the overestimated daytime yield. The annual effect will depend on the average magnitude as well as the prevalence of atmospheric stability.

**IX.2 Effect of atmospheric stability on ambient background sound**

The change in wind profile at night also results in lower ambient background levels than expected: at night the wind velocity near the ground may be lower than expected from logarithmic extrapolation of the wind velocity at 10 m, resulting in lower levels of wind induced sound from low vegetation. The contrast between wind turbine and ambient sound levels is therefore at night more pronounced.

**IX.3 Wind noise on a microphone**

To avoid high wind induced pressure levels in windy conditions, outdoor measurements are best performed with a large diameter wind screen. The overall reduction from a bigger wind screen relative to a smaller one is
determined by the ratio of the screen diameters. A wind screen does not reduce noise from atmospheric turbulence at very low frequencies.¹

In a stable atmosphere the low near-ground wind velocity creates less wind noise on the microphone. As a result, sound measurements during a stable night are much less influenced by wind induced microphone noise (and other sounds as well, since nights are usually more quiet) than in a neutral or unstable atmosphere. The results in this book shows that wind turbine sound can be measured accurately at great distances (up to 2 km) if the atmosphere is stable.

The model developed in this thesis shows that, in order to reduce wind induced sound, it helps to measure over a low roughness surface and in a stable atmosphere, as both factors help to reduce turbulence, even if the average wind velocity on the microphone does not change. But in a stable atmosphere near-ground wind velocities will usually be low, decreasing wind induced noise further. With increasing stability, wind induced pressure levels will drop and finally reach a low level determined by turbulence in the wake of the wind screen.

**IX.4 Degree of atmospheric stability**

Stability is a property of the atmosphere, in principle occurring all over the earth. It depends on surface properties and weather conditions which determine the magnitude and evolution over time of the heat balance in the atmospheric boundary layer. Most important are differences in heat transfer at the surface (water, soil) and in the atmosphere (atmospheric humidity and clouds, wind mixing). With current knowledge, the effects of stability on the wind profile over flat ground can be modelled satisfactorily. In mountaineous areas terrain induced changes on the wind profile influence the stability related changes and the outcome is less easily predicted: these changes can weaken as well as amplify the effect of atmospheric stability.

¹ frequencies below $V/(3D)$, where $V$ is the wind speed at the microphone and $D$ the wind screen diameter
Results from various onshore, relatively flat areas show that in daytime the ratio of the wind velocity at 80 m (hub height) and the wind velocity at reference height of 10 m is 1.25 to 1.5. This ratio is in agreement with the usual logarithmic wind profile for low roughness lengths (low vegetation). At night the situation is quite different and the ratio has a much wider range with values from 1.7 to 4.3. At night high altitude wind velocities thus can be (much) higher than expected from logarithmic extrapolation of 10-m wind velocities.

IX.5 Measures to mitigate stability related effects

Presently available measures to decrease the immission sound level from modern turbines are to create more distance to a receiver or to slow down the rotor, preferably by an optimized control mechanism. Quieter blades as such will always be advantageous, but expected changes are modest and will not eliminate the beating or thumping character due to atmospheric stability.

Controlling the stability related sound emission requires a new strategy in wind turbine control and wind farm design. In the present situation there is usually more latitude for sound (and energy) production in daytime, but less during quiet nights. A strategy for onshore wind farms might be to use more of the potential in daytime, less at night.

A control strategy may depend on whether the legally enforced limit is a 10-m wind velocity or an ambient background sound level dependent limit. The 10-m wind velocity or the background sound level can act as the control system input, with blade pitch the controlled variable. In both cases a suitable place must be chosen to measure the input parameter. For background sound level as input it is probably necessary to use two or more inputs to minimize the influence of local (near-microphone) sounds.

An ambient background controlled emission level may be the best strategy in relatively quiet areas as it controls an important impact parameter: the level above background or intrusiveness of the wind turbine sound.

Even if the sound emission level does not change, annoyance may be diminished by eliminating the rhythm due to the beating character of the sound. A solution is to continuously change the blade pitch, adapting the
angle of attack to local conditions during rotation. This will probably also be an advantage from an energetic point of view as it optimizes lift at every rotor angle, and it will decrease the mechanical load ‘pulses’ on the blades accompanying the sound pulses.

Increased fluctuation due to the interaction of sound from different turbines can be eliminated by adding small random variations to the blade pitch or rotor load, mimicking the random variations imposed by atmospheric turbulence in daytime when this effect does not occur.

**IX.6 Recommendations**

When night time is the critical noise period, wind turbine sound levels should be assessed taking into account stable atmospheric conditions. When the impulsive character of the sound is to be assessed, this should be carried out in times of a stable atmosphere, as that is the relevant condition for impulsiveness.

When ambient sound is considered as a sound masking wind turbine sound, neither sound should be related to wind velocity at 10 meter reference height via a (possibly implicit) neutral or ‘standard’ wind profile. A correct, stability dependent wind profile should be used. In flat and certainly in mountainous terrain one should determine directly the relationship between hub height wind velocity on the one hand and ambient background sound at an immission location on the other hand, in order to eliminate any badly correlated, intermediate wind velocity.

Also, in the assessment of wind turbine electrical power production the sole use of a neutral wind profile (a ‘standard atmosphere’) should be abandoned as it yields data that are not consistent with reality.

When comparing stable and unstable atmospheric conditions, the difference in sound power as well as in sound limits can lead to new control strategies and onshore wind farm concepts. Presently only distance is a factor used to minimize noise impact. A wind farm can be optimized with a strategy that maximizes power output while keeping sound power within limits. When daytime immission levels do comply with the noise
limits, but nighttime immission levels do not, a control system can be implemented to reduce the turbine speed when necessary.
In new turbine designs continuous blade pitch control could be applied to increase energy yield and reduce annoyance at the same time by eliminating the thumping character of the emitted sound.
This is the end of my tour of discovery, of over two years of reading about and trying to understand atmospheric physics and wind turbines, of measurements and theory, of applying knowledge and expertise in physics and acoustics to a new topic. Of course there is much more to discover: indeed, it looks like wind turbines have become more fascinating now their sound has proved to be more complex than a simple constant noise from the sky, driven only by wind with a constant profile. This may motivate researchers and consultants to put more effort in better predictions of wind turbine noise, and considering again noise exposure to local residents.

This period began with publishing the results of the measurements at the Rhede wind farm and it ended, seemingly symbolically, with the first International Conference on Wind Turbine Noise in Berlin in October 2005. At that conference there was a general acknowledgment that wind turbine sound is not the simple issue we once thought it was. At the conference many delegates agreed that, looking back, the internationally used ‘standard wind profile’ might have been misleading people by suggesting it was, everywhere and always, the best wind profile. Although the widely used IEC-61400 standard certainly does not state that, a less careful reader might think it did, finding no alternative profile in the standard. Thus, it becomes a question of careful communication and taking into account that acoustic consultants do (did?) not have the knowledge to apply the standard in ‘non standard’ conditions. Paul Botha [2005] proposed to do away with 10-m wind velocities entirely and relate background sound directly to hub height wind velocity. This is a sensible idea as it relates the two factors that are most relevant, wind turbine sound and ambient sound, without an intermediate variable (10-m wind velocity). It will lead to better insight in the masking capability of background sound: the ability to mask (= make inaudible) unwanted sound is not only dependent on wind velocity, but also on atmospheric stability and wind direction.
The Berlin conference helped me solve a riddle. Malcolm Hayes had written me before that according to his observations blade swish is caused by the blade that is going down, not by the blade being in the downward position (passing the mast). This seems contradictory to my conclusion that blade beating is due to blades passing the mast. Oerlemans [2005] showed that close to the tower Malcolm was right, but this could not explain blade swish far away from a turbine. So what we heard depended on the distance to the turbine, which is also true for other sound phenomena: further away from the turbine the sound has a lower pitch, the pulses can be amplified by synchronicity of turbines and it can be louder under an inversion layer. This point again illustrates that one must be careful when generalizing observations.

I don’t expect the problem of the distinct, beating character of wind turbine sound to be solved easily. Though I am convinced the sound character is a major factor in wind turbine noise annoyance, a 5 dB penalty for an impulsive character of the sound may indeed impede wind farm projects as a wind farm will need more ‘empty space’. Also, the sound is not as impulsive as gun shots or hammering are, giving way to a discussion on whether it is ‘really’ impulsive (5 dB penalty) or not (no penalty). Is it possible to have a truly independent opinion in a legally created dichotomy with such significant consequences? Several technical possibilities to minimize the noise have been outlined in this book, but we need not just depend on technical solutions. A change in public relations can also make a difference: proponents must accept that wind turbine noise is not (always) ‘benign’, that the noise may affect people, and that people who are complaining are not always just a nuisance. And no, we still do not understand wind turbine noise immission entirely, so proponents should watch their WARYDU attitude.
“..... about 80 per cent of the population supports wind power in the surveys investigated in this paper. On the local level the support of wind power in areas with operating wind power plants is equally high. (....) This, however, does not mean that protests will not appear. It takes only one devoted opponent to start for instance a legal procedure against a planning permit. This is one of the reasons why public conflicts over wind power plants have become the rule rather than the exception. Lack of communication between the people who shall live with the turbines, and the developers, the local bureaucracy, and the politicians seems to be the perfect catalyst for converting local scepticism, and negative attitudes into actual actions against specific projects. Conversely, information and dialogue is the road to acceptance.”

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I want to express my sincere gratitude to a number of people. Foremost is Diek (prof dr ir H Duifhuis) who for 20 years has been a true, though most of the time distant colleague, and who immediately responded positive to my request to be my promotor. The same, but for a shorter period, holds for my second promotor Ton (prof dr A J M Schoot Uiterkamp). Both were very confident in my capacities and my work and they also allowed me to change the subject without much ado.¹ “The point is, Frits”, they lectured me, “that you demonstrate your scientific capability, the subject as such is not that important”. And it is not by obligation that I also mention my wife Luci, who for two years gracefully took most of my share in cooking and household chores.

There are some others I would like to commemorate. My colleague Aart van der Pol who was always interested in new work and new ideas and sharp in his comments. Terry Mazilsky, one of the beleaguered residents, who read the first and the final chapters for language and clarity. My friend Dorothé Faber with whom I discussed my work from a very different, non-technical perspective. And the students and secretaries who had to abide with too little attention devoted to them and me again and again forgetting things. I hope they felt I did try to support them even though I was busy or being elsewhere in my thoughts. Much the same goes for my daughters Inge and Maya, who will meet me at the end of this period as my beloved paranymps.

Two organizations have supported my research. The province of Groningen has subsidized the measurement project at the Rhede wind farm to support the residents and thus helped to produce a report and after that a scientific publication on the effect of atmospheric stability. The British Renewable Energy Foundation gave a grant to elaborate my thoughts on

¹ Originally my thesis was to be about Sound monitoring in quiet areas, and Dick, Ton and I discussed the nature of sound and quietness and our response to it; this changed when I became ever more involved in wind turbines.
the beating character of wind turbine sound, which led to my second publication on wind turbine noise. The KNMI gave support by providing data that helped me prepare the presentation about the wind statistics.

My Faculty of Mathematics and Natural Sciences allowed me to spend half of my working time for two years on this promotion. Unfortunately there was no money for a substitute, so I had to fit my usual work in the other half of the time. It has been busy sometimes, but I am lucky to have work that I like (well, most of it) and is worthwhile, so I never count the hours from 9 to 5. It is, I think, significant that I have a position in a University based Science Shop, because that position enables me to spend time on projects for the benefit of citizen groups with no net financial return. This book shows that citizens may need the help of science to support their claims and improve their situation, which is not available elsewhere except perhaps at high costs. On behalf of the people I have helped I am grateful our University and Faculty are still firm supporters of the Science Shop idea.¹

¹ Don’t count your chickens before they are hatched! At the time this book was finalized, 21 years after the Science Shop for Physics started, it was announced that because of large financial deficits the Faculty executive board proposed to close down all four faculty science shops.
This study was started after complaints of residents that the sound of a wind farm was louder and more annoying than predicted, especially when there was little wind in the evening or at night. The explanation appeared to be the occurrence of another wind profile than that used to predict the noise impact (the wind profile describes how the wind velocity increases with height). There are probably several reasons why this was not found earlier: 1) because wind turbines become taller, there is a growing discrepancy between prediction and practice; 2) measurements are usually done in daytime when the wind profile resembles more closely the commonly used standard profile; 3) based on the sound that occurs in daytime, it is hard to imagine the sound can be so different at night; 4) “there are always people complaining”, so complaints are not always a reason for a thorough investigation; 5) at least some wind energy proponents prefer to downplay the disadvantages rather than solve them.

According to Dutch legislation and international guidelines the sound production of a wind farm can only be checked by measurements when the wind farm operator cooperates. The consequence is an implicit partiality in favor of the operator detrimental to independent verification. Because of the level of detail of instructions measurements and assessments are hampered and there is no margin for the very expertise of an investigator. For a lay person understanding the jargon was already utterly impossible and he cannot but hire an expensive expert to argue his case.

From this study one can conclude that through the use of a restricted model of reality, viz. a forever neutral atmosphere, experts have lost sight (temporarily) of the true reality in which a neutral atmosphere is not very prevalent. It is precisely the occurrence of complaints that may indicate such errors.

The sound of modern wind turbines is generated mainly by the flow of the wind along the blades. In this process a turbulent boundary layer develops at the rear side of the blade where trailing edge sound of relatively high
frequencies originates and which is radiated into the environment. This turbulent boundary layer becomes thicker and produces more sound when the wind flows in at a greater angle.

The inflowing wind is turbulent itself. The blade cuts through these turbulent movements and as a result again sound is generated: in-flow turbulence sound. Here lower frequencies dominate. Finally a blade also radiates sound when the forces on the blade change because of a local variation in wind velocity. This happens every time the blade passes the tower because there the wind is slowed down by the tower. On the one hand this causes more trailing edge sound due to the change in inflow angle, on the other hand more infrasound is generated because of the sudden sideways movement at the rate of the blade passing frequency.

For all these sounds loudness increases when the speed increases. Because the tip has the highest speed the sound of a wind turbine mainly comes from the blade tips. Moreover, for human hearing the trailing edge sound is most important because it is in an area of frequencies that we can hear well.

It is often assumed that there is a fixed relation between the wind velocity at hub height and at a reference height of 10 meter. This is the relation valid in a neutral or ‘standard’ atmosphere. No other relations are given in legislation or international guidelines for wind turbine sound that are valid in other conditions of the atmosphere, viz. the stable and unstable conditions.

The atmosphere is unstable when in daytime the air near the ground is relatively warm from contact with the surface heated by solar insolation. In that case vertical air movements originate and the wind profile is not equal to the profile in a neutral atmosphere, though it does not differ strongly. A stable atmosphere however has a markedly different wind profile. The atmosphere is stable when the air close to the ground is relatively cold due to contact with the ground surface when this cools down at night by radiating heat. A stable atmosphere occurs especially in nights with a partial or no cloud cover and the wind is not too strong (close to the ground). In a stable atmosphere the turbulence has decreased substantially
and as a result layers of air are less strongly coupled. The lower layer of air is thus less taken along with the wind that at higher altitudes keeps on blowing, giving rise to greater differences between wind velocities at different heights.

The present study was performed mainly near the Rhede wind farm close to the Dutch – German border. The farm consists of 17 1.8 MW turbines of 98 m hub height and three 35 m blades. The level of the incoming sound has been measured at a number of locations. The sound could be measured up to a distance of 2 km. It proved that, contrary to predictions, already at a weak wind (at 10 m height) the turbines could rotate at almost top speed and consequentially produce much sound.

It appeared that a wind profile proper to stable conditions could explain the measured sound levels excellently. At the same wind velocity at a reference height of 10 meter, wind turbines in a stable atmosphere generate more sound than in a neutral atmosphere, while at the same time the wind velocity near the ground is so low that the natural ambient sound due to rustling vegetation is weaker. As a result the contrast between wind turbine sound and natural ambient sound is more pronounced in stable conditions than it is in neutral conditions.

When the wind profile after sunset changes while the atmosphere becomes more stable, the difference in wind velocity over the rotor increases. This causes a change in the level of the trailing edge sound. At the low tip this is reinforced because the inflow angle already was less favourable due to the wind being slowed down by the presence of the mast. The differences in wind speed lead to variations in the sound radiated by the blade tips that reach their highest values when a tip passes the mast. For a modern, tall wind turbine the calculated variation is approximately 5 dB at night, whereas it is approximately 2 dB in daytime. This is perceived as a more pronounced fluctuation of the sound.

A more stable atmospheric boundary layer moreover implies that there is less atmospheric turbulence, so wind turbines in a farm will experience a more equal and constant wind. As a result, in a stable atmosphere wind turbines can, more than in daytime, run almost at the same speed and then
diverge again. With several turbines the fluctuations in sound can reinforce one another when they reach the ear of an observer simultaneously. With two turbines (at the same distance) this leads to an increase in level of 3 dB, with three turbines to an increase of 5 dB.

In measurements this reasoned upon effect indeed occurred. With a single 45 m high wind turbine at a distance of 280 m at night variations of 6 dB were found. Near the wind farm the variations were usually 5 dB, but they could rise to approximately 9 dB, as expected when the fluctuations of several turbines coincide.

From other research and from descriptions of residents one can establish that the sound of a wind turbine or wind farm becomes more annoying because of 'swishing', 'sloshing', 'clapping', 'beating' or 'thumping'. All descriptions mention a periodic variation on top of a constant noisy sound. This corresponds to the calculated and measured modulation of trailing edge sound. From psycho-acoustic research it has been shown earlier that human sensitivity to sound fluctuations is high at frequencies that occur in the night time sound of modern wind turbines. If this fluctuating sound is sufficiently loud in a bedroom it can cause sleep disturbance.

In the temperate climate zone a stable atmosphere is to be expected between sunset and sunrise over land if there is a -partly- clear sky (because clouds hinder the radiation of heat) and the wind is not too strong (because a strong wind promotes vertical heat exchange). From an analysis of measurements of the KNMI at Cabauw, in the central part of the Netherlands, up to an altitude of 200 m, it appears that there is a diurnal and seasonal pattern in the wind profile that correlates with the diurnal and seasonal variation in the heat exchange between the earth’s surface and the atmosphere. The fact that at sunset the wind often lies down is a consequence of the increasing atmospheric stability, and this decrease in wind velocity close to the ground is accompanied by an increase at higher altitudes. This has significant consequences for the energy production of wind turbines, where the rotor height plays an important part. If one starts from the measured wind velocities at Cabauw at 10 m height and a forever neutral atmosphere, the annually averaged electrical power generated by a 80 m high, 2 MW (reference) wind turbine would amount to almost
500 kW. However, based on the real, measured wind speed at 80 m height the annual power in reality amounts to 600 kW. So, because of atmospheric stability there is, relative to a neutral atmosphere, a significantly higher yield at night time hours, that even amply compensates for the lower yield in daytime hours.

The higher wind velocity at night on the rotor also causes a higher level of generated sound. If again one starts from the measured wind velocities at Cabauw at 10 m height and an atmosphere assumed to be neutral, the average sound power level generated by the reference wind turbine is 102 dB(A). In reality, however, it is 2 dB higher. This is also an average over an entire year; in separate nights the difference can be substantially higher, e.g. when a turbine rotates at (almost) top speed at a time it was expected to not produce at all because of the low 10 m wind velocity.

The degree of atmospheric stability at Cabauw is hardly different from what was observed at the Rhede wind farm. At other locations in countries in the temperate zone stability occurs to a similar extent. The consequences of atmospheric stability as described here, will thus occur at many wind farms that exist or are to be built in the temperate zone. However, above large bodies of water stability is rather a seasonal than a diurnal phenomenon, en in mountainous terrain the consequences of stability on the wind profile can be strengthened as well as weakened due to changes induced by height variations in the area.

The sound of a wind turbine or wind farm can thus become more annoying after sunset for two reasons: it becomes louder and the sound exhibits stronger fluctuations. At a given rotor diameter a blade can only be made less noisy with a different design or by slowing down the speed. A decrease in speed however reduces the generated electrical power and must therefore be applied only when necessary. To achieve this a control can be applied that lowers the speed when a noise limit is exceeded, increasing the speed again when the limit allows. This control could work on the generator and/or the pitch angle of the blades.
By changing the pitch angle while the blades rotate, the wind can flow in at an optimal angle at any position on the rotor, by which the energetic efficiency will increase on the one hand and the fluctuation strength of the sound will decrease on the other hand, even rendering the fluctuations inaudible. The total sound power will then decrease even relative to a neutral atmosphere, because the in-flow turbulence sound level will be lower due to the relative absence of atmospheric turbulence. Tilting the rotor to change the pitch angle during rotation does not appear to be a fruitful strategy: the tilt must be so great that the disadvantages will dominate.

The fluctuations near a wind farm can be stronger due to interference from the fluctuations of several turbines. This can be prevented by desynchronizing the turbines, as it happens in daytime by large scale atmospheric turbulence, by adding small and uncorrelated variations in the load of the rotors or the pitch angle of the blades of the individual turbines.

Controlling the sound production thus requires a new strategy for managing wind turbines: in daytime there is often more margin available for sound production than at night and this margin can be used in daytime in exchange for more restrictions at night.

Finally another, very different problem was addressed: the influence of wind on a microphone in or without a wind screen. When there is sufficient wind the microphone signal contains a low frequency, rumbling sound disturbing the measurement of ambient sound. This rumble is not sound from the environment, but is generated by pressure fluctuations caused by turbulent wind velocity variations. With a pressure sensitive microphone these pressure variations are not distinguishable from acoustical pressure variations. It appears that a wind screen is effective only by damping contributions of small turbulent eddies. A wind screen has no effect when eddies are bigger than the wind screen.

The strength of atmospheric turbulence does not only depend on the (average) wind velocity, but also on the local roughness of the earth surface and the stability of the atmosphere. These last two factors cause friction and thermal turbulence, respectively. The turbulence strength is
well known for an unobstructed wind flow over flat land. Turbulence is weaker in a stable and stronger in an unstable atmosphere. The ‘sound’ pressure level based on atmospheric turbulence appears to agree well with measured and published levels of wind induced pressure levels. Thus the influence of wind on a sound measurement in wind can be calculated. In reverse this calculation model yields a new method to measure the strength of atmospheric turbulence.

To conclude, it can be stated that with respect to wind turbine sound an important phenomenon has been overlooked: the change in wind after sunset. This phenomenon will be more important for modern, tall wind turbines and in view of the many wind farms that are planned. If this problem is not recognized and solved it will hamper the expansion of wind energy.
Bobby vraagt: 'Hoort u de windmolens wel eens?'

'Wat voor geluid maken ze?'

'Net als op elkaar slaand metaal, maar als er een echt harde wind staat worden de wieken vager en begint de lucht te schreeuwen van pijn.' Hij siddert.

'Waar zijn de windmolens voor?'

'Ze zorgen dat alles ’t doet. Als je je oor tegen de grond houdt kun je ze horen.'

'Wat bedoel je met alles?'

'De lichten, de fabrieken, de spoorwegen. Zonder de windmolens staat alles stil.'

Dit onderzoek is tot stand gekomen na klachten van bewoners dat het geluid van een windpark luider en hinderlijker was dan voorspeld, vooral als er ’s avonds of ’s nachts weinig wind was. De verklaring hiervoor bleek het optreden van een ander windprofiel dan werd gehanteerd bij de voorspelling van de geluidsbelasting (het windprofiel beschrijft hoe de windsnelheid toeneemt met de hoogte). Dat dit niet eerder is gevonden heeft waarschijnlijk meerdere redenen: 1) doordat windturbines hoger en groter worden is er een groeiende kloof tussen voorspelling en praktijk; 2) er wordt normaliter overdag gemeten wanneer het windprofiel meer lijkt op het gewoonlijk gebruikte standaardprofiel; 3) men kan zich, op grond van het overdag optredende geluid, moeilijk voorstellen dat het ’s nachts zo anders kan zijn; 4) “er zijn altijd wel mensen die klagen”, dus klachten zijn niet altijd een reden tot grondig onderzoek; 5) tenminste een aantal voorstanders van windenergie bagatelliseert liever de nadelen dan ze op te lossen.

Volgens de Nederlandse wetgeving en internationale richtlijnen kan de geluidsproductie van een windpark alleen door metingen gecontroleerd worden als de exploitant meewerkt. Het gevolg is een impliciete partijdigheid ten gunste van de exploitant en ten nadele van onafhankelijke

1 'The suspect', door Michael Robotham, Time Warner Paperbacks, 2003 (p. 151), vertaling G.P. van den Berg
controle. Ook door de gedetailleerdheid van voorschriften worden metingen en beoordelingen bemoeilijkt en is er geen ruimte meer voor de eigen deskundigheid van een onderzoeker. De burger kan het jargon al helemaal niet meer volgen en moet een dure deskundige inhuren om zijn zaak te beargumenteren.

Bij dit onderzoek kan men constateren dat deskundigen door het gebruik van een beperkt model van de werkelijkheid, namelijk een eeuwig neutrale atmosfeer, (tijdelijk) het zicht hebben verloren op de echte werkelijkheid waarin die neutrale atmosfeer niet zo vaak voorkomt. Juist klachten kunnen helpen om dergelijke dwalingen aan te wijzen.

Het geluid van moderne windturbines wordt vooral opgewekt door de stroming van de wind langs de wieken. Daarbij ontwikkelt zich een turbulente grenslaag aan de achterkant van de wiek waarin relatief hoogfrequent achterrandgeluid (‘trailing edge sound’) ontstaat dat wordt uitgestraald naar de omgeving. Deze turbulente grenslaag wordt dikker en produceert meer geluid als de wind onder een grotere hoek instroomt.

De instromende wind is zelf ook turbulent. De wiek snijdt door deze turbulente bewegingen heen waarbij weer geluid ontstaat: instromingsturbulentiegeluid (‘in-flow turbulent sound’). Hierin domineren lagere frequenties. Tenslotte straalt een wiek ook geluid af als de krachten op de wiek veranderen doordat de windsnelheid lokaal varieert. Dit gebeurt telkens als de wiek de mast passeert omdat daar de wind is afgeremd door de mast. Enerzijds ontstaat daarbij meer achterrandgeluid omdat de instromingshoek verandert, anderzijds ontstaat er ook infrageluid door de plotselinge zijwaartse beweging in het tempo van de wiekpasseerfrequentie.

Bij al deze geluiden neemt de sterkte ervan toe naarmate de snelheid groter is. Omdat de tip de hoogste snelheid heeft is het geluid van een windturbine vooral van de wiektips afkomstig. Voor het menselijk gehoor is bovendien het achterrandgeluid het belangrijkst omdat dat in een frequentiegebied ligt dat wij goed kunnen waarnemen.
Vaak wordt aangenomen dat er een vaste relatie is tussen de wind op as hoogte en op een referentiehoogte van 10 meter. Dit is de relatie die geldig is in een neutrale of ‘standaard’ atmosfeer. Er worden geen andere relaties gegeven in de wetgeving en in internationale richtlijnen die gelden bij andere toestanden van de atmosfeer, namelijk de stabiele en instabiele toestand.

De atmosfeer wordt instabiel als overdag de lucht nabij de grond relatief warm is door contact met het door zoninstraling verwarmde aardoppervlak. Er ontstaan dan verticale luchtbewegingen en het windprofiel is niet meer gelijk aan dat in een neutrale atmosfeer, maar wikt daar niet sterk vanaf. Een stabiele atmosfeer kent echter een duidelijk afwijkend windprofiel. De atmosfeer is stabiel als de lucht nabij de grond relatief koud is door contact met het door warmte-uitstraling afkoelende aardoppervlak ’s nachts. Een stabiele atmosfeer treedt vooral op tijdens niet gedeeltelijk of geheel onbewolkte nachten met niet teveel wind (aan de grond). In een stabiele atmosfeer is de turbulentie sterk verminderd met als gevolg dat luchtlagen minder sterk gekoppeld zijn. De onderste luchtlag wordt daardoor minder meegenomen door de wind die op grotere hoogte gewoon blijft doorwaaien, waardoor er grotere verschillen zijn tussen windsnelheden op verschillende hoogten.

Het hier beschreven onderzoek is grotendeels uitgevoerd bij windpark Rhede vlakbij de Duits-Nederlandse grens. Het park telt 17 1,8 MW turbines met een as hoogte van 98 m en drie wieken van 35 m lengte. Op een aantal punten is het niveau van het invallende geluid langdurig gemeten. Het geluid kon tot op 2 km afstand worden gemeten. Bij een zwakke wind (op 10 m hoogte) bleken de turbines, anders dan voorspeld, al op vrijwel topsnelheid te kunnen draaien en dientengevolge veel geluid te produceren. Een windprofiel dat bij stabiele omstandigheden past bleek de gemeten geluidsniveaus uistekend te kunnen verklaren. Bij een gelijke windsnelheid op een referentiehoogte van 10 meter, produceren windturbines in een stabiele atmosfeer meer geluid dan in een neutrale atmosfeer, terwijl dan tegelijkertijd de windsnelheid nabij de grond zo laag is dat het natuurlijke omgevingsgeluid van ruisende vegetatie zwakker is.
Het contrast tussen windturbinegeluid en natuurlijk omgevingsgeluid is daardoor bij stabiele omstandigheden groter dan bij instabiele.

Als het windprofiel na zonsondergang verandert door een stabielere wordende atmosfeer, wordt het verschil in windsnelheid over de rotor groter. Dit veroorzaakt een verandering in de sterkte van het achterrandgeluid. Bij de lage tip wordt dit nog versterkt doordat de instromingshoek al ongunstiger was vanwege de door de mast verlaagde windsnelheid. De verschillen in windsnelheid leiden tot variaties in het door de tips afgestraalde geluid die het grootst zijn als een tip de mast passeert. Voor een moderne, hoge windturbine bedraagt de berekende variatie ongeveer 5 dB ’s nachts, terwijl dit overdag ca. 2 dB is. Dit wordt ervaren als een duidelijker fluctuatie van het geluid.

Een stabielere atmosferische grenslaag betekent bovendien dat er minder atmosferische turbulentie is waardoor windturbines in een park een meer gelijke en meer constante wind ervaren. In een stabiele atmosfeer kunnen windturbines daardoor, méér dan overdag, een tijd nagenoeg gelijk lopen en weer langzaam uiteenlopen. Bij meerdere turbines kunnen de fluctuaties in het geluid elkaar versterken als ze het gehoor van een waarnemer gelijktijdig bereiken. Bij twee turbines (op gelijke afstand) leidt dit tot een 3 dB hoger niveau van de fluctuaties, bij drie turbines tot een 5 dB hoger niveau.

Bij metingen bleek dit beredeneerde effect daadwerkelijk voor te komen. Bij een enkele windturbine van 45 m as hoogte werden op een afstand van 280 m ’s nachts variaties gevonden van 6 dB. Bij het windpark bedroegen de variaties meestal 5 dB, maar ze konden oplopen tot ongeveer 9 dB, zoals verwacht wordt bij het samenvallen van de fluctuaties van meerdere turbines.

Uit onderzoek elders en uit beschrijvingen van omwonenden kan men constateren dat het geluid van een windturbine of windpark vooral na zonsondergang hinderlijker wordt door het ‘zoeven’ of ‘klotsen’, ‘klappen’, ‘slaan’ of ‘bonken’. De omschrijvingen vermelden steeds een periodieke variatie bovenop een constant ruisachtig geluid. Dit correspondeert met de berekende en gemeten modulatie van het achterrandgeluid. Uit psycho-akoestisch onderzoek is veel eerder al
gebleken dat de menselijke gevoeligheid voor geluidsfluctuaties hoog is bij frequenties die juist voorkomen in het nachtelijke geluid van moderne turbines. Als dit fluctuerende geluid voldoende luid doordringt in een slaapkamer kan het tot slaapverstoring leiden.

In de gematigde klimaatzone kan men tussen zonsondergang en zonsopgang boven land een stabiele atmosfeer verwachten als er een gedeeltelijk onbewolkte hemel is (bewolking verhindert de warmte-uitstraling) en een niet te harde wind (veel wind bevordert de verticale warmtevereffening). Uit een analyse van metingen van het KNMI bij Cabauw, in het midden van Nederland, tot op 200 m hoogte blijkt dat er een dagelijkse en jaarlijkse gang is in het windprofiel die samenhangt met de dagelijkse en seizoensvariatie in de warmte-uitwisseling tussen aardoppervlak en atmosfeer. Dat bij zonsondergang de wind vaak gaat liggen is een gevolg van de toenemende atmosferische stabiliteit, en deze windsnelheidsafname nabij de grond gaat gepaard met een toename van de windsnelheid op grotere hoogte.

Dit heeft belangrijke gevolgen voor de energieproductie van een windturbine, waarbij bovendien de rotorhoogte een rol speelt. Als wordt uitgegaan van de gemeten windsnelheden bij Cabauw op 10 m hoogte en een altijd neutrale atmosfeer, dan zou het over een jaar gemiddelde opgewekte elektrische vermogen van een 80 m hoge 2 MW windturbine bijna 500 kW bedragen. Gebaseerd op de werkelijke, gemeten windsnelheid op 80 m hoogte bedraagt het over een jaar gemiddelde vermogen echter 600 kW. Door atmosferische stabiliteit is er dus, ten opzichte van een neutrale atmosfeer, een aanmerkelijk hogere opbrengst in de nachturen, waardoor zelfs de lagere opbrengst overdag ruim wordt gecompenseerd.

De hogere windsnelheid ’s nachts op de rotor veroorzaakt echter ook een hogere geluidsproductie. Als weer wordt uitgegaan van windsnelheden op 10 m hoogte en een neutraal veronderstelde atmosfeer, dan bedraagt het geluidsvermogen van de turbine ’s nachts gemiddeld ca. 102 dB(A). In werkelijkheid is het ruim 2 dB hoger. Ook dit is een gemiddelde over een heel jaar; in afzonderlijke nachten kan het verschil veel groter zijn,
bijvoorbeeld als een windturbine op (vrijwel) topsnelheid draait, terwijl verwacht was dat deze, gezien de lage windsnelheid op 10 m hoogte, helemaal niet zou produceren. Dit gebeurt vooral in het zomerhalfjaar.

De mate waarin atmosferische stabiliteit optreedt bij Cabauw blijkt nauwelijks te verschillen van wat bij windpark Rhede is waargenomen. Op andere locaties in landen in de gematigde zone blijkt stabiliteit in vergelijkbare mate voor te komen. De beschreven gevolgen van atmosferische stabiliteit zullen dus bij veel windparken optreden die in de gematigde zone staan of nog gebouwd worden. Echter, boven grote wateroppervlakken is stabiliteit eerder een seizoens- dan een dagelijks verschijnsel, en in bergachtig gebied kunnen de gevolgen van stabiliteit op het windprofiel zowel versterkt als verzwakt worden door veranderingen tengevolge van hoogteverschillen in het gebied.

Geluid van een windturbine of windpark wordt dus om twee redenen na zonsondergang hinderlijker: het wordt luider en het geluid vertoont sterkere fluctuaties. Bij een gegeven rotordiameter kan een wiek alleen stiller worden door een ander ontwerp of door de snelheid te verlagen. Snelheidsverlaging gaat echter ten koste van het opgewekte elektrische vermogen en moet daarom liefst alleen worden toegepast wanneer dat nodig is. Daartoe kan een regeling worden toegepast die de snelheid verlaagt wanneer een geluidslimiet wordt overschreden, en deze weer verhoogt wanneer de limiet dat toelaat. De regeling zou kunnen ingrijpen op de generator en/of de vaanstand van de wieken.

Door de vaanstand tijdens de rotatie van de wieken te variëren kan op elke positie de wind onder een optimale hoek de rotor instromen, waardoor enerzijds het energetisch rendement toeneemt en anderzijds de fluctuatiesterkte van het geluid afneemt en de fluctuaties zelfs onhoorbaar kunnen worden. Het totale geluidsvermogen zal afnemen, zelfs ten opzichte van een neutrale atmosfeer, omdat het instromingsturbulentiegeluid zal verminderen door de relatieve afwezigheid van atmosferische turbulentie. Het kantelen van de rotor waardoor tijdens een rotatie de vaanstand verandert lijkt geen vruchtbare strategie: de kanteling moet zo groot zijn dat de nadelen overheersen.

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Bij een windpark kunnen de fluctuaties sterker zijn door interferentie van de fluctuaties van meerdere turbines. Dit kan worden voorkomen door de turbines te desynchroniseren, zoals dat overdag gebeurt door grootschalige atmosferische turbulentie, door kleine en ongecorreleerde variaties in de belasting van de rotors of in de vaanstand van de wieken van de afzonderlijke turbines.

Het beheersen van de geluidsproductie vergt derhalve een nieuwe strategie bij de regeling van windturbines: overdag is er vaak meer geluidsruimte beschikbaar dan ’s nachts en die ruimte kan overdag gebruikt worden als er ’s nachts beperkingen worden opgelegd.

Als laatste is nog een geheel ander probleem onderzocht: de invloed van wind op een microfoon, al of niet in een windbol. Bij voldoende wind bevat het microfoonsignaal een laagfrequent, rommelend geluid waardoor de meting van omgevingsgeluid wordt verstoord. Deze ‘rumble’ is geen geluid uit de omgeving, maar ontstaat door drukvariaties tengevolge van turbulente windsnelheidsvariaties. Met een drukgevoelige microfoon zijn deze drukvariaties niet te onderscheiden van akoestische drukvariaties. Het blijkt dat een windbol alleen effectief is doordat de bijdragen van kleine turbulente wervels worden gedempt. Een windbol heeft geen effect bij wervels die groter zijn dan de windbol.

De sterkte van atmosferische turbulentie hangt niet alleen af van de (gemiddelde) wind snelheid, maar ook van de lokale ruwheid van het aardoppervlak en de stabiliteit van de atmosfeer. De twee laatste factoren veroorzaken respectievelijk wrijvingsturbulentie en thermische turbulentie. De turbulentiesterkte is in de literatuur goed bekend bij een vrije aanstroom van wind over vlak land. De turbulentie is zwakker in een stabiele, sterker in een instabiele atmosfeer.

Het op atmosferische turbulentie gebaseerde ‘geluids’drukniveau blijkt goed overeen te komen met gemeten en gepubliceerde niveaus van door wind geïnduceerde drukniveaus. De invloed van wind op een geluidsmeting in wind kan dus worden berekend. Omgekeerd levert het rekenmodel een nieuwe methode om de sterkte van de atmosferische turbulentie te meten.
Tot slot kunnen we concluderen dat er bij het geluid van windturbines een belangrijk fenomeen over het hoofd is gezien: de verandering van de wind na zonsondergang. Dit fenomeen zal belangrijker worden voor moderne, hoge windturbines en met het oog op de vele windparken die worden gepland. Als dit probleem niet wordt onderrond en opgelost zal het de uitbreiding van windenergie bemoeilijken.
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Appendix A
List of symbols

Symbol: definition [unit]

\( \alpha \): angle of attack [radian] or [degree];
also: rotor pitch angle [radian] or [degree]
also: constant relating wind velocity to pressure [-]

\( \delta_t^* \): displacement thickness of turbulent boundary layer [m]

\( \eta_k \): Kolmogorov size [m]

\( \kappa \): von Karman’s constant [0.4]

\( \nu \): kinematic viscosity of air \([m^2\cdot s^{-1}]\)

\( \rho \): correlation coefficient \((1/3\) octave band level vs. \( L_A \)) [-];
also: air density \([kg/m^3]\)

\( \Psi(\zeta) \): stability function [-]

\( \theta \): rotor tilt angle [radian] or [degree]

\( \zeta \): dimensionless height \((h/L)\) [-]

\( \Omega \): turbine rotor angular velocity \([\text{rad}\cdot s^{-1}]\)

\( a \): induction factor \((1 - V_b/V_h)\) [-]

\( b \): correction factor for boundary layer thickness (value: \( 2 - 4 \))

\( c \): velocity of sound in air \([m\cdot s^{-1}]\)

\( C \): blade chord length [m]; also: air density dependent constant
\[ C = 20 \cdot \log(0.215 \kappa \alpha \rho V_o^2 / p_{ref}) \text{ [dB]} \]

\( C_p \): constant \((C_p = 20 \cdot \log(0.215 \kappa \alpha) - 9.5) \text{ [dB]} \)

\( D \): diameter [m]

\( D_\eta \): directivity function [-]

\( D_{j,k} \): decrease in octave band sound level \( j \) of turbine \( k \) with distance [dB]

\( D_{\text{geo}} \): decrease in sound level due to geometrical spreading [dB]

\( D_{\text{air}} \): decrease in sound level due to air absorption [dB]

\( D_{\text{ground}} \): decrease in sound level due to ground absorption and reflection [dB]

\( \text{dB(A)} \): unit of level after A-weighting

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dB(G): unit of level after G-weighting

\( f \): frequency [Hz]

\( f_{mod} \): modulation frequency [Hz]

\( f_{peak,TE} \): peak frequency of trailing edge sound [Hz]

\( f_{peak,if} \): peak frequency of in-flow turbulence sound [Hz]

\( f_m \): middle frequency of 1/3 octave band [Hz]

\( f_b \): blade passing frequency [Hz]

\( f_c \): screen size related corner frequency \((f_c = 0.3V_1D)\) [Hz]

\( f_\alpha \): \( \alpha \)-dependent factor for TE layer thickness [-]

\( f_{\text{log}} \): ratio \( v_{98}/v_{10} \) valid in a neutral atmosphere [-]

\( f_{(un)stable} \): ratio \( v_{98}/v_{10} \) valid in an (un)stable atmosphere [-]

\( F_{bb} \): fluctuation strength [vacil]

\( F(z) \): turbulence related function:

\[
F(z) = -20 \cdot \log \left( (z/D)^{1/3} \cdot (\ln(z/z_0) - \Psi) \right) \text{ [dB]}
\]

\( G(z) \): turbulence related function:

\[
G(z) = -20 \cdot \log \left( 0.2 \cdot (z/\ell_o)^{1/3} \cdot (\ln(z/z_0 - \Psi) \right) \text{ [dB]}
\]

\( h \): height [m]

\( H \): turbine height [m]

\( h_{\text{ref}} \): reference height for wind velocity (and direction) [m]

\( k \): integer number (of harmonic frequency) [-]; also: exponent of wind velocity in relation with associated turbulent pressure [-]

\( K_1 \): constant (128.5 dB)

\( K_\alpha \): \( \alpha \)-dependent increase in trailing edge sound level [dB]

\( \ell \): eddy length scale [m]

\( \Delta L \): increase in sound level [dB]

\( L \): Monin-Obukhov length [m]

\( L_A \): broad band A-weighted sound level [dB(A)]

\( L_{AS} \): 5-percentile of broad band sound levels over a period [dB(A)]

\( L_{AS95} \): 95-percentile of broad band sound levels over a period [dB(A)]

\( L_{at}(u) \): pressure level due to atmospheric turbulence [dB]

\( L_{at,1/1(f)} \): pressure level due to turbulent wind per octave band [dB]

\( L_{at,1/3(f)} \): pressure level due to turbulent wind per 1/3 octave band [dB]

\( L_{at,A} \): broad band A-weighted pressure level [dB]

\( L_{imm} \): immission sound level [dB(A)]

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\( L_{eq}: \) equivalent sound level; \( L_{eq,T}: \) over time \( T \) [dB(A)]

\( L_{p,1/3}: \) turbulent pressure level at microphone per 1/3 octave band [dB]

\( L_{red,1/3}: \) ‘meteorologically reduced’ 1/3 octave band pressure level [dB]

\( L_{red,1/1}: \) ‘meteorologically reduced’ octave band pressure level [dB]

\( L_w: \) sound power level [dB(A)]

\( L_{wj}: \) j-th octave band sound power level [dB(A)]

\( M: \) Mach number = air flow velocity/c (at radius \( R: \) \( M = \Omega R/c \)) [-]

\( m: \) stability exponent [-]

\( m_{h1,h2}: \) \( m \) determined between heights \( h_1 \) and \( h_2 \) [-]

\( mf: \) modulation factor [-]

\( n: \) dimensionless frequency (\( n = f z/V \)) [-]

\( N: \) number of blades [-]; rotational speed (\( \Omega R/2\pi \)) [s\(^{-1}\)]

\( Ph: \) Power at height \( h; \) \( Ph,lpp; Ph,hp \) [W]

\( p: \) (sound) pressure [Pa]

\( p_{0}: \) rms pressure in narrow frequency band centered at frequency \( f \) [Pa]

\( p_{0.1/3}: \) rms pressure in 1/3 octave band [Pa]

\( p_{\text{ref}}: \) reference (sound) pressure [20 \( \mu \)Pa]

\( p(0): \) rms pressure at center of wind screen [Pa]

\( r: \) distance [m]

\( R: \) rotor radius = blade length [m]

\( \Delta R: \) increment in \( R \) [m]

\( R_X: \) range between maximum and minimum sound levels

\( (X= bb \text{ or } f) \) [dB]

\( R_{X,90}: \) range between 5- and 95-percentile of sound levels

\( (X= bb \text{ or } f) \) [dB]

\( Re: \) chord based Reynolds number (\( Re = \Omega RC/v \)); wind screen diameter based Reynolds number [-]

\( S: \) ratio of distance along blade and chord length [-]

\( Sp: \) 1/3 octave band weighing function for TE sound [dB]

\( SPL_i: \) sound pressure level of source \( i \) [dB]

\( Sr: \) Strouhal number [-]

\( u: \) longitudinal (along wind) component of turbulent wind velocity [m\cdot s\(^{-1}\)]
\( u_f \): rms longitudinal component of turbulent wind velocity per unit frequency [\( \text{m} \cdot \text{s}^{-3/2} \)]

\( u^* \): friction velocity [\( \text{m} \cdot \text{s}^{-1} \)]

\( U \): instantaneous wind velocity: \( U = <U> + u \) [\( \text{m} \cdot \text{s}^{-1} \)]

\( V \): air flow velocity or wind velocity [\( \text{m} \cdot \text{s}^{-1} \)]

\( V_o \): reference velocity [\( \text{1 m} \cdot \text{s}^{-1} \)]

\( V_b \): induced wind velocity at turbine blade [\( \text{m} \cdot \text{s}^{-1} \)]

\( V_h, V_{xx} \): wind velocity at height \( h \) or height \( xx \) m [\( \text{m} \cdot \text{s}^{-1} \)]

\( V_{h,b}, V_{xx,b} \): induced wind velocity at turbine blade or height \( h \) [\( \text{m} \cdot \text{s}^{-1} \)]

\( V_{hub} \): wind velocity at wind turbine hub height \( h \) [\( \text{m} \cdot \text{s}^{-1} \)]

\( V_i \): local (induced) velocity at blade \( \approx 2V/3 \) [\( \text{m} \cdot \text{s}^{-1} \)]

\( V_{ref} \): wind velocity at reference height [\( \text{m} \cdot \text{s}^{-1} \)]

\(<x>\): time average of variable \( x \)

\( z_o \): roughness height; altitude [\( \text{m} \)]

Subscripts:

1/1: frequency octave band

1/3: 1/3 frequency octave band

A: A-weighted

at: atmospheric turbulence

bb: broad band

f: at frequency of (1/3) octave band

h: at height \( h \), hub

i: component of TE sound (\( i = p, s, a \))

if: in-flow

p: pressure, pressure side

ref: reference

s: suction side

TE: trailing edge

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Appendix B

Dominant sources of wind turbine sound

With modern wind turbines there are three important mechanisms that produce sound. These will be reviewed here up to a detail that is relevant to the text in this book.

B.1 Infrasound: thickness sound

When a blade moves through the air, the air on the forward edge is pushed sideways, moving back again at the rear edge. For a periodically moving blade the air is periodically forced, leading to ‘thickness sound’. Usually this will not lead to a significant sound production as the movement is smooth and thus accelerations relatively small.

When a blade passes the turbine tower, it encounters wind influenced by the tower: the wind is slowed down, forced to move sideways around the tower, and causes a wake behind the tower. For a downwind rotor (i.e. the wind passes the tower first, then the rotor) this wake causes a significant change in blade loading.

The change in wind velocity near the tower means that the angle of attack of the air on a blade changes and lift and drag on the blade change more or less abruptly. This change in mechanical load increases the thickness sound power level at the repetition rate of the blade passing frequency $f_B$. For modern turbines $f_B = N \cdot \Omega / (2\pi)$ typically has a value of approximately 1 Hz. As the movement is not purely sinusoidal, there are harmonics with frequencies $k \cdot f_B$, where $k$ is an integer. Harmonics may occur up to 30 Hz, so thickness sound coincides with the infrasound region (0–30 Hz).

Measured levels at 92 m from the two-bladed 2 MW WTS-4 turbine showed that measured sound pressure levels of the individual blade harmonics were less than 75 dB, and well predicted by calculations of wind-blade interaction near the turbine tower [Hubbard et al 2004, Wagner et al 1996]. The envelope of the harmonics peaks at the fifth harmonic ($k = 5$ with $f_B = 1$ Hz), indicating a typical pulse time of $(5 \text{ Hz})^{-1} = 0.2$ s which is 20% of the time between consecutive blade passages. The WST-4 is a
downwind turbine with an 80 m tubular tower, where the wind velocity deficit was estimated to be 40% of the free wind velocity [Hubbard et al, 2004]. For modern, upwind rotors the velocity deficit in front of the tower is smaller. As a consequence the change in blade loading is less than for downwind turbines. From data collected by Jakobsen it appears that the infrasound level at 100 m from an upwind turbine is typically 70 dB(G) or less, whereas near downwind turbines it is 10 to 30 dB higher. As 95 dB(G) corresponds to the average infrasound hearing threshold [Jakobsen, 2004], infrasound from (upwind) wind turbines does not appear to be so loud that it is directly perceptible.

5.2 Low frequencies: in-flow turbulent sound

Because of atmospheric turbulence there is a random movement of air superimposed on the average wind velocity. The contribution of atmospheric turbulence to wind turbine sound is named ‘in-flow turbulence sound’ and is broad band sound stretching over a wide frequency range. For turbulent eddies larger in size than the blade this may be interpreted as a change in the direction and/or velocity of the incoming flow, equivalent to a deviation of the optimal angle of attack. This leads to the same phenomena as described in section B.1, but changes will be random (not periodic) and less abrupt. For turbulent eddies the size of the chord length and less, effects are local and do not occur coherently over the blade. When the blade cuts through the eddies, the movement normal to the wind surface is reduced or stopped, given rise to high accelerations and thus sound.

In-flow turbulence sound has a maximum level in the 1/3 octave band with frequency

\[ f_{peak,if} = \frac{(St \cdot 0.7R \cdot \Omega)}{(H - 0.7R)} \]  

(B.1)

where Strouhal number St is 16.6 [Grosveld, 1985, Wagner et al, 1996]. Most sound is produced at the high velocity, outer parts of the blades. For a modern, tall, three-bladed wind turbine with hub height \( H = 100 \) m, blade length \( R = 35 \) m and angular velocity \( \Omega = 2\pi f_0 / 3 = 2 \) rad/s \( (20 \) rpm), \( f_{peak,if} = 11 \) Hz which is in the infrasound region. Measured fall-off from \( f_{peak,if} \) is
initially approx. 3 dB per octave, increasing to 12 dB per octave at frequencies in the audible region up to a few hundreds of hertz [Grosveld 1985, Wagner et al 1996].

B.3 High frequencies: trailing edge sound

Several flow phenomena at the blade itself or in the turbulent wake behind a blade cause high frequency sound (‘airfoil self-noise’). Most important for modern turbines is the sound from the turbulent boundary layer at the rear of the blade surface where the boundary layer is thickest and turbulence strength highest. Trailing edge sound has a maximum level in the 1/3 octave band with frequency

$$f_{\text{peak,TE}} = 0.02 \cdot \Omega \cdot R / (\delta^* \cdot M^{0.6})$$  \hspace{1cm} (B.2)

where Mach number $M$ is based on airfoil velocity. The displacement thickness of the turbulent boundary layer is:

$$\delta^* = b \cdot 0.37 \cdot C \cdot \text{Re}^{-0.2}/8$$  \hspace{1cm} (B.3)

for a zero angle of attack. Re is the chord based Reynolds number [Brooks et al 1989]. The experimental factor $b$ accounts for the empirical observation that the boundary layer is a factor 2 to 4 thicker than predicted by theory [Lowson 1995, Wagner et al 1996]. For air of 10 °C and atmospheric pressure, a typical chord length $C = 1$ m, and other properties as given above (section B.2), $f_{\text{peak,TE}} = 1700/a$ Hz. With $b = 2$ to 4, $f_{\text{peak,TE}}$ is 450 – 900 Hz. The spectrum (see $S_p$ below) is symmetrical around $f_{\text{peak,TE}}$ and decreases with 3 dB for the first octave, 11 dB for the next; the contribution from further octave bands is negligible [Brooks et al 1989].

According to Brooks et al [1989] trailing edge sound level can be decomposed in components $S P L_p$ and $S P L_s$ due to the pressure and suction side turbulent boundary layers with a zero angle of attack of the incoming flow, and a component $S P L_\alpha$ that accounts for a non-zero angle of attack $\alpha$. For an edge length $\Delta R$ each of the three components of the immission sound level at distance $r$ can be written as [Brooks et al 1989]:

$$S P L_i = 10 \cdot \log(\delta_i^* \cdot M^5 \cdot \Delta R \cdot D_h / r^2) + S p_i + K_i - 3 + K_i$$  \hspace{1cm} (B.4)

B-3
and total trailing edge immission sound level as:

\[ \text{SPL}_{\text{TE}} = 10 \cdot \log(\Sigma_{i} \cdot 10^{\text{SP}_{i}/10}) \]  

(B.5)

where the index \( i \) refers to the pressure side, suction side or angle of attack part \((i = p, s, \alpha)\). The directivity function \( D_{h} \) equals unity at the front of the blade \((\theta = 180^\circ)\) and falls off with \( \sin^{2}(\theta/2) \). Because of the strong dependence on \( M \) (~ \( M^{5} \), equation B.4) trailing edge sound is dominated by sound produced at the high velocity parts: the blade tips.

\( \text{SP}_{i} \) gives the symmetrical spectral distribution of the trailing edge sound spectrum centered on \( f_{\text{peak,TE}} \) and is maximum \((0 \, \text{dB})\) at this centre frequency. The constant \( K_{1} - 3 = 125.5 \, \text{dB} \) applies when the chord based Reynolds number exceeds \( 8 \cdot 10^{5} \) and the pressure-side turbulent boundary displacement thickness \( \delta_{i} > 1 \, \text{mm} \), as is the case for modern tall turbines. \( K_{i} \) is non-zero only if \( i = \alpha \).

For positive angles of attack \( \alpha < 10^\circ \) the boundary layer thickness \( \delta_{*} \) shrinks with a factor \( f_{p} = 10^{-0.042\alpha} \) at the pressure-side and \( \delta_{*} \) grows at the suction-side with a factor \( f_{s} = 10^{0.068\alpha} \). Because \( \delta_{*} = \delta_{*} \), \( f_{p} = f_{s} \). \( K_{\alpha} \) has a large negative value for \( \alpha = 0 \). For \( 1^\circ < \alpha < 10^\circ \) and \( M = 0.2 \) the calculated values of \( K_{\alpha} \) (see formula 49 in [Brooks et al. 1989]) with \( K_{\alpha} = K_{2}-K_{1}+3 \) are plotted in figure B.1 and these can be approximated by:

\[ K_{\alpha} = -0.35 \cdot \alpha^{2} + 5.5 \cdot \alpha - 14.4 \quad (\alpha \text{ in degrees}) \]  

(B.6)

With equation B.4, equation B.5 can be rewritten as:

\[ \text{SPL}_{\text{TE}} = 10 \cdot \log(\delta_{*}^{*} \cdot M^{5} \cdot \Delta R \cdot D_{h}/r^{2}) + K_{1} - 3 + 

+ 10 \cdot \log(\Sigma_{i} \cdot 10^{(10 \cdot \log(f) + \text{SP}_{i} + K_{i})/10}) \]  

(B.7)

The last term in B.7 is the \( \alpha \)-dependent part. For the peak frequency 1/3 octave band level \( (\text{SP}_{i} = 0) \) the last term in equation B.7 is 3 \( \text{dB} \) for \( \alpha = 0 \) and 3.4 \( \text{dB} \) for \( \alpha = 1^\circ \), then increasing with 1.5 \( \text{dB} \) per degree to 14.5 \( \text{dB} \) at \( \alpha = 9^\circ \). The level increase \( \Delta \text{SPL}_{\text{TE}}(\alpha) = \text{SPL}_{\text{TE}}(\alpha) - \text{SPL}_{\text{TE}}(\alpha = 0) \) is given in table B.1 and plotted in figure B.1. The best linear approximation in the range \( 1^\circ < \alpha < 10^\circ \) is:

B-4
\[ \Delta SPL_{TE}(\alpha) = 1.5 \cdot \alpha - 1.2 \text{ (dB)} \]  

(B.8)

with \( \alpha \) in degrees (or \( \Delta SPL_{TE}(\alpha) = 86 \cdot \alpha - 1.2 \text{ dB with } \alpha \text{ in radians} \)).

**Table B1: increase of trailing edge sound level with angle of attack \( \alpha \)**

<table>
<thead>
<tr>
<th>( \alpha )</th>
<th>1°</th>
<th>2°</th>
<th>3°</th>
<th>4°</th>
<th>5°</th>
<th>6°</th>
<th>7°</th>
<th>8°</th>
<th>9°</th>
</tr>
</thead>
<tbody>
<tr>
<td>( \Delta SPL_{TE}(\alpha) ) (dB)</td>
<td>0.4</td>
<td>1.4</td>
<td>2.9</td>
<td>4.6</td>
<td>6.4</td>
<td>8.0</td>
<td>9.4</td>
<td>10.6</td>
<td>11.5</td>
</tr>
</tbody>
</table>

The blade swish that is audible near a turbine is a variation in level of less than 3 dB (in daytime) [ETSU 1996]. It must correspond to a change in sound level of 1 dB to be heard at all. An increase of 1 dB corresponds to an increase in \( \alpha \) with 0.7°, an increase of 3 dB corresponds to 2.9°. So, for a swish level of 2 ± 1 dB, we estimate the change in \( \alpha \) at the tower passage as 1.8° ± 1.1°. Part of this is due to the lower wind velocity at the lower blade tip relative to the rotor average, the rest is due to the slowing down of the wind by the tower.

For small angles the change of wind velocity with angle of attack \( \alpha \) at radius \( R \) is \( dV_{wind} = \Omega \cdot R \cdot d\alpha \), or

\[ dV_{wind} = 0.017 \cdot \Omega \cdot R \cdot d\alpha \]  

(B.9)

with \( \alpha \) in degrees.
So for a modern turbine at high speed (Ω·R ≈ 70 m/s at tip at 20 rpm) the wind velocity deficit where the blade tip passes the tower and α = 1° (0.017 radians) is 1.2 m/s. In a free 14 m/s wind, i.e. 9.3 m/s at the rotor, this is 13%. This deficit is due to the influence of the tower as well as the (daytime) wind profile.
Appendix C

Simultaneous registrations of sound immission level

Additional information to section IV.10: measurements at locations A and P through X (see map figure IV.2) in year 2002. Graphs show measured values of $L_{eq,5\text{min}}$ at locations near Rhede wind farm and differences relative to measured value at location A. Wind velocity and wind direction and time of measurement are mentioned in the figures.
Appendix D
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Waar of niet: mogelijk of vermeend, NVS-Nieuws 18e jaargang nr. 2 (april 1993) pp. 8-9 (reactie op artikel)
Noise from the Marnewaard shooting range; a review of sound and annoyance measurements, proceedings Internoise93, Leuven, pp. 1145-1148
A home kit for road traffic noise, proceedings Euronoise95, Lyon, pp. 163-168
Laagfrequent geluid-een onderschat probleem, Geluid mrt 1996, pp.14-18
Straling - één pot nat; Psychologische effecten van elektromagnetische straling, NVS-Nieuws, 22e jaargang nr.4, oktober 1997, pp. 11-14
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Long range outdoor propagation and interference of low frequency tonal sound, proceedings Internoise 1998, Christchurch
Case control study in low frequency sound measurements, proceedings Internoise 1999, Fort Lauderdale
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Windturbines: een verschil van dag en nacht, Geluid, jaargang 27, nr. 1 (2004)
Do wind turbines produce significant low frequency sound levels?, proceedings 11th International Meeting on Low Frequency Noise and Vibration and its Control, Maastricht (2004)
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Mitigation measures for nighttime wind turbine noise, proceedings First International Meeting on Wind Turbine Noise: Perspectives for Control, Berlin (2005)


Monitoring van geluid in stille gebieden, proceedings NAG-lezingendag, Utrecht (2005)


D1.2 Co-author


Penetration of radon from crawl space to higher levels in dwelling, E. Veermans, G.P. van den Berg, R.J. de Meijer en L.W. Put; Radiation Protection Dosimetry vol.30 (1), pp. 45-50 (1990)


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Stellingen behorende bij het proefschrift “The sound of high winds: the effect of atmospheric stability on wind turbine sound and microphone noise” van Godefridus Petrus van den Berg

1. Exploitanten van windparken weten verbazend weinig van wind.

2. Het ‘klappen’ van de wieken is ’s nachts het hinderlijkste aspect van windturbinegeluid en vereist een net zo voortvarende aanpak als eerder het geval was bij het machinale, tonale geluid van windturbines.

3. Het optimaliseren van windparken op zoveel mogelijk opgewekte energie én zo weinig mogelijk overlast voor omwonenden is noodzakelijk om de exploitatie van windenergie op land te laten slagen.

4. De meest effectieve manier om bij geluidsmetingen aan windturbines het windgeluid op de microfoon te reduceren is over het hoofd gezien: ’s nachts meten.

5. Het materiaal van een windbol voor een microfoon doet er niet toe, als het maar geluid doorlaat en wind tegenhoudt.

6. Van teveel luisteren naar hun opdrachtgevers worden geluidsadviseurs slechthorend.

7. Van marktwerking in het openbaar vervoer is alleen de overheid beter geworden.

8. Duurzame groei van welvaart bestaat niet.


10. Ingewikkelde wetgeving is goed voor de werkgelegenheid van adviseurs.